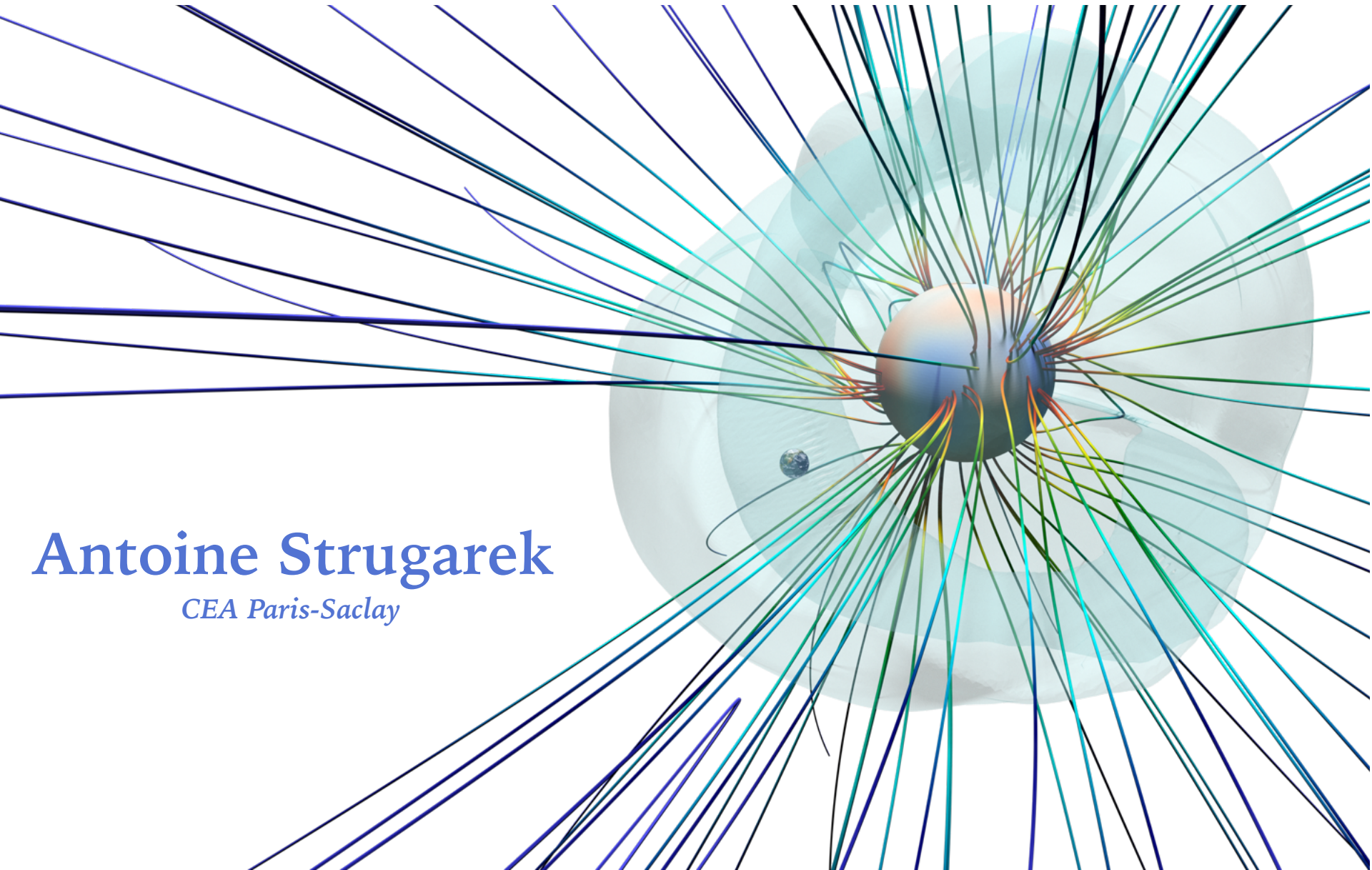


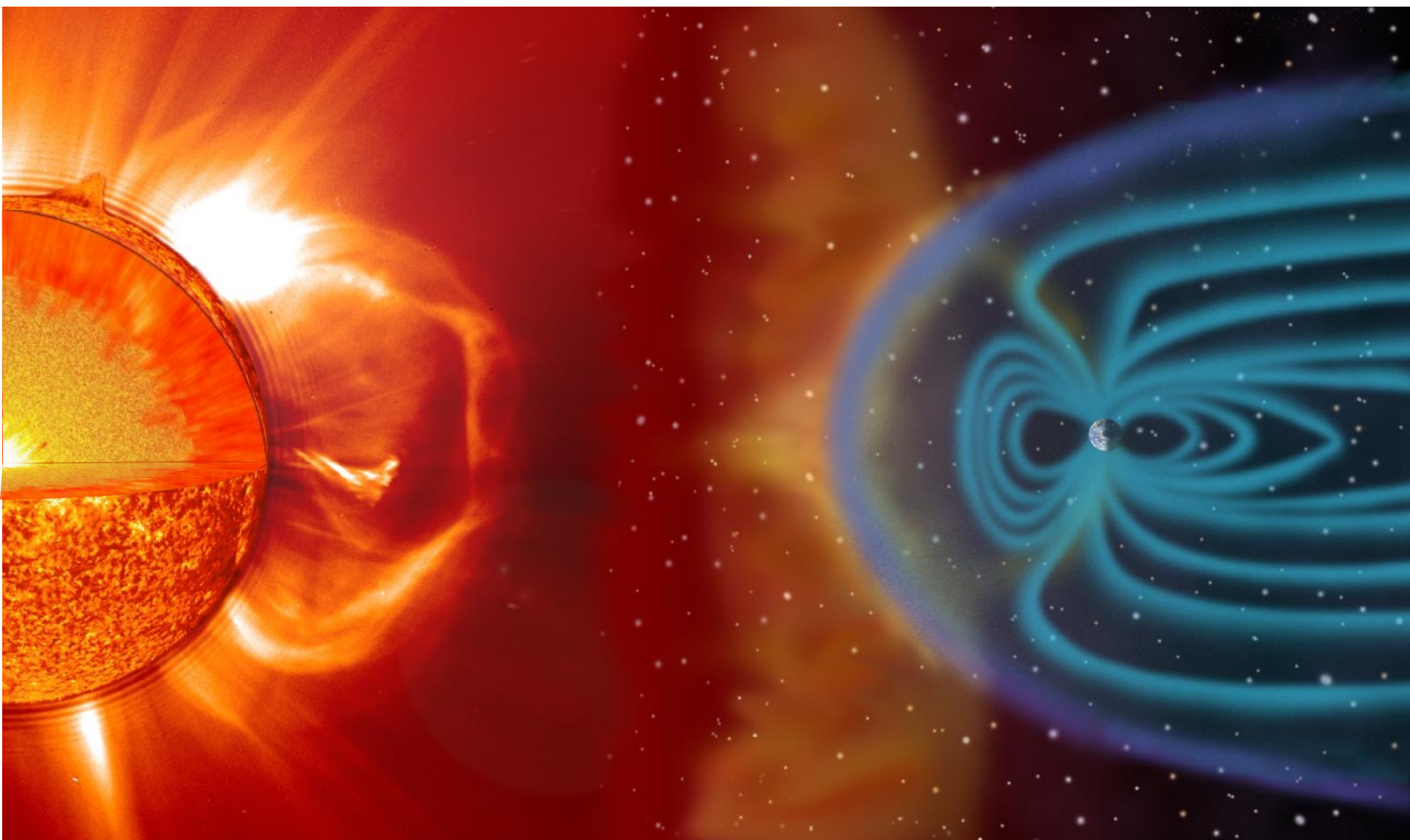
# Star-planet magnetic interactions



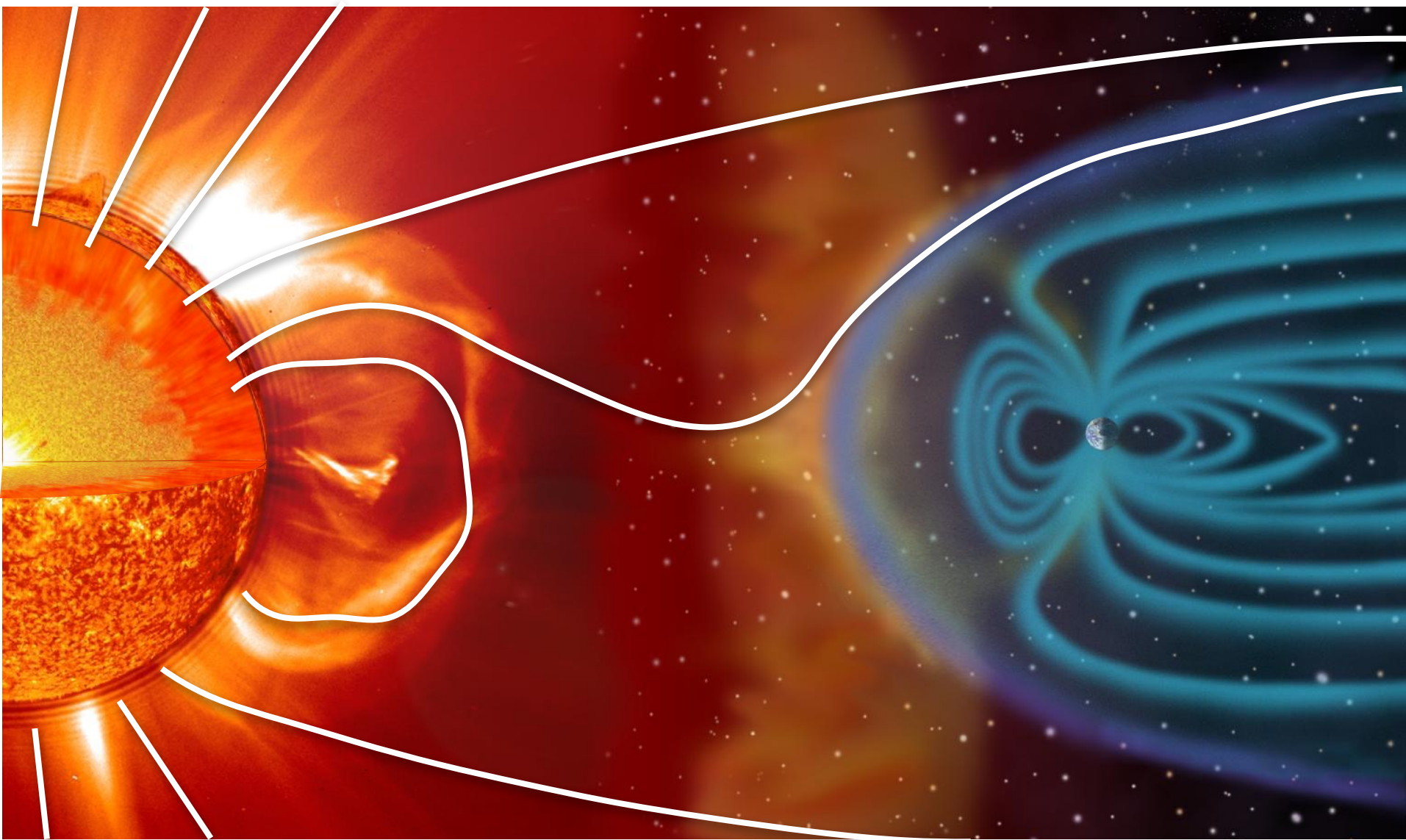
**Antoine Strugarek**

*CEA Paris-Saclay*

# What are star-planet magnetic interactions?



# What are star-planet magnetic interactions?



# Outline

## I. Setting the scene

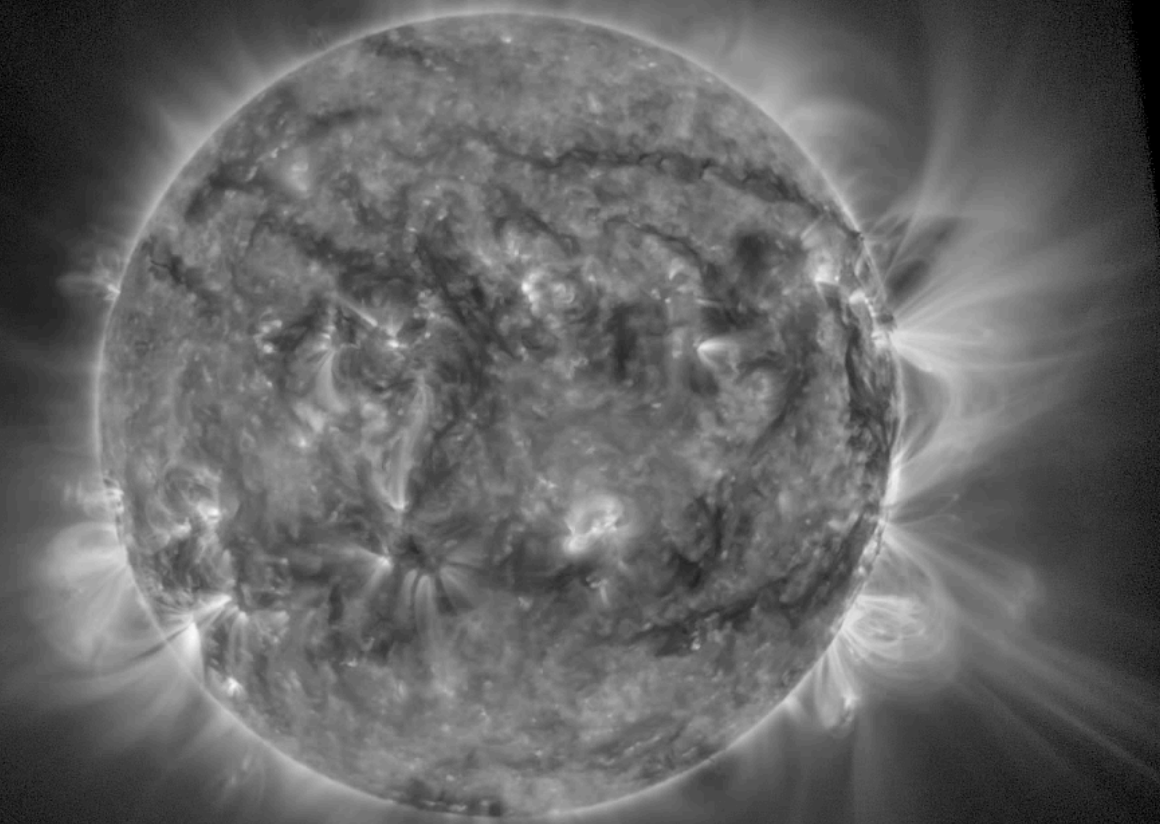
A. Stellar magnetism

B. Stellar winds

C. Planetary magnetism

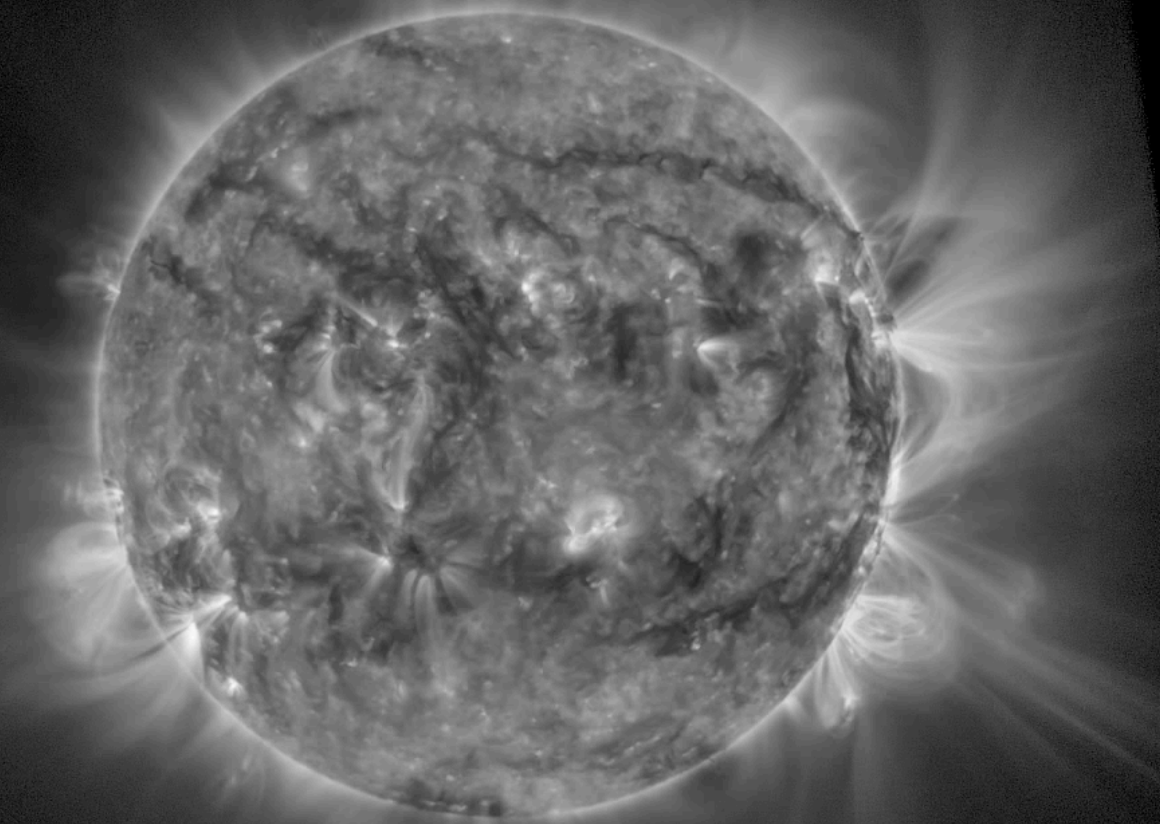
## II. Star-planet magnetic interactions

# Solar Magnetism: what do we know?



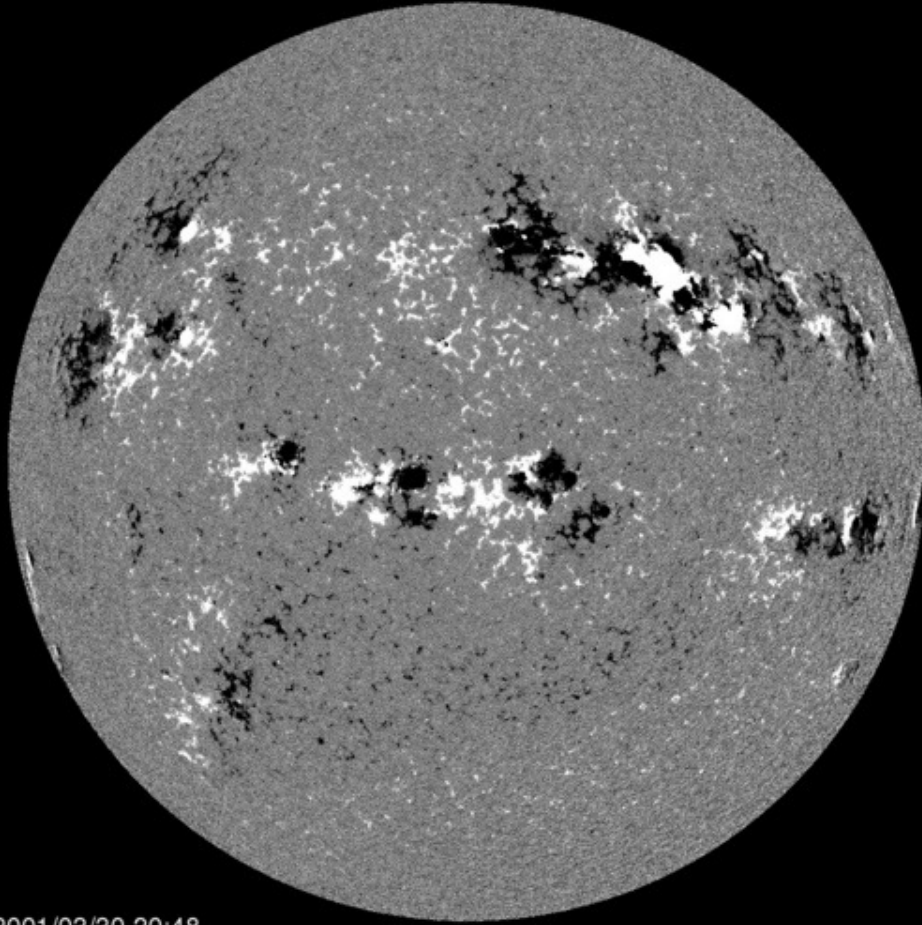
SWAP/PROBA2 17.4 nm 2012-06-21 06:10:32 CR 2125

# Solar Magnetism: what do we know?

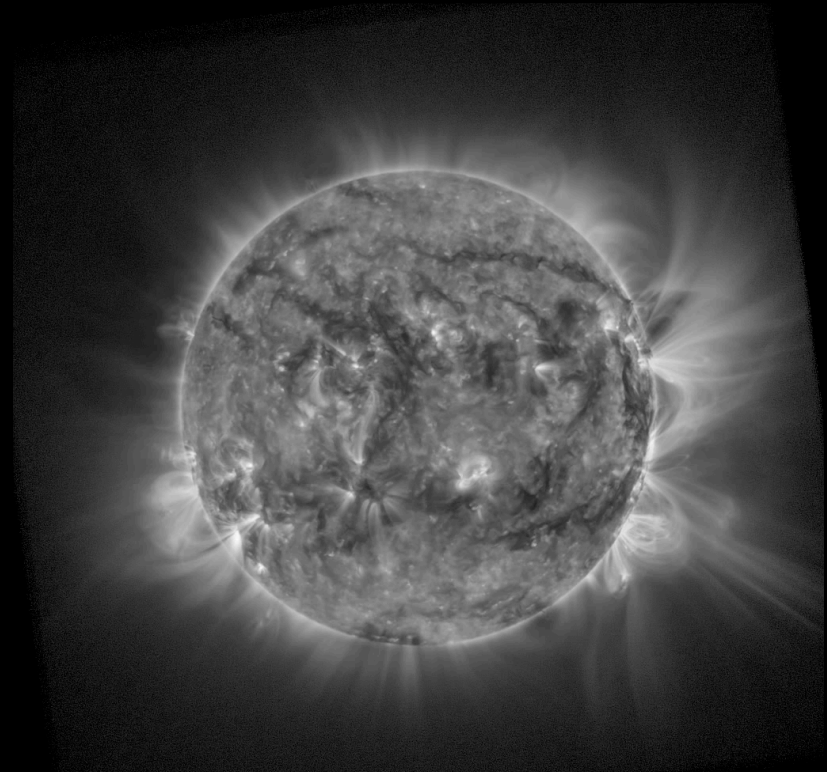


SWAP/PROBA2 17.4 nm 2012-06-21 06:10:32 CR 2125

# Solar Magnetism: what do we know?



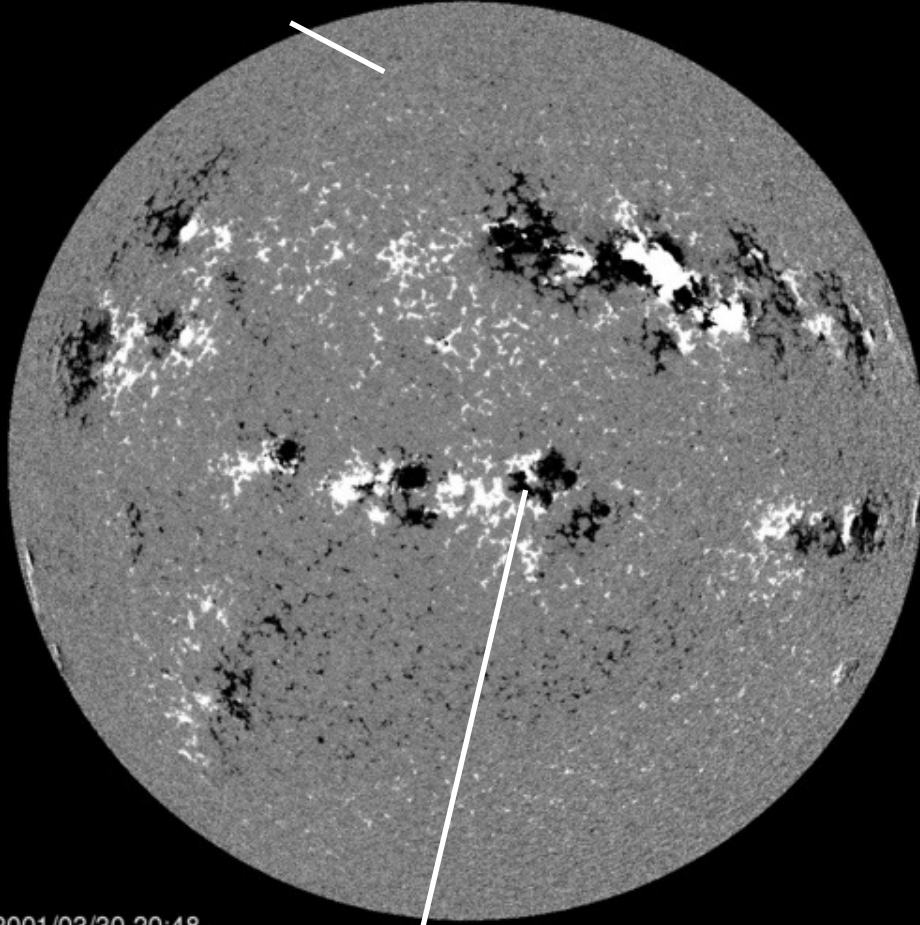
2001/03/30 20:48



SWAP/PROBA2 17.4 nm 2012-06-21 06:10:32 CR 2125

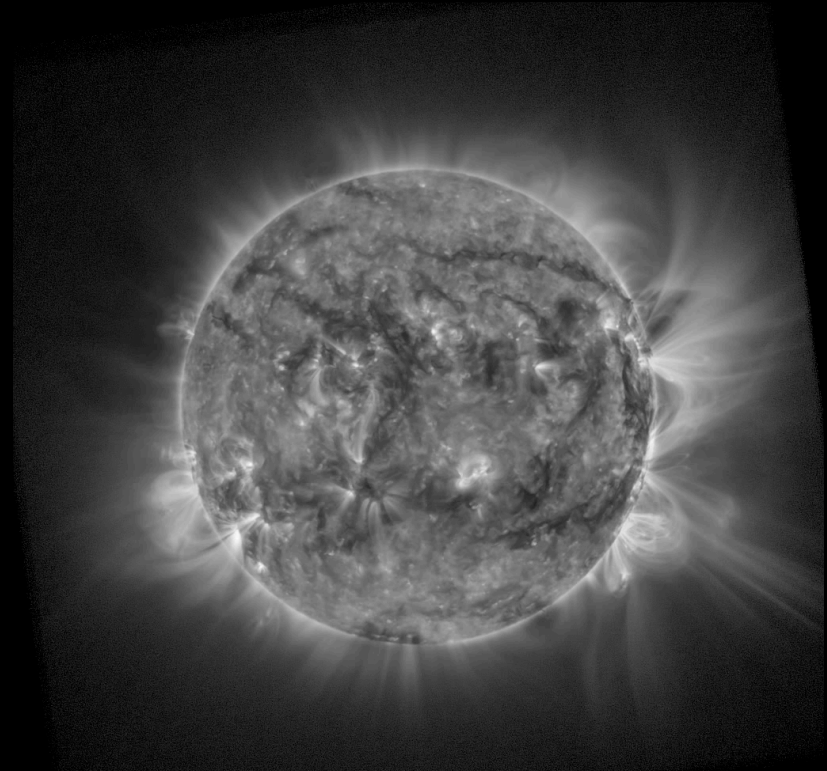
# Solar Magnetism: what do we know?

Large scale field  $\sim 2\text{-}5\text{ G}$



2001/03/30 20:48

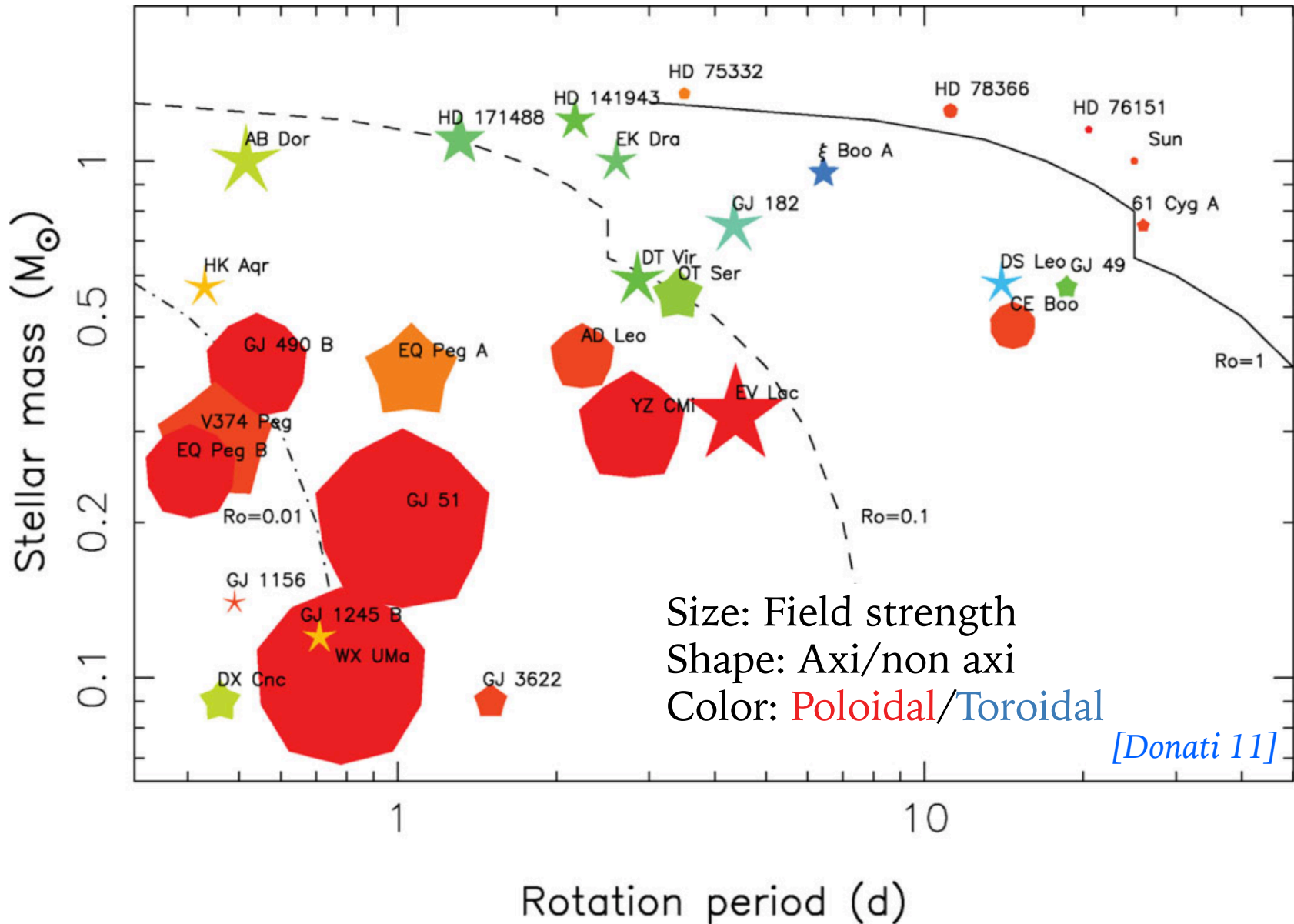
In active regions  $\sim 10^4\text{ G}$



SWAP/PROBA2 17.4 nm 2012-06-21 06:10:32 CR 2125

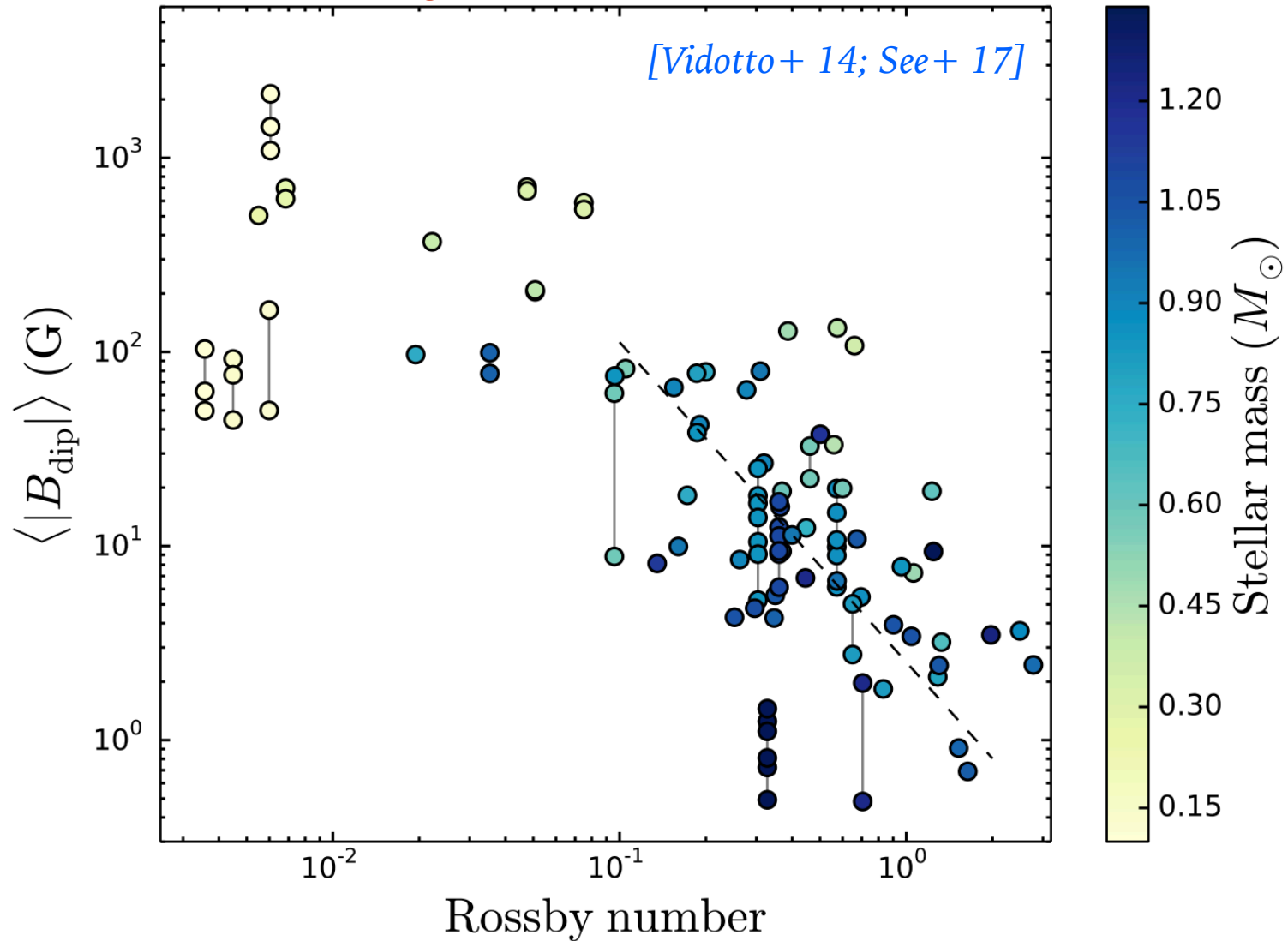


# Stellar Magnetism: what do we know?



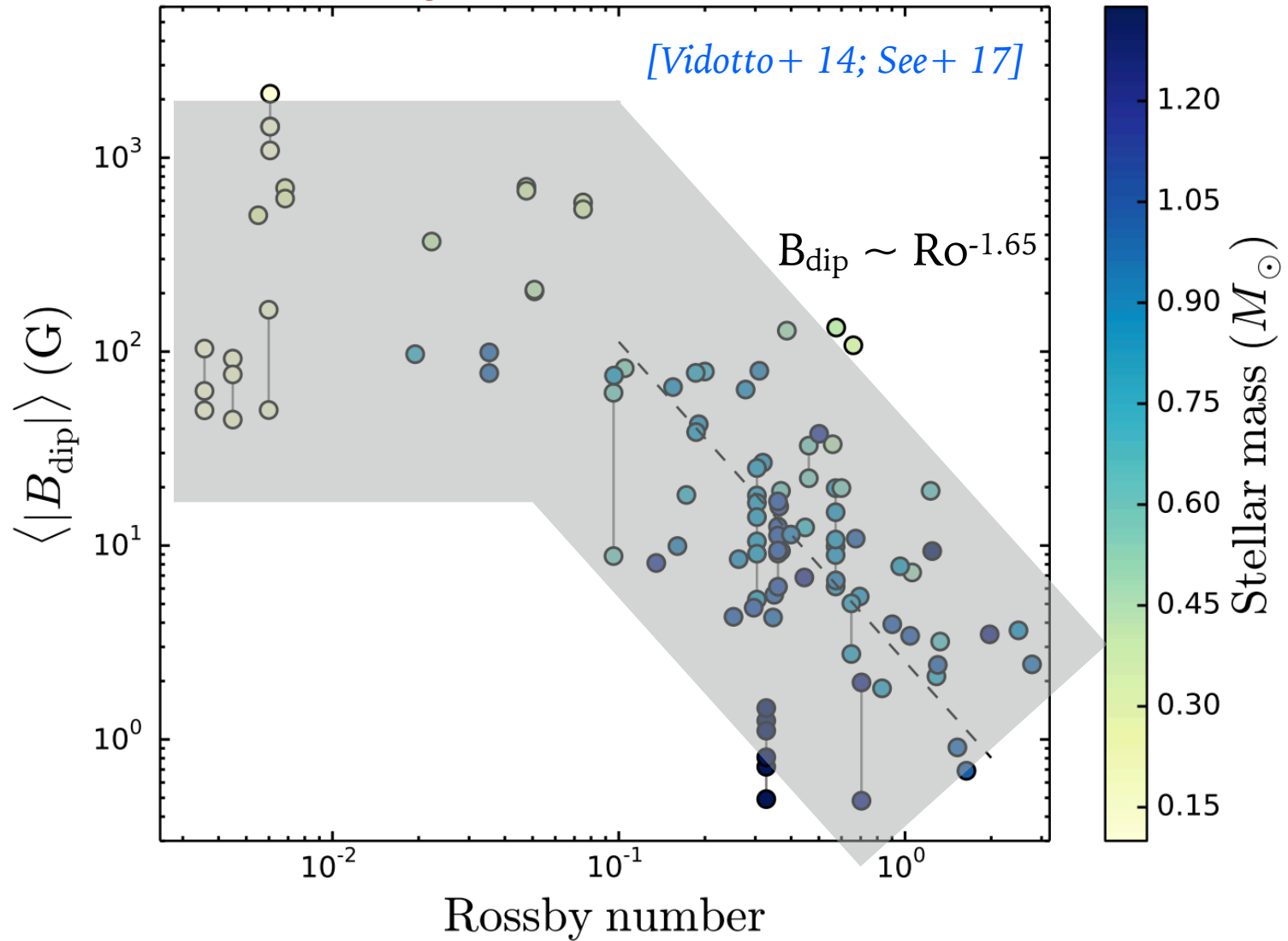
# Stellar Magnetism: what do we know?

*Caution: strong bias in stellar mass*



# Stellar Magnetism: what do we know?

*Caution: strong bias in stellar mass*



# Outline

## I. Setting the scene

A. Stellar magnetism

**B. Stellar winds**

C. Planetary magnetism

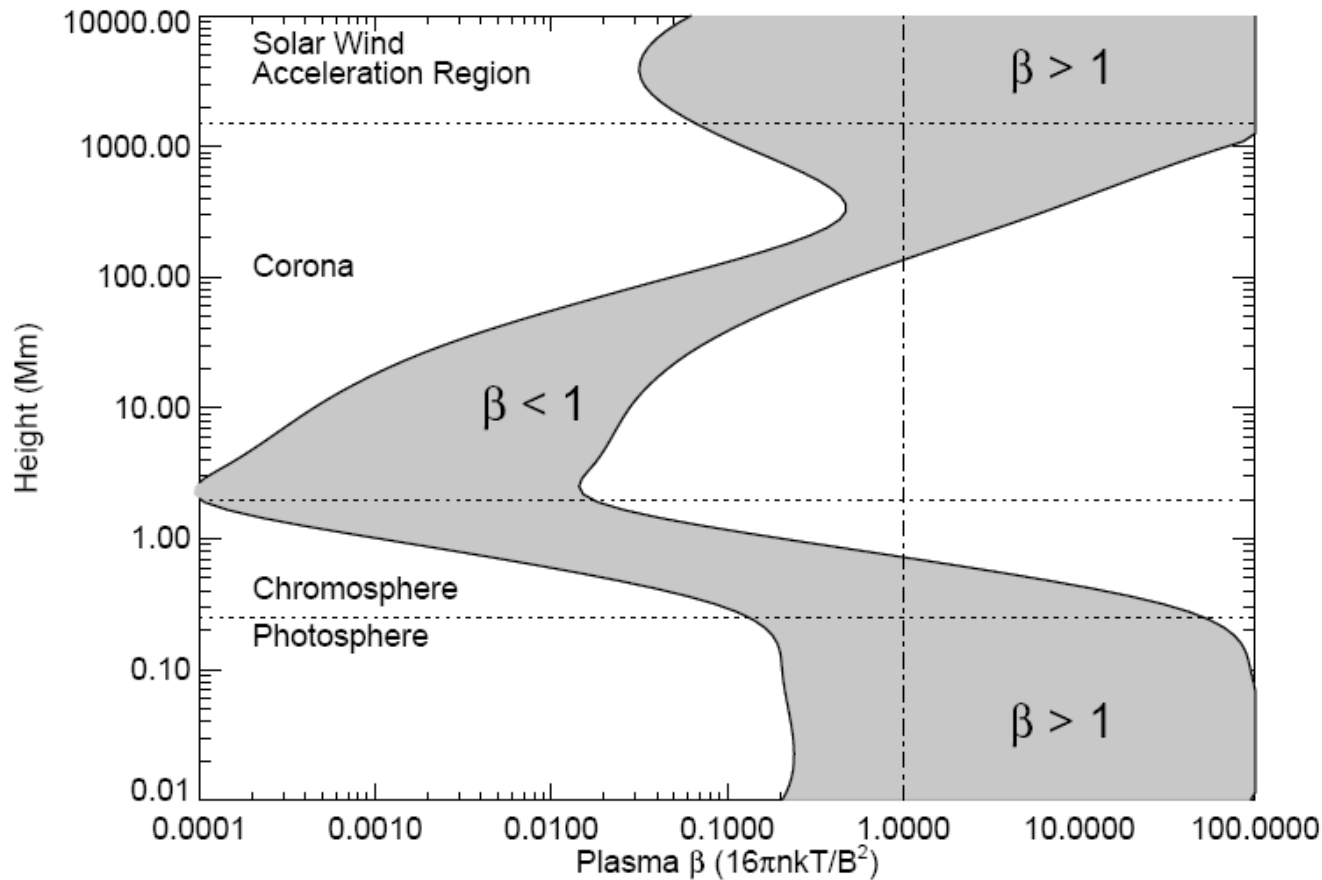
## II. Star-planet magnetic interactions

# Stellar magnetism sculpts the stellar environment (I)

Total solar eclipse in Atacama desert, 2<sup>nd</sup> July 2019



# Stellar magnetism sculpts the stellar environment (II)



$$\beta = \frac{8\pi nk_B T}{B^2}$$

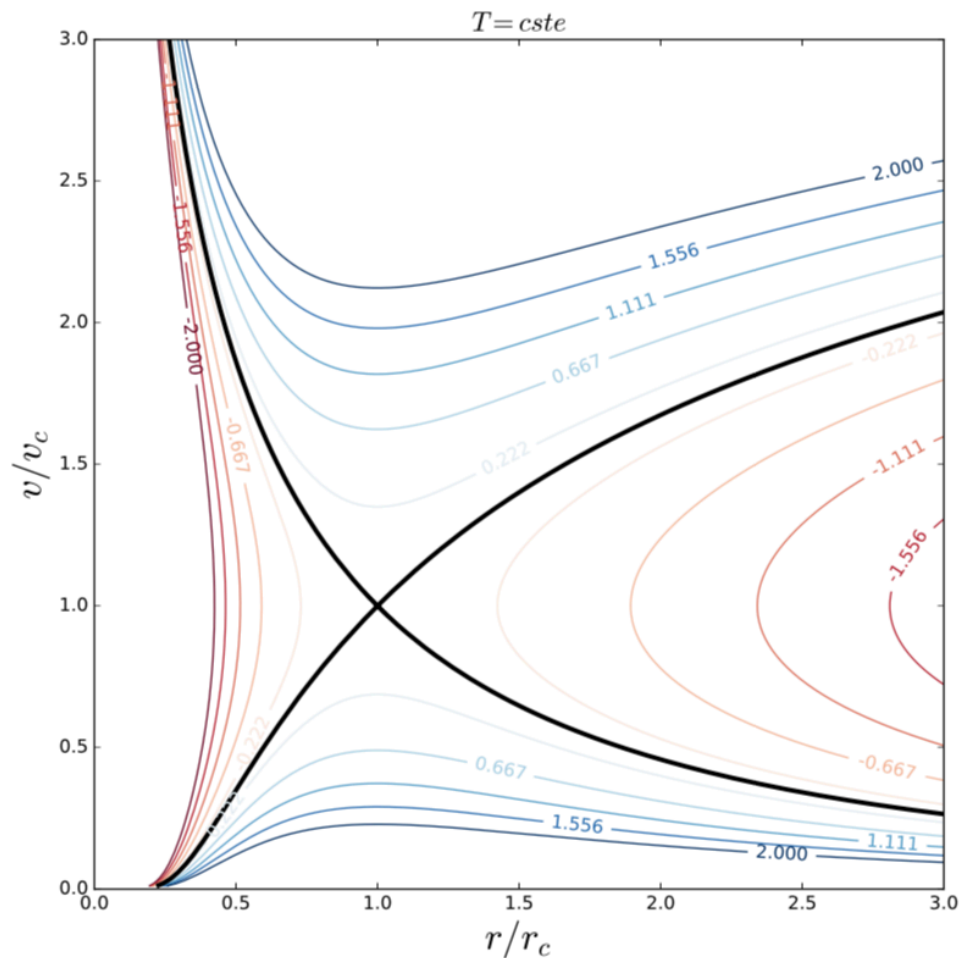
Figure 1.22: Plasma  $\beta$  in the solar atmosphere for two assumed field strengths, 100 G and 2500 G. In the inner corona ( $R \lesssim 0.2R_\odot$ ), magnetic pressure generally dominates static gas pressure. As with all plots of physical quantities against height, a broad spatial and temporal average is implied (Gary, 2001).

# The winds from cool stars

Let us assume a static ( $u=0$ ) and isothermal ( $T\sim 1\text{MK}$ ) solar corona. A back-of-the-envelope calculation gives a ratio between the solar surface pressure ( $\sim 100\text{ mb}$ ) and pressure infinitely away from the Sun to be  $P_{\text{inf}}/P_{\text{surf}} \sim 10^{-4}$ . But estimates from interstellar pressure give  $P_{\text{inf}}/P_{\text{surf}} \sim 10^{-14}!!$

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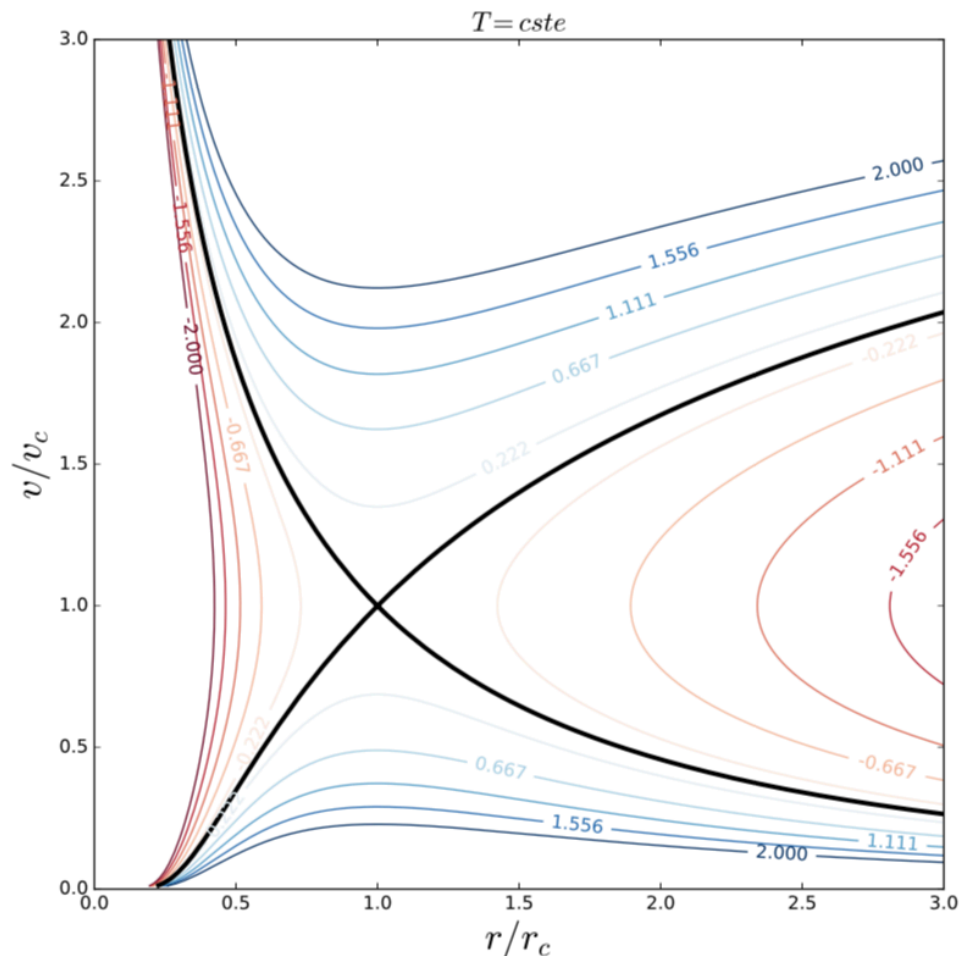




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In **1958**, E. Parker postulates the existence of a **stationary wind** in the corona of the Sun to solve this paradox (this was also hinted by cometary tails observed by Biermann 1951)





# The winds from cool stars

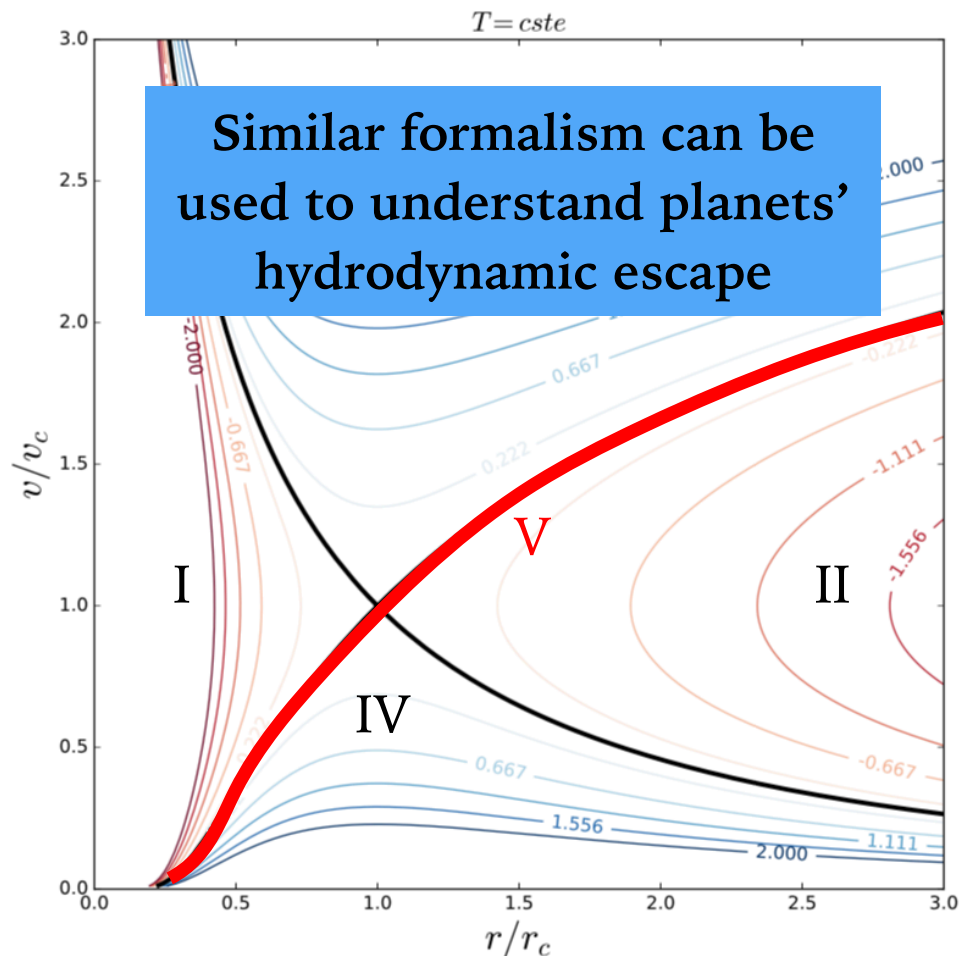
Let us assume a static ( $u=0$ ) and isothermal ( $T \sim 1\text{MK}$ ) solar corona. A back-of-the-envelope calculation gives a ratio between the solar surface pressure ( $\sim 100\text{ mb}$ ) and pressure infinitely away from the Sun to be  $P_{\text{inf}}/P_{\text{surf}} \sim 10^{-4}$ . But estimates from interstellar pressure give  $P_{\text{inf}}/P_{\text{surf}} \sim 10^{-14}!!$

In 1958, E. Parker postulates the existence of a **stationary wind** in the corona of the Sun to solve this paradox (this was also hinted by cometary tails observed by Biermann 1951)

**III:** Not observed

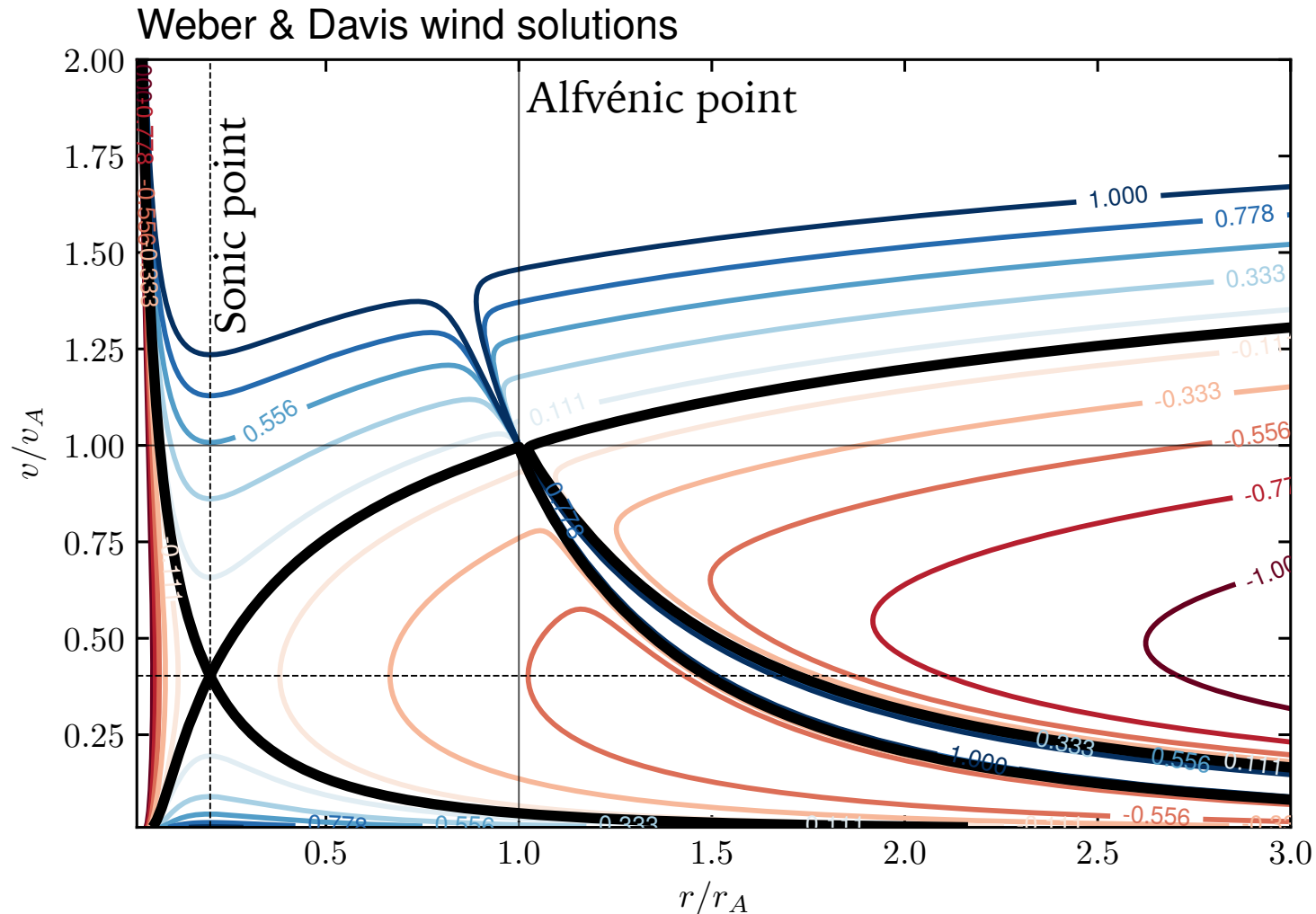
**IV:**  $P_{\text{inf}}$  too high

**V:** Parker's solution



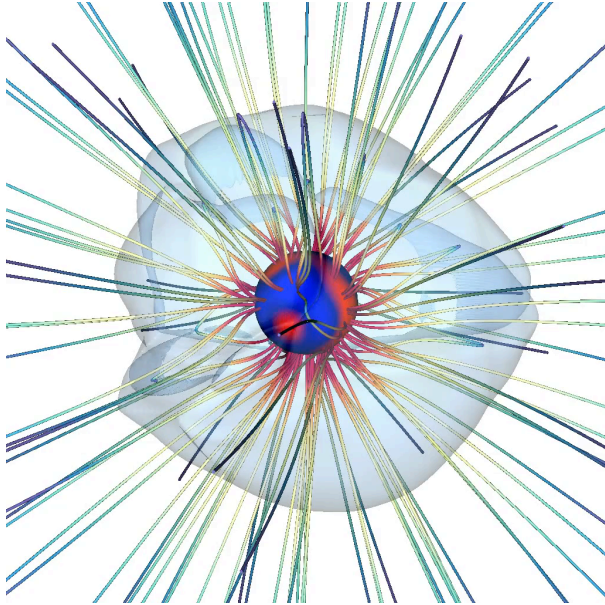
# Magnetized stellar winds

But the wind develops in a magnetized medium. Hence, Weber & Davis (1967) expanded Parker's theory in the MHD case



# Stellar magnetism is directly related to stellar evolution

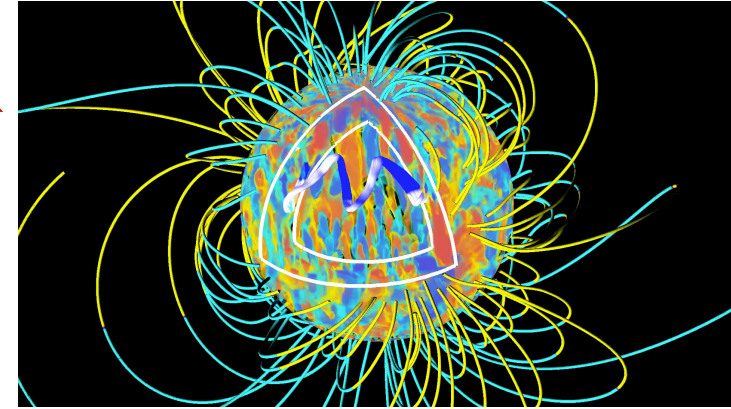
Rotation



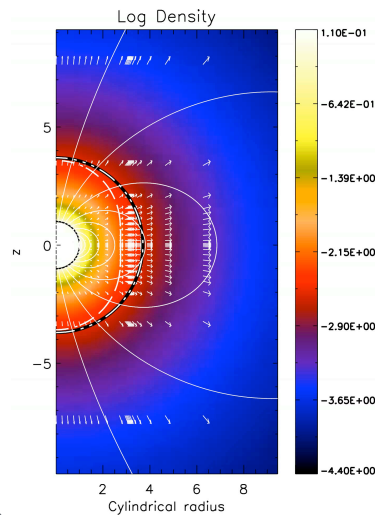
Dynamo



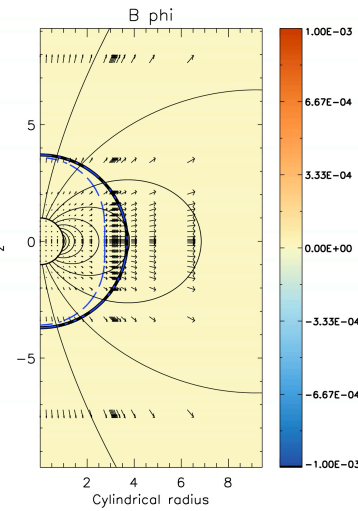
Magnetic field



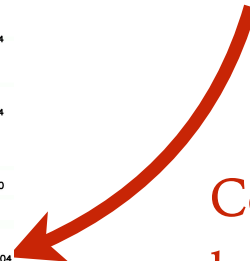
Torques



Wind

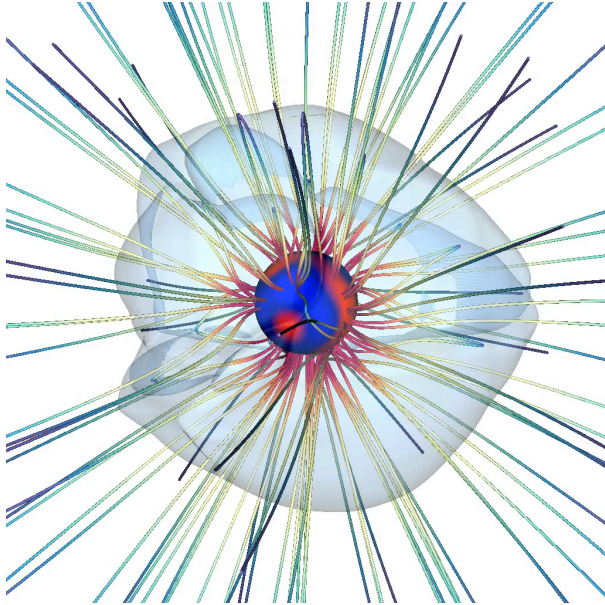


Coronal heating



# Stellar magnetism is directly related to stellar evolution

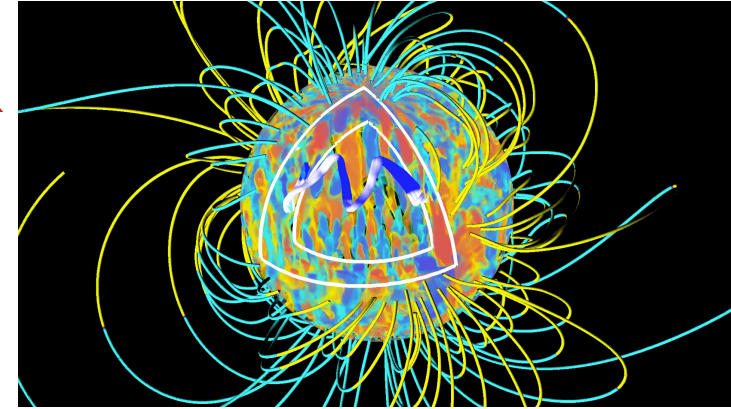
Rotation



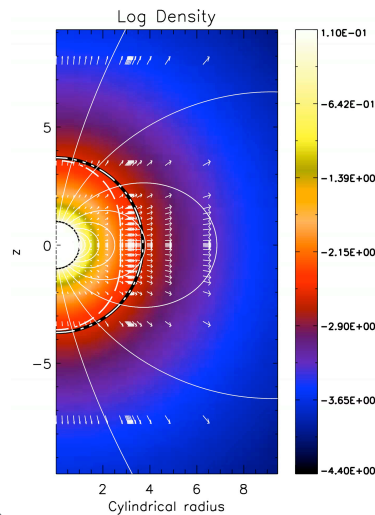
Dynamo



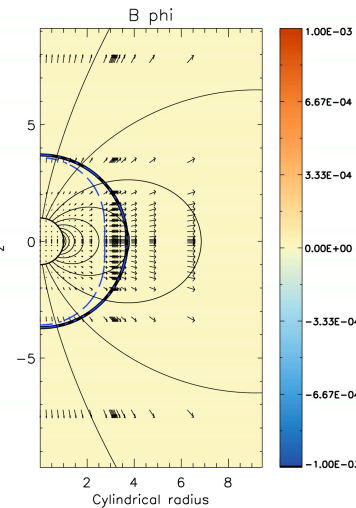
Magnetic field



Torques



Wind



Coronal heating



# Outline

## I. Setting the scene

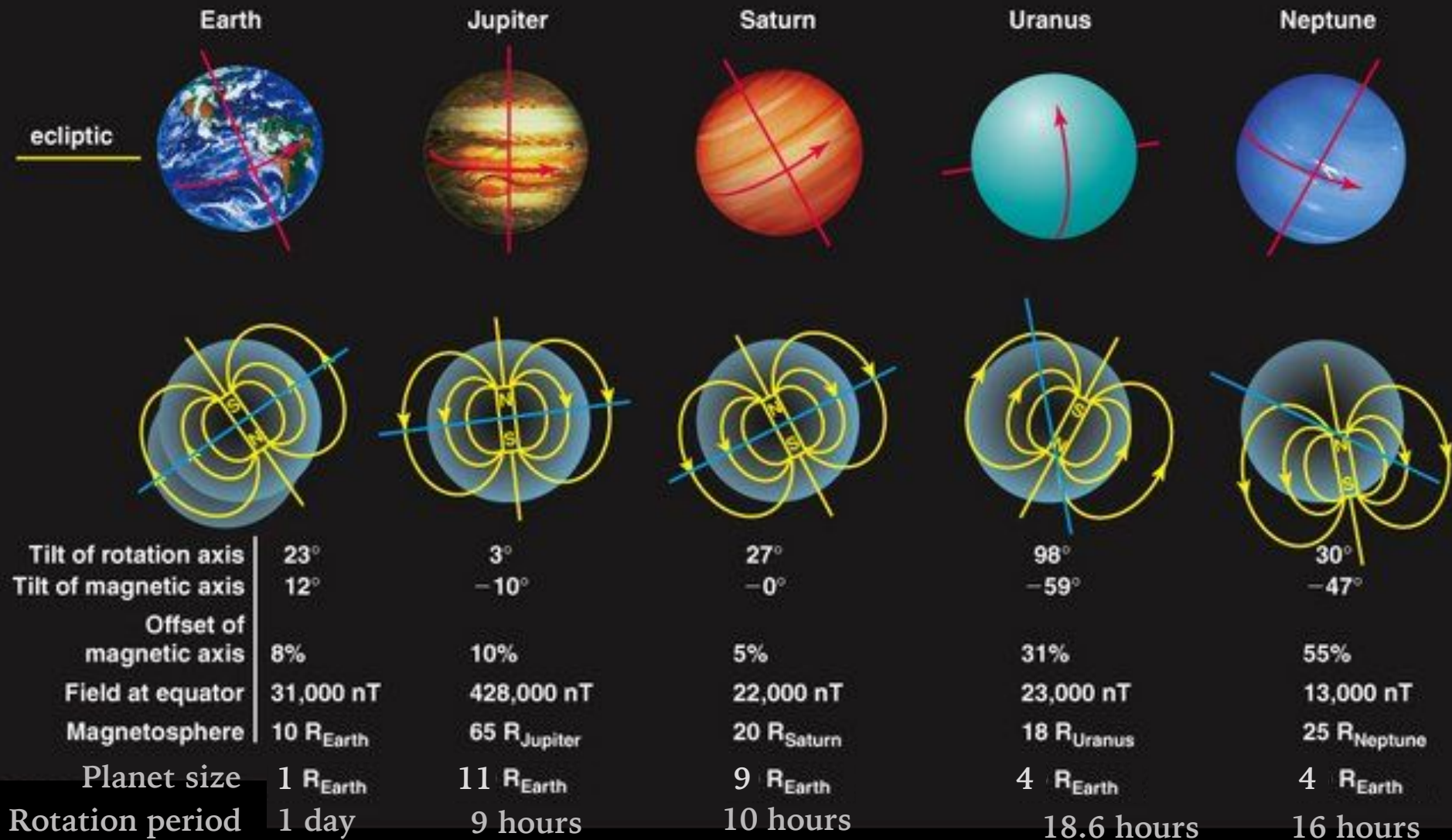
A. Stellar magnetism

B. Stellar winds

**C. Planetary magnetism**

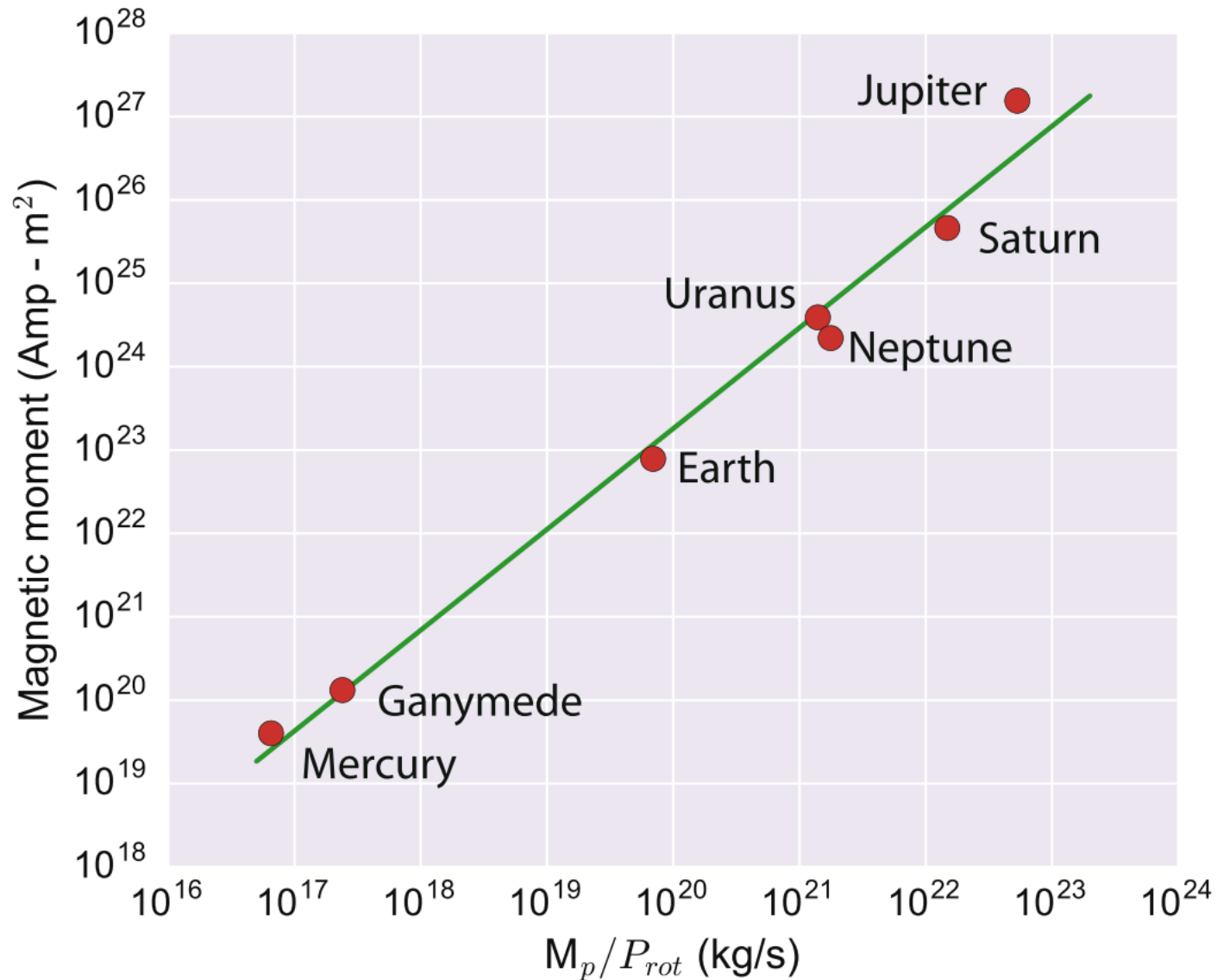
## II. Star-planet magnetic interactions

# What do we know about planet magnetism?

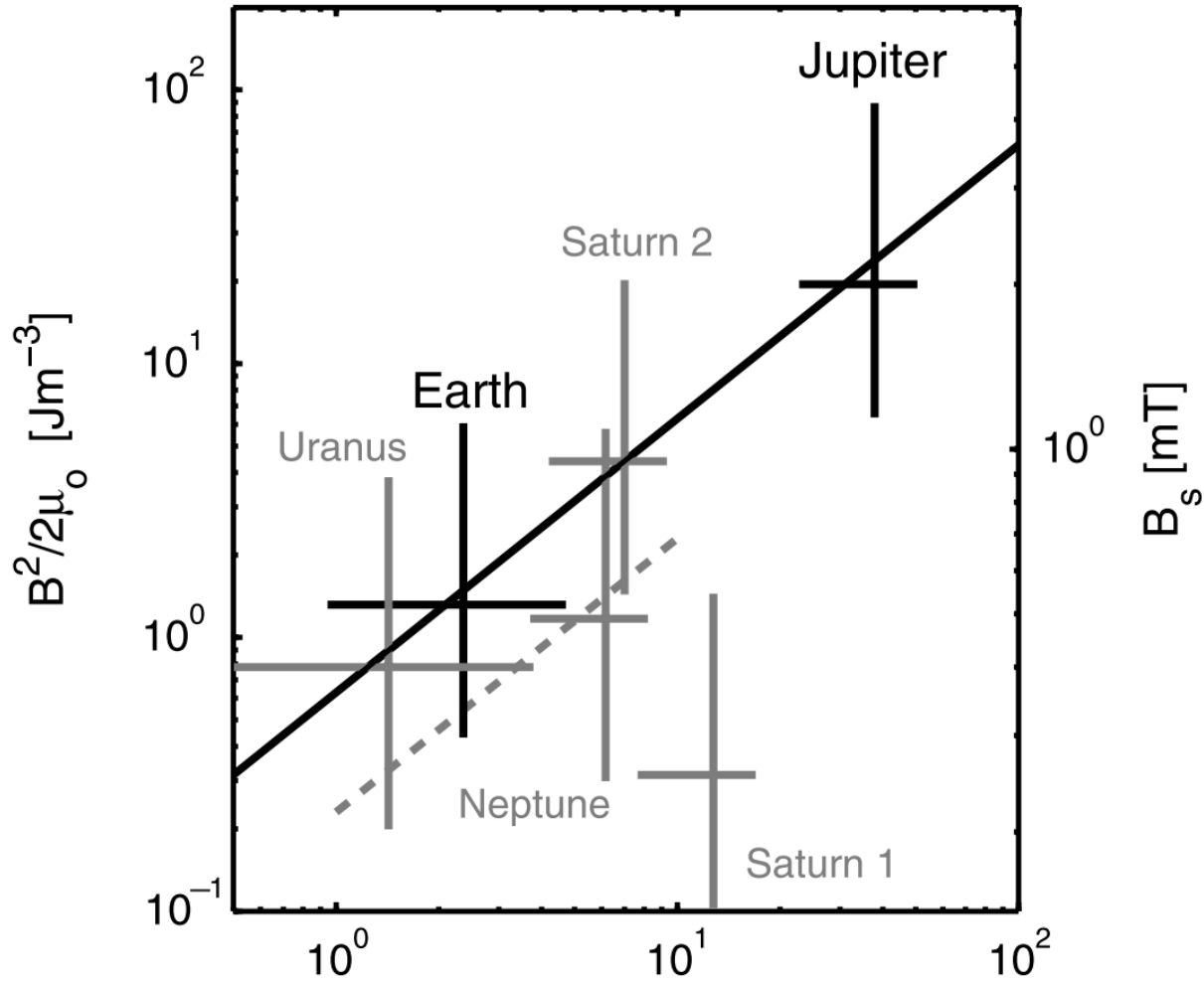




# What do we know about planet magnetism?



# What do we know about planet magnetism?



ohmic/total dissipation ———  $f_{\text{ohm}} \rho^{1/3} (F q_0)^{2/3}$  [Jm<sup>-3</sup>]

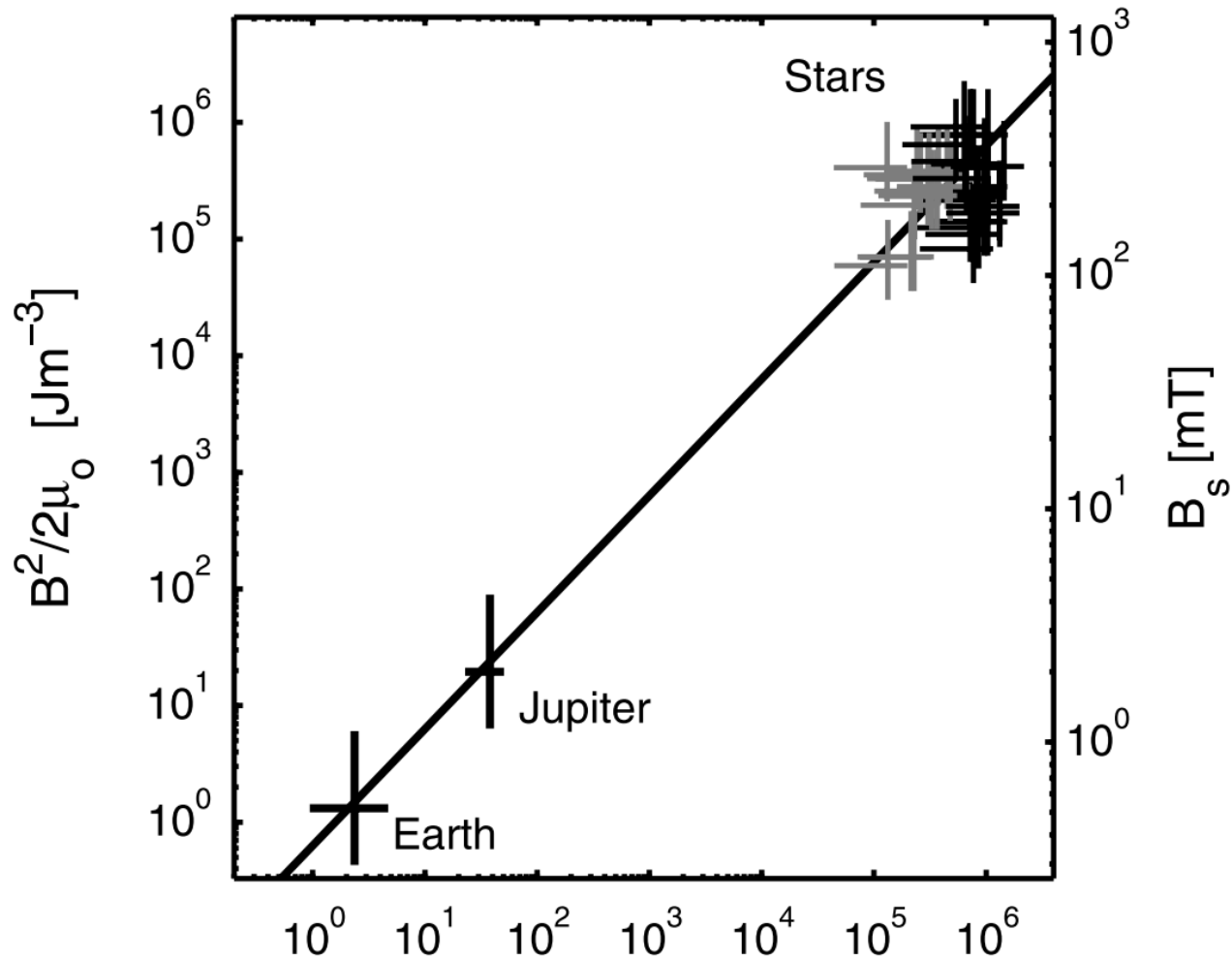
average density ———  $\rho$

geometric factor  $\sim 1$  ———  $f_{\text{ohm}}$

Bolometric flux ———  $(F q_0)^{2/3}$

[Christensen 2010]

# What do we know about planet magnetism?

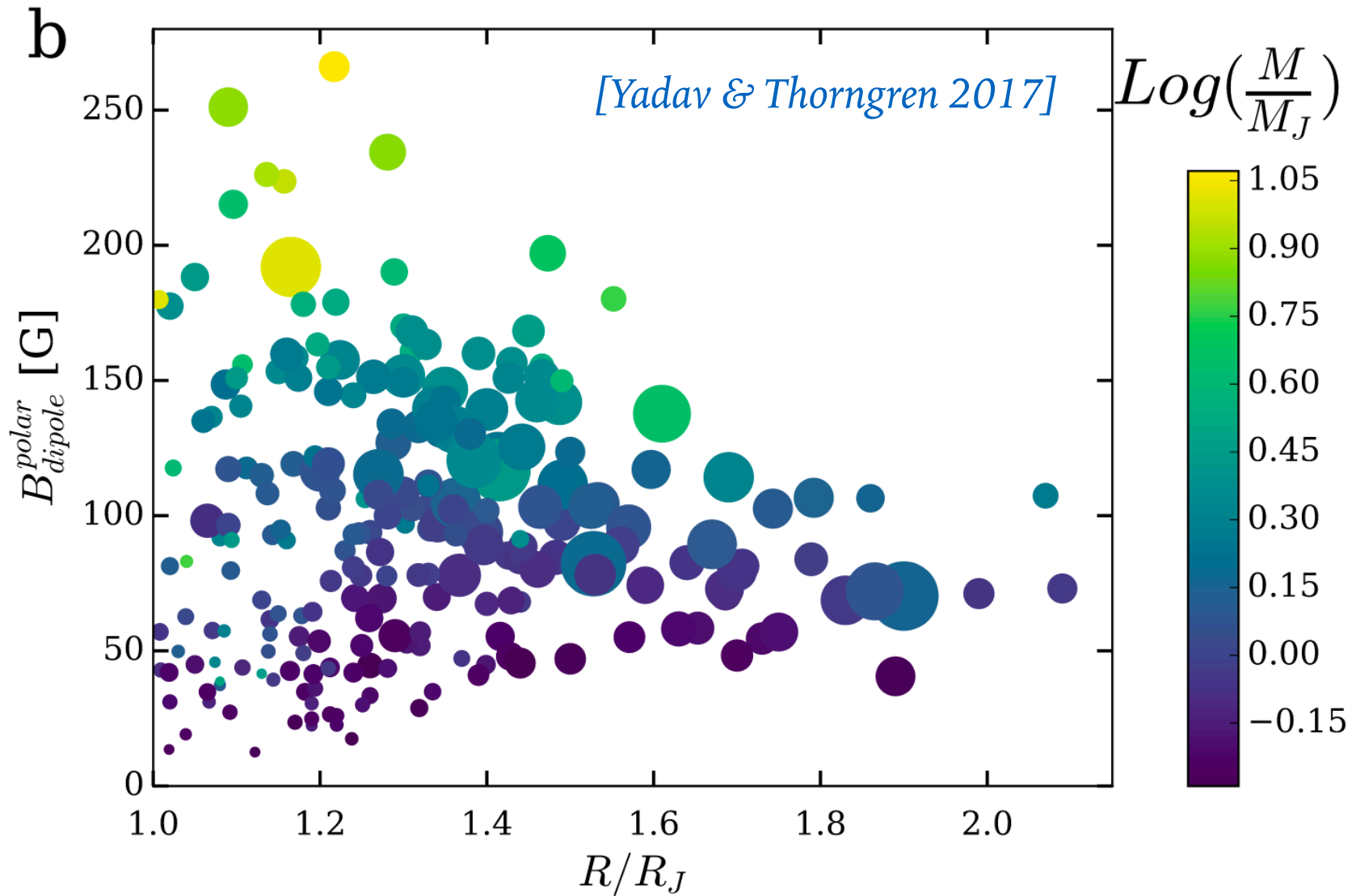


ohmic/total dissipation  $\sim f_{\text{ohm}} \rho^{1/3} (F q_o)^{2/3}$  [Jm<sup>-3</sup>]

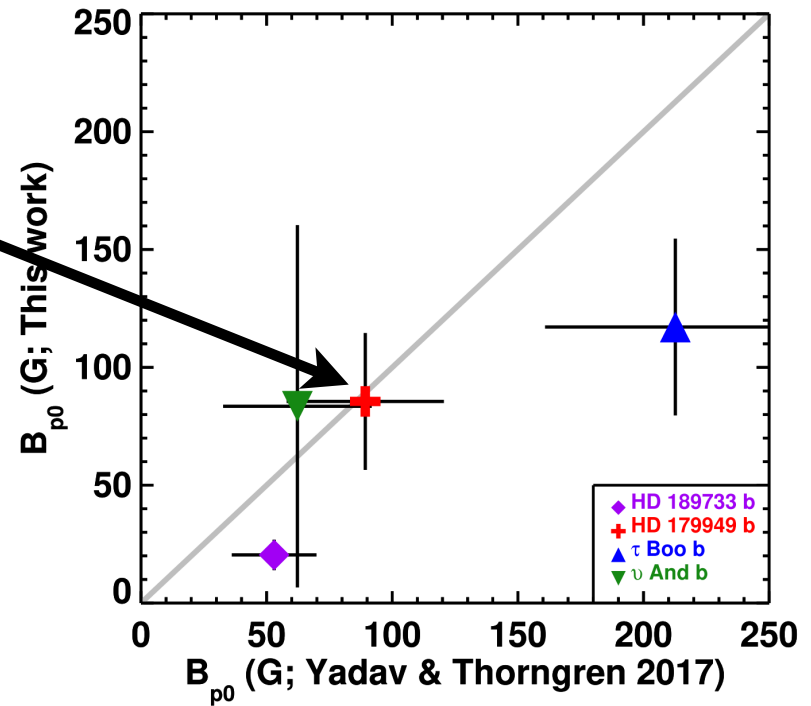
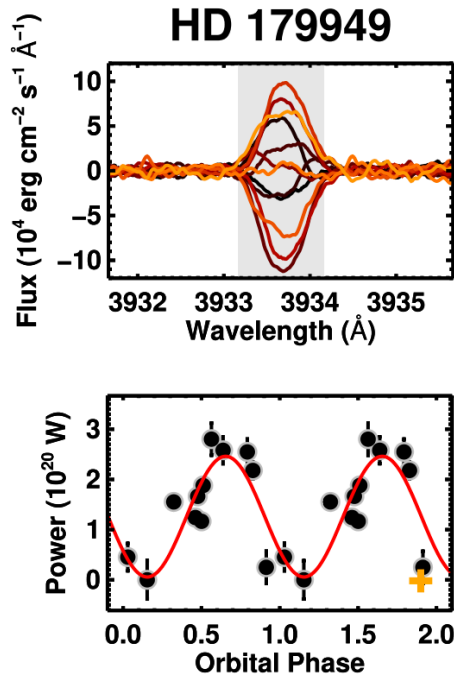
average density  $\rho$       geometric factor  $\sim 1$       Bolometric flux  $F q_o$

[Christensen 2010]

# Planetary dynamos powered by the stellar incident radiation

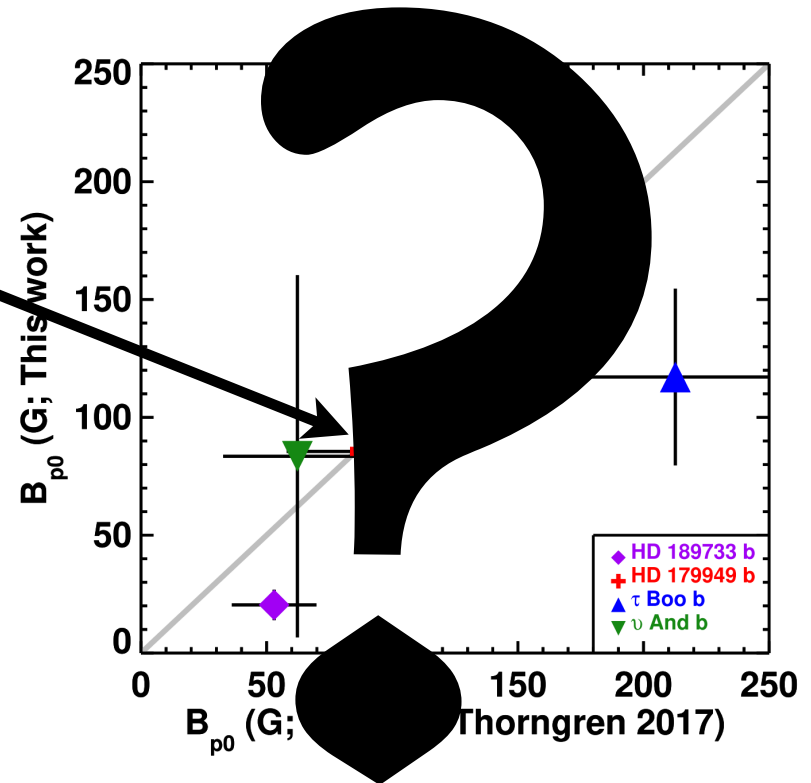
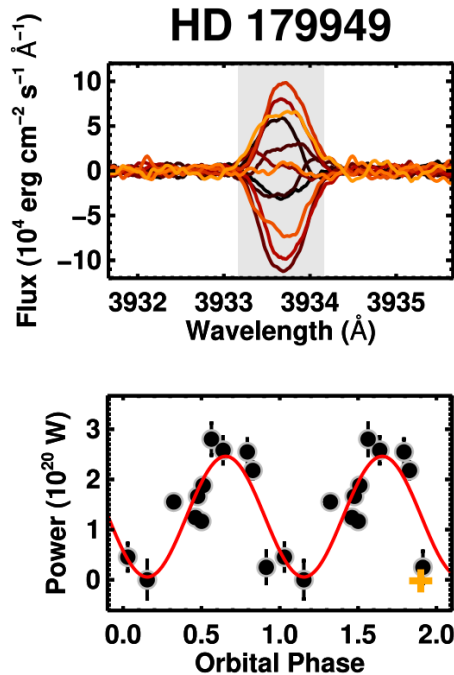


# What do we know about exoplanets' magnetism?



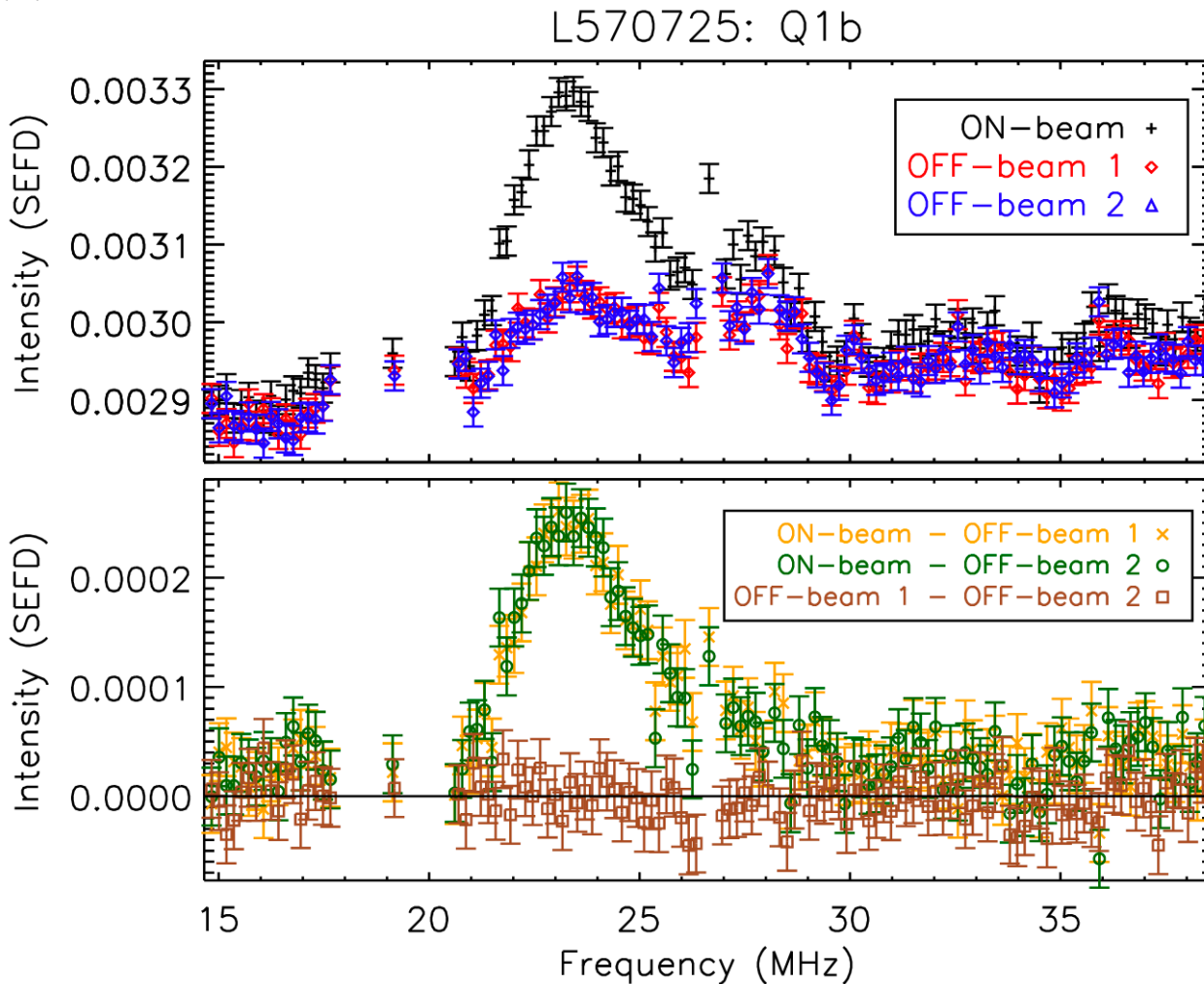
A signal in the Ca II H&K bands of 4 host stars was observed to be correlated with the orbital phase of the hot exoplanet (generally once the rotational modulation of the star is removed)

# What do we know about exoplanets' magnetism?



A signal in the Ca II H&K bands of 4 host stars was observed to be correlated with the orbital phase of the hot exoplanet (generally once the rotational modulation of the star is removed)

# The promise of radio detections: the case of $\tau$ Boo



[Turner + 21]

First tentative detection!

$B_p \sim 5-10$  G

[See also Callingham + 21, Vedantham + 21]

# Outline

## I. Setting the scene

A. Stellar magnetism

B. Stellar winds

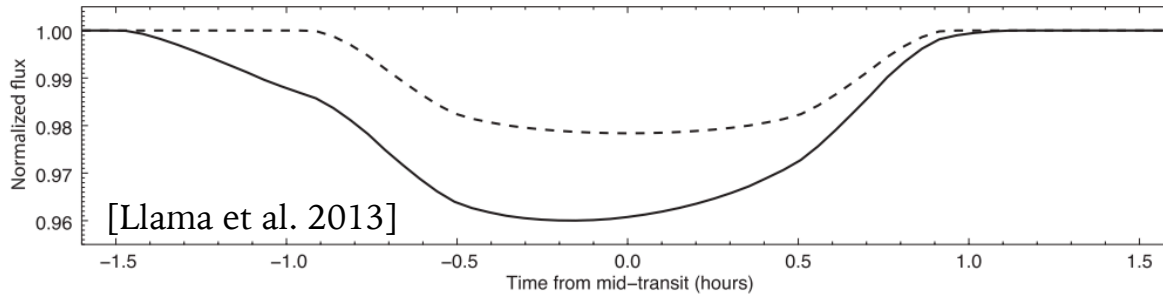
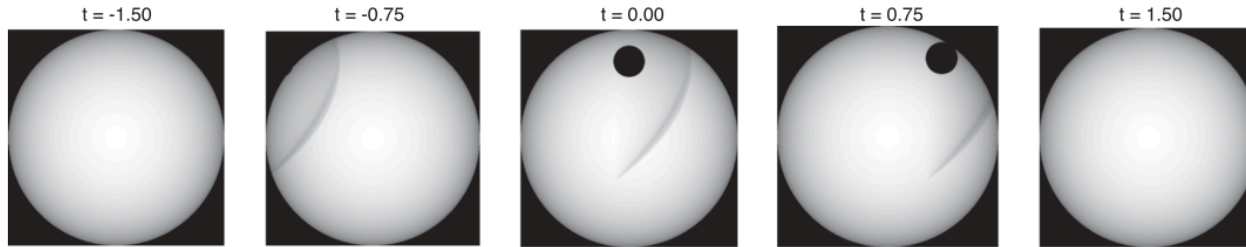
C. Planetary magnetism

## II. Star-planet magnetic interactions



# Star-planet (magnetic) interactions

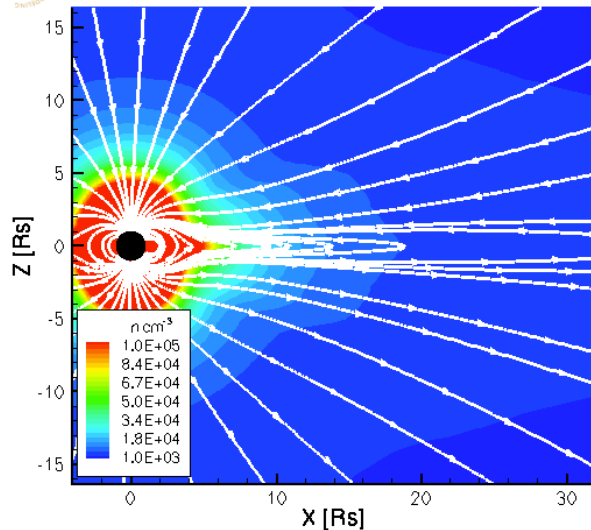
## Transients and short-timescale effects



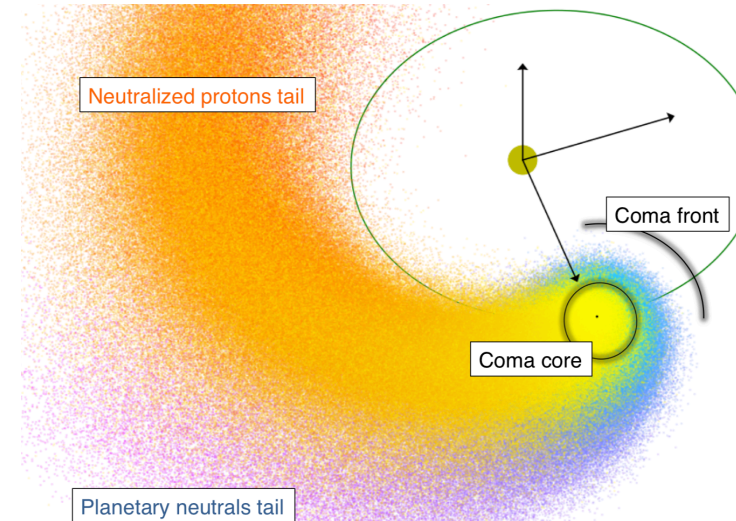
Pre-shock absorption

Atmospheric escape

University of Michigan  
Manchester et. al.  
2003



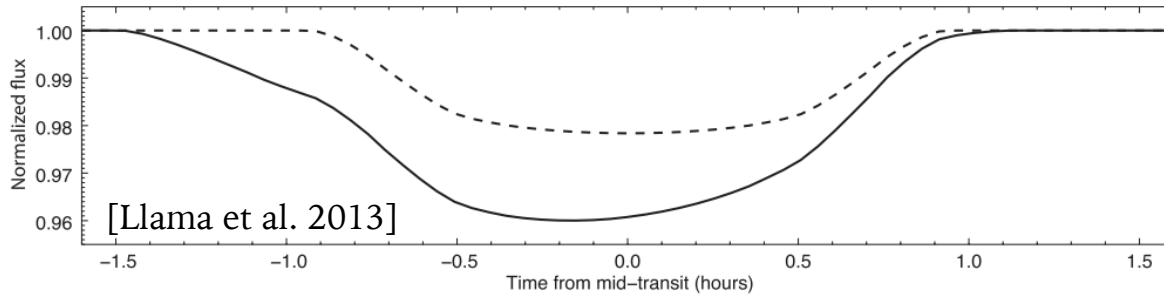
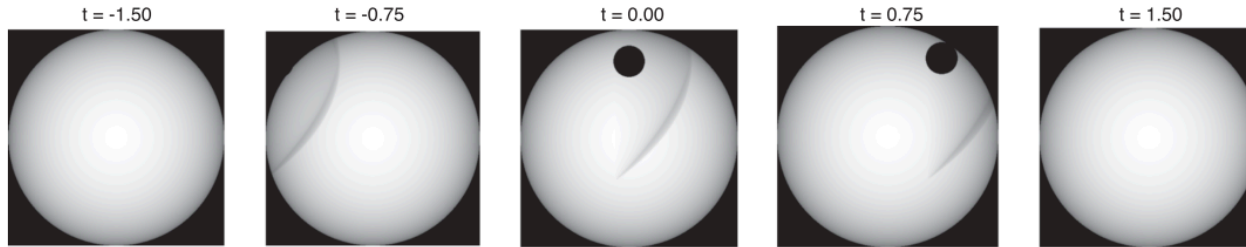
Transient interactions



GJ 436b, Bourrier + 2016

# Star-planet (magnetic) interactions

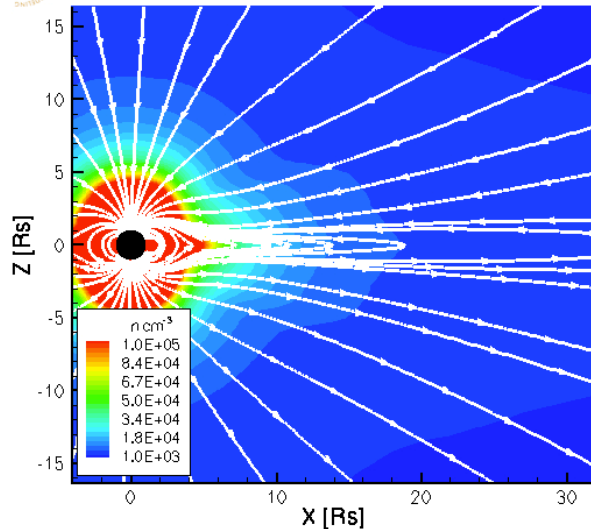
## Transients and short-timescale effects



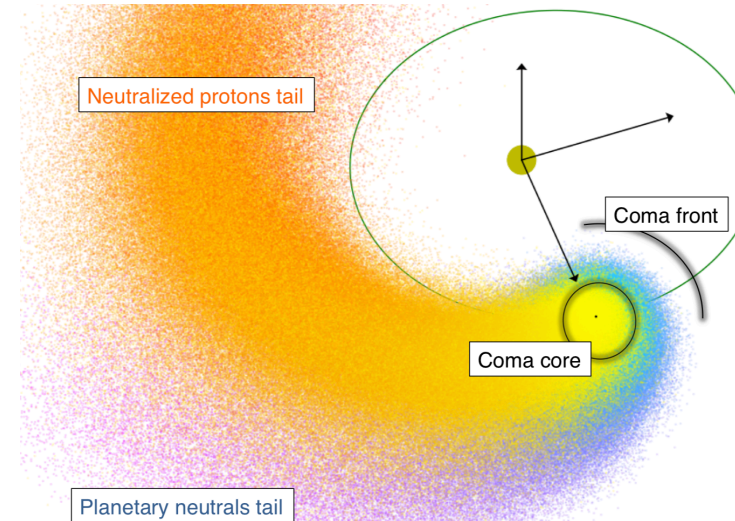
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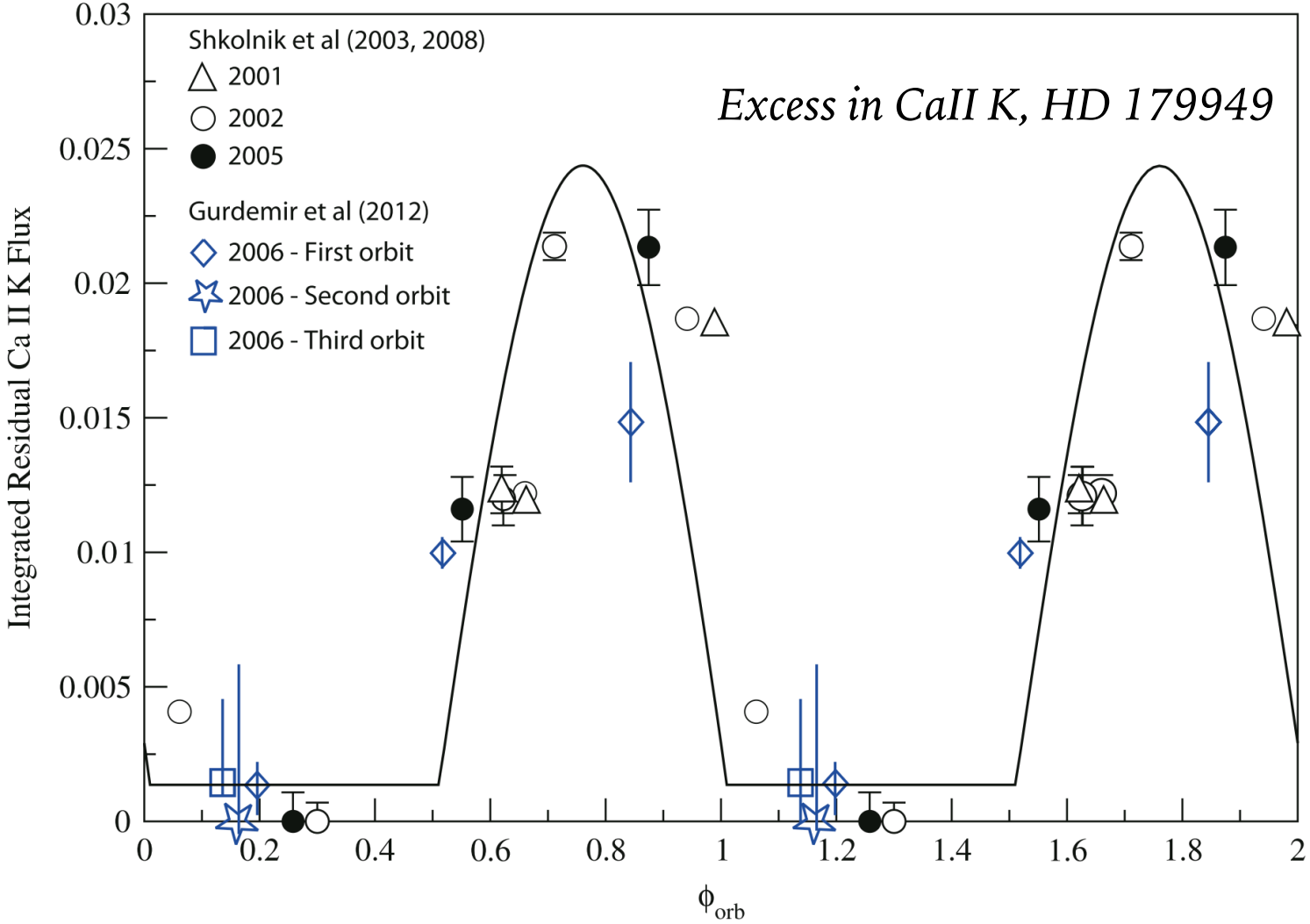
Transient interactions



GJ 436b, Bourrier + 2016

# Star-planet (magnetic) interactions

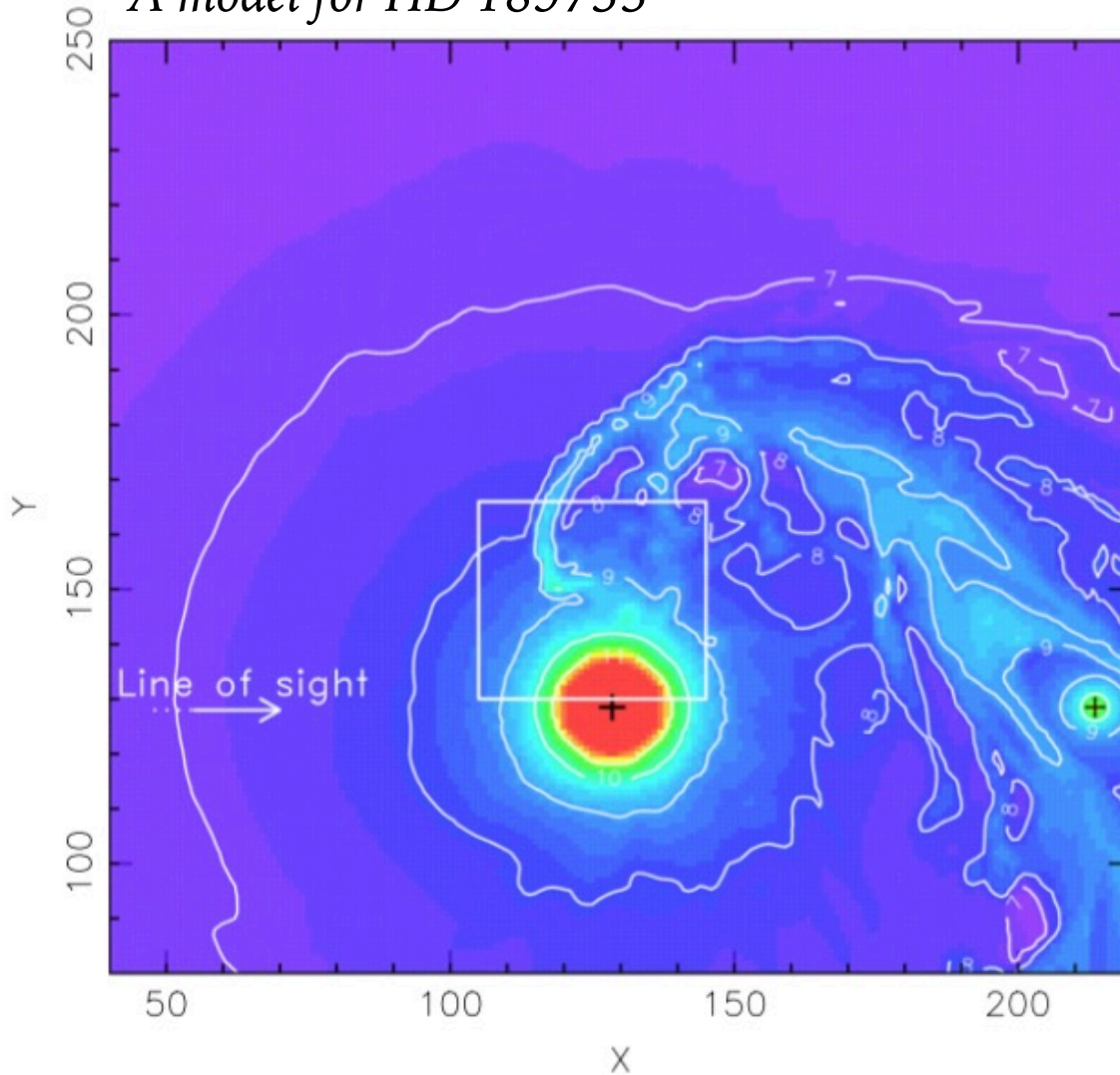
## Transients and short-timescale effects



# Star-planet (magnetic) interactions

## Transients and short-timescale effects

*A model for HD 189733*



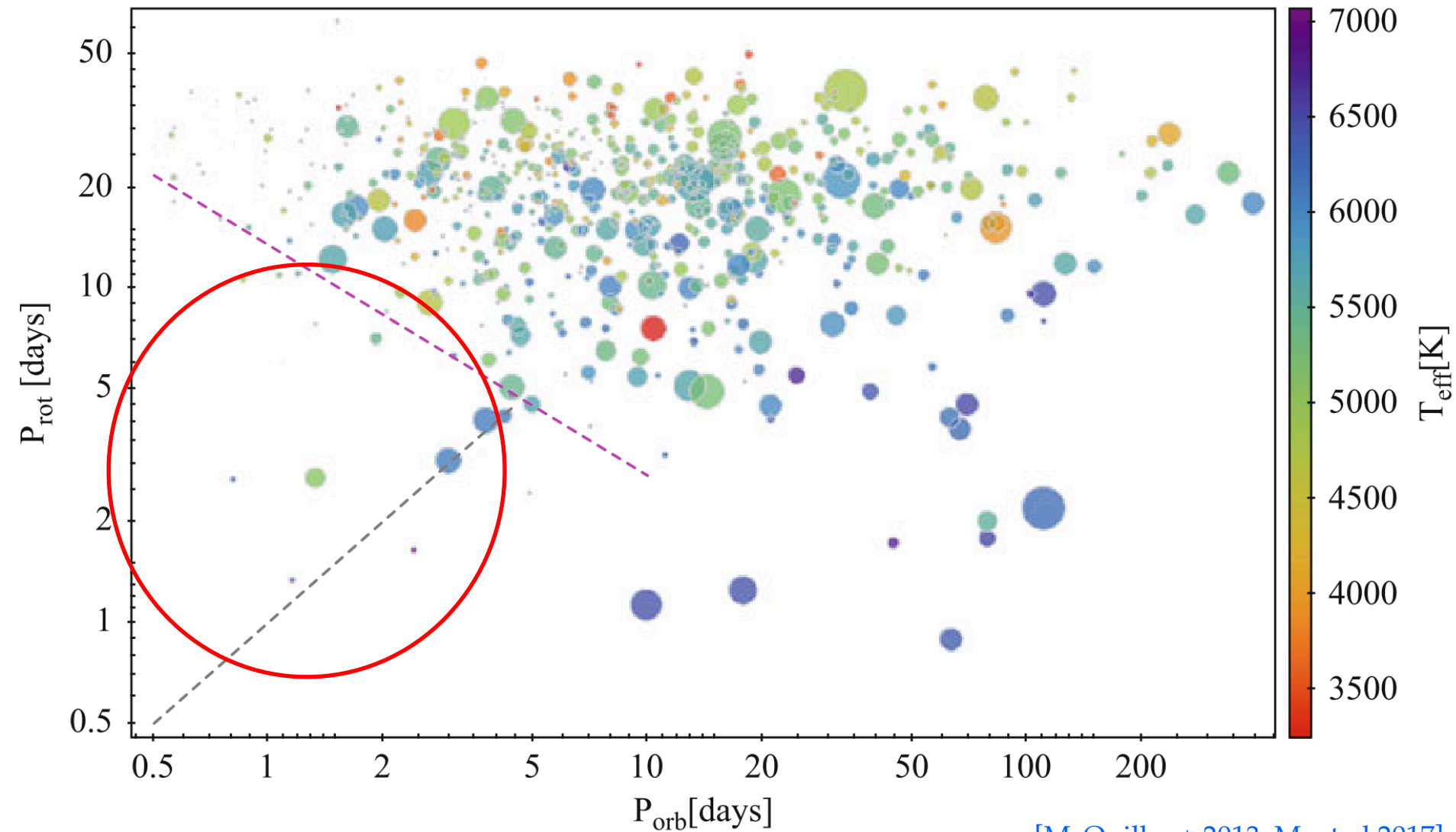
Back-reaction on  
the central star,  
potentially  
observable?

[Matsakos+ 15; Pillitteri+ 15]

# Star-planet (magnetic) interactions

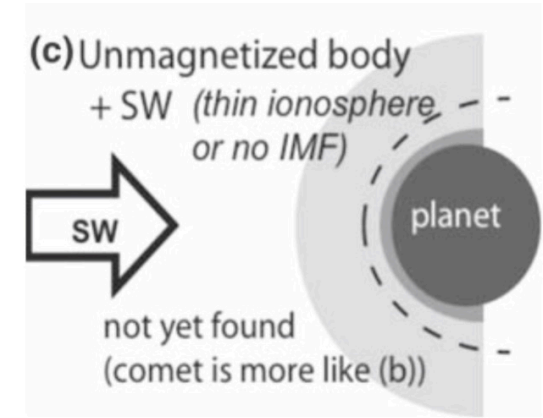
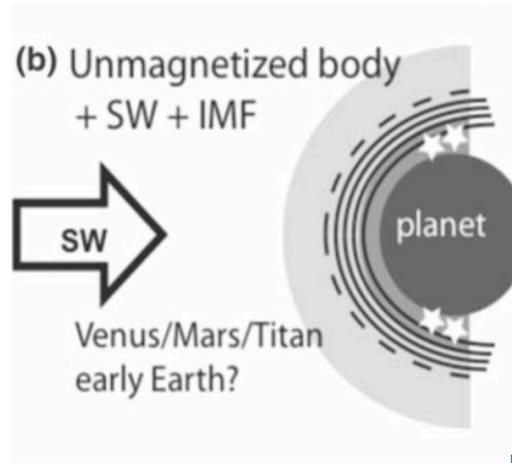
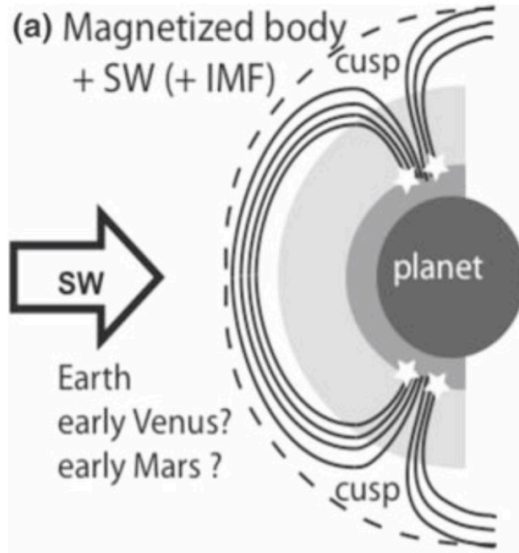
## Secular effects

737 Kepler Objects of Interest (KOI)

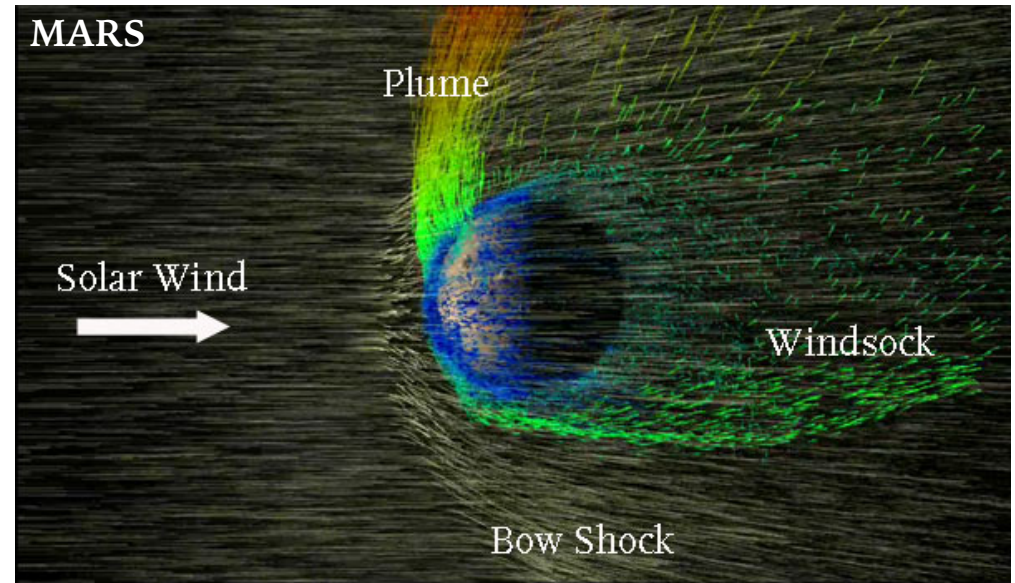
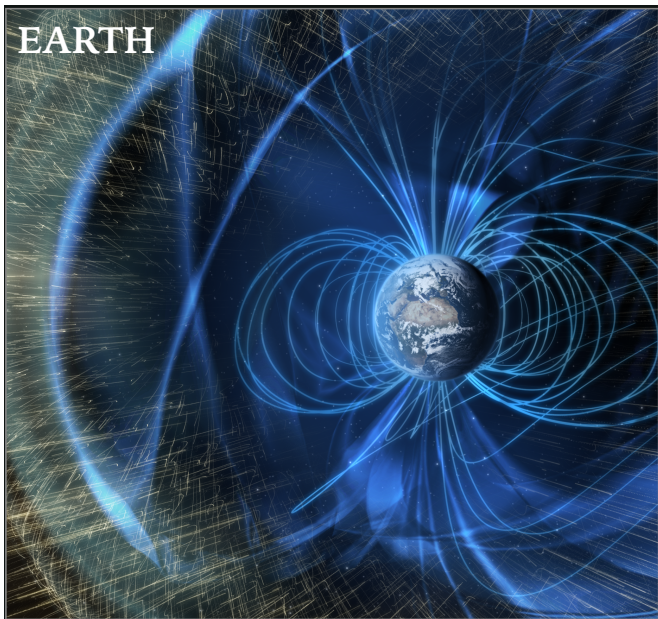


# Star-planet (magnetic) interactions

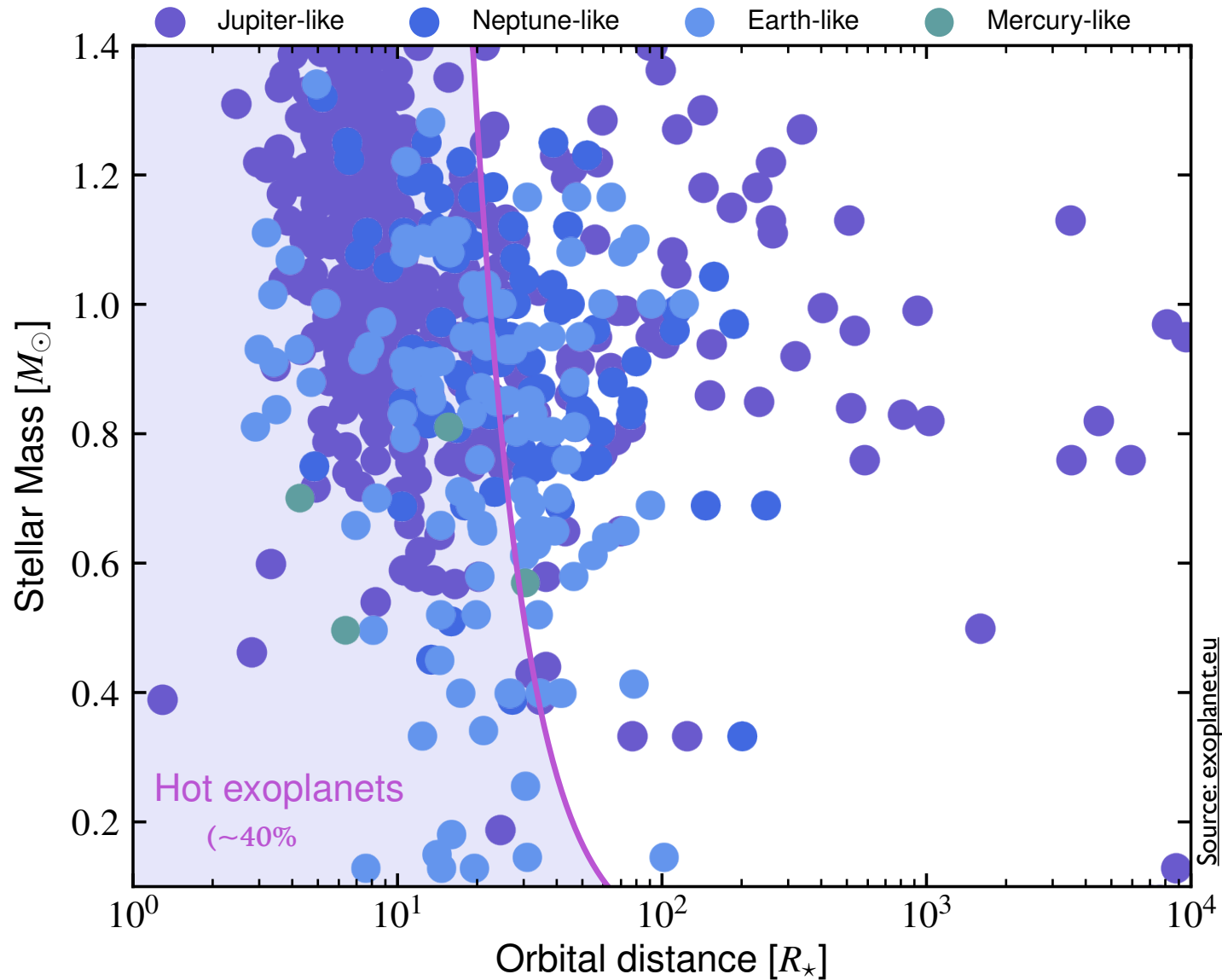
## towards habitability?



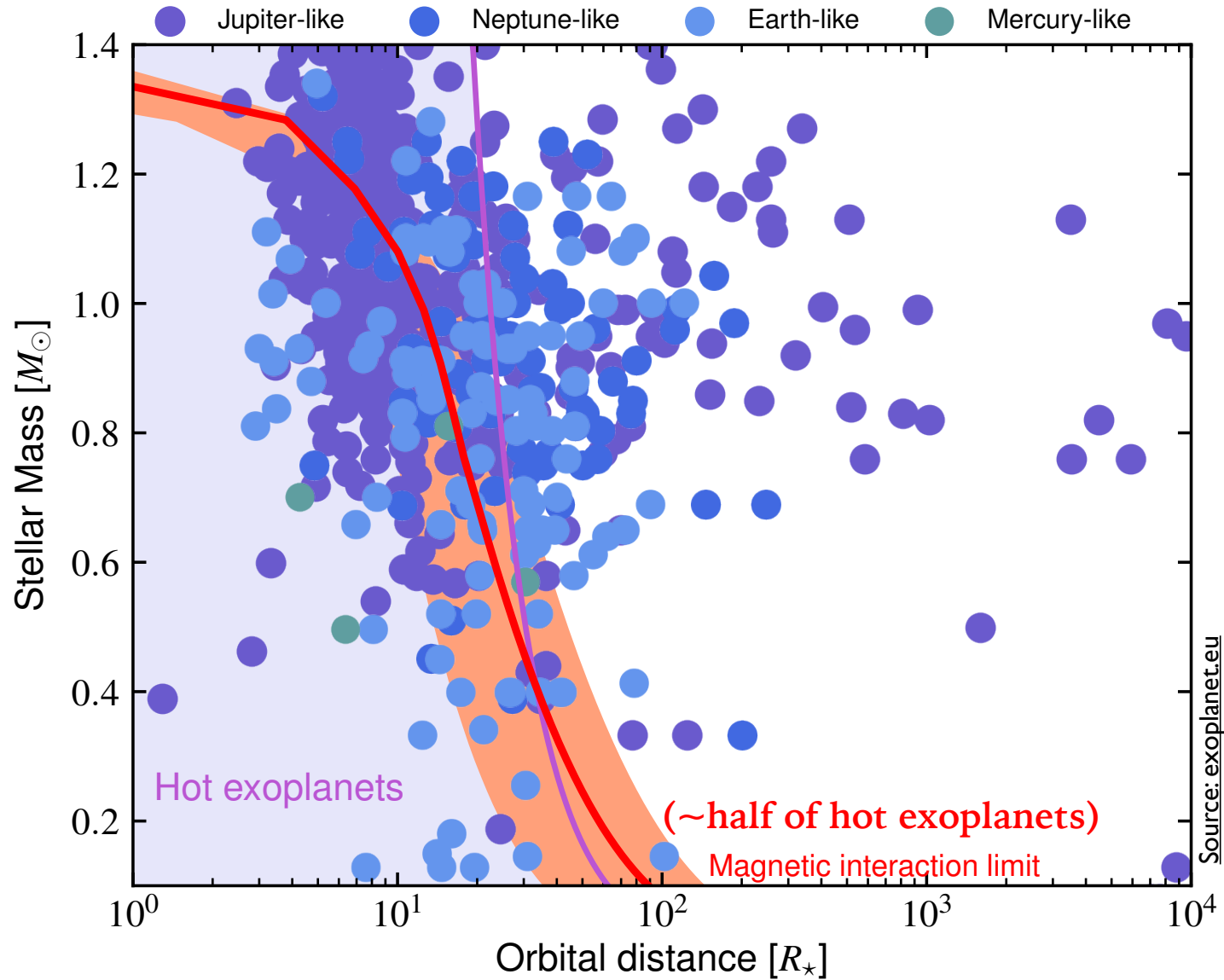
[Yamauchi and Wahlund 2007]



# Magnetic connection in hot exosystems



# Magnetic connection in hot exoplanets

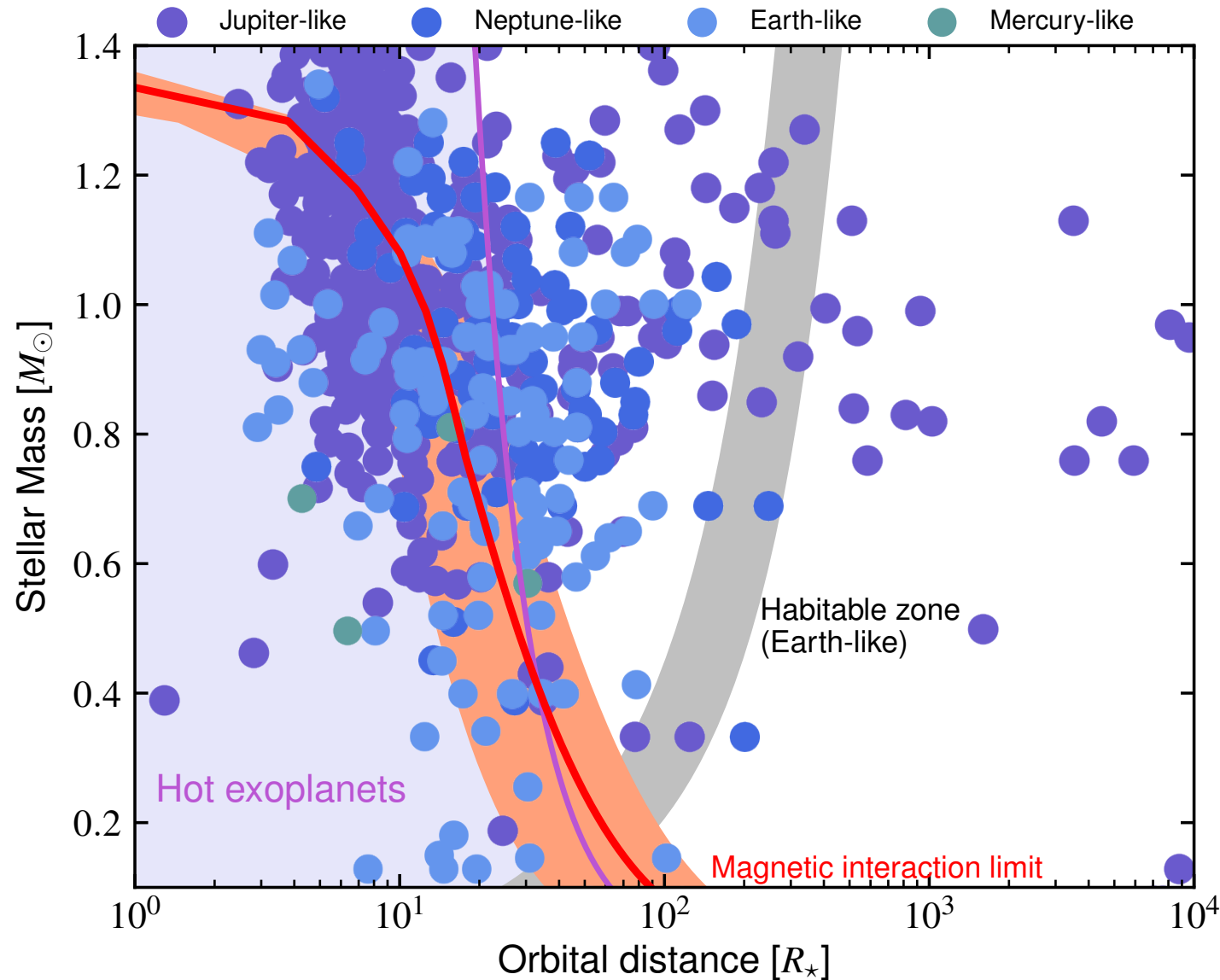


starAML 1D code  
Réville+ 2015





# Magnetic connection in hot exoplanets

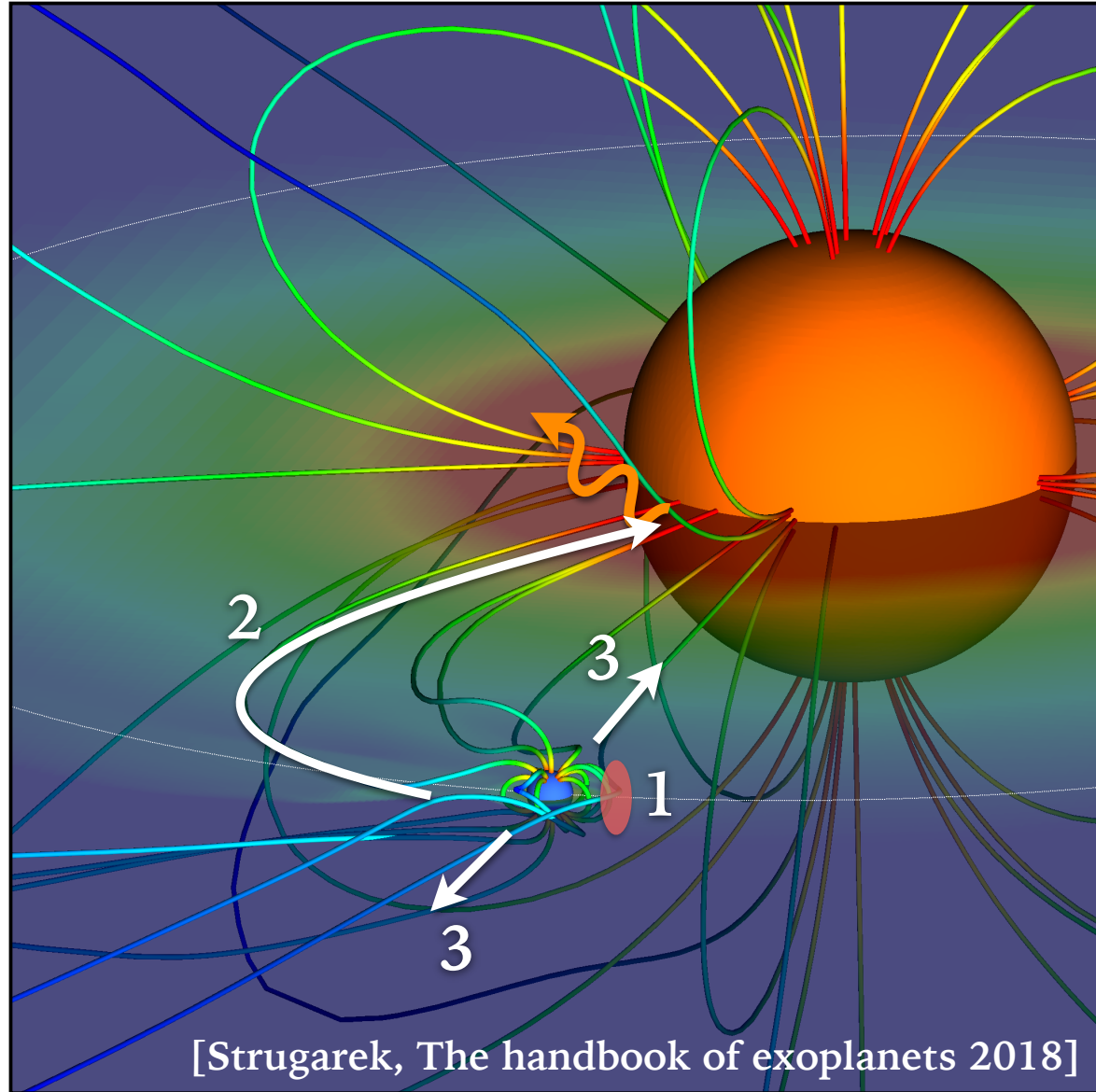
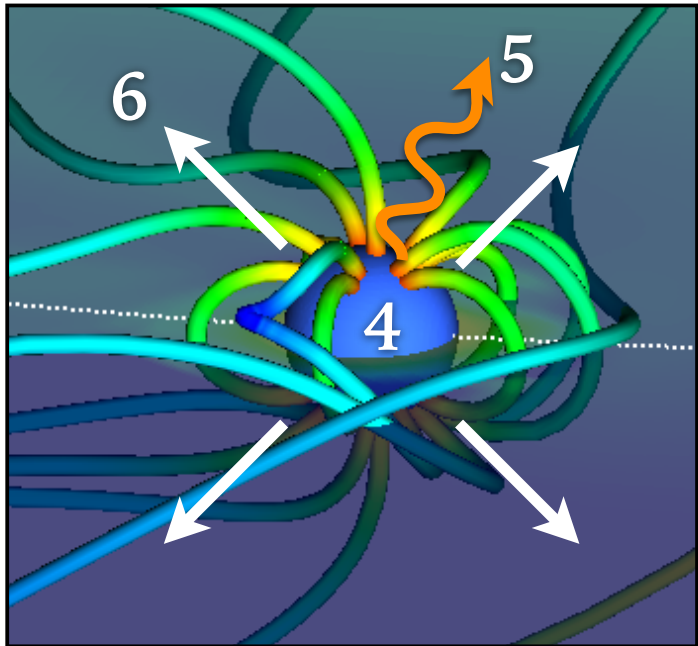


starAML 1D code  
Réville+ 2015



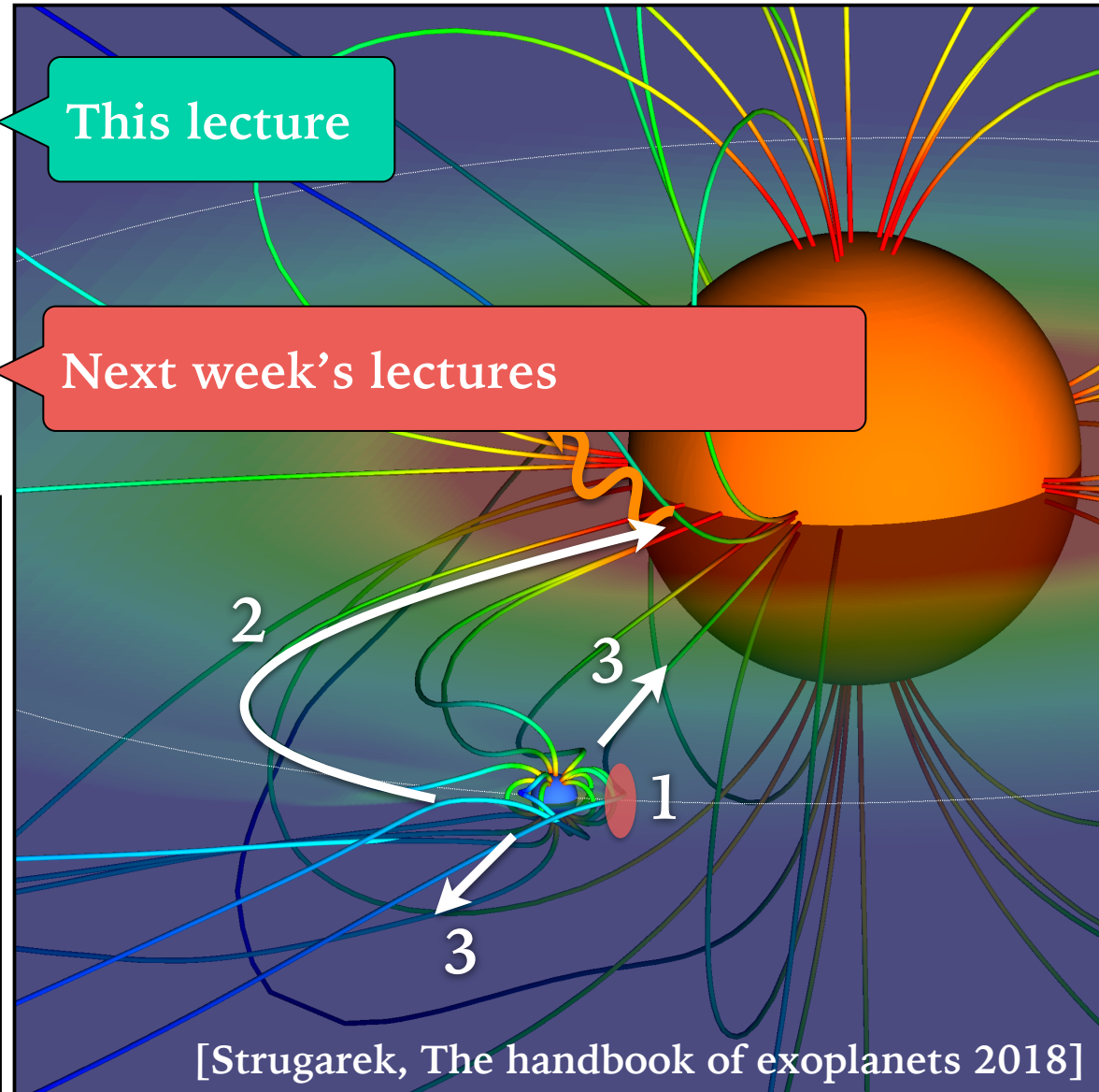
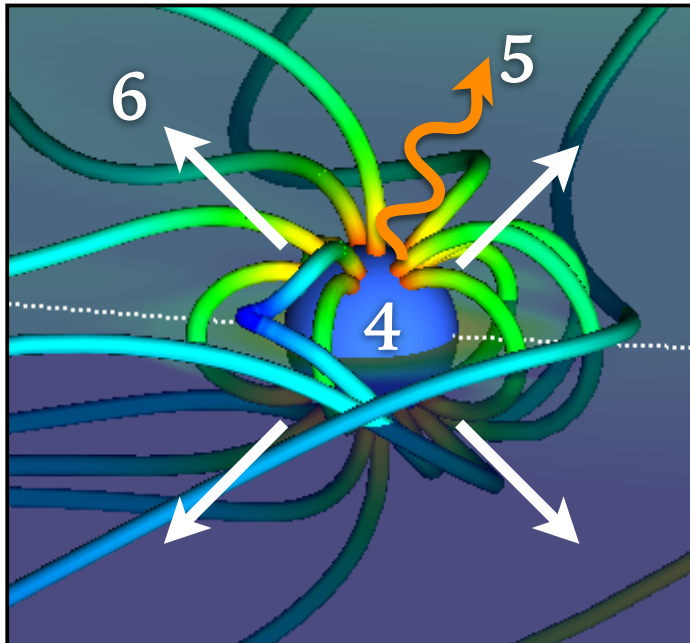
# Major effects of star-planet magnetic interactions

1. Shocks & geometry
2. Energy channeling
3. Planet migration
4. Planet heating
5. Planet emissions
6. Atmospheric escape



# Major effects of star-planet magnetic interactions

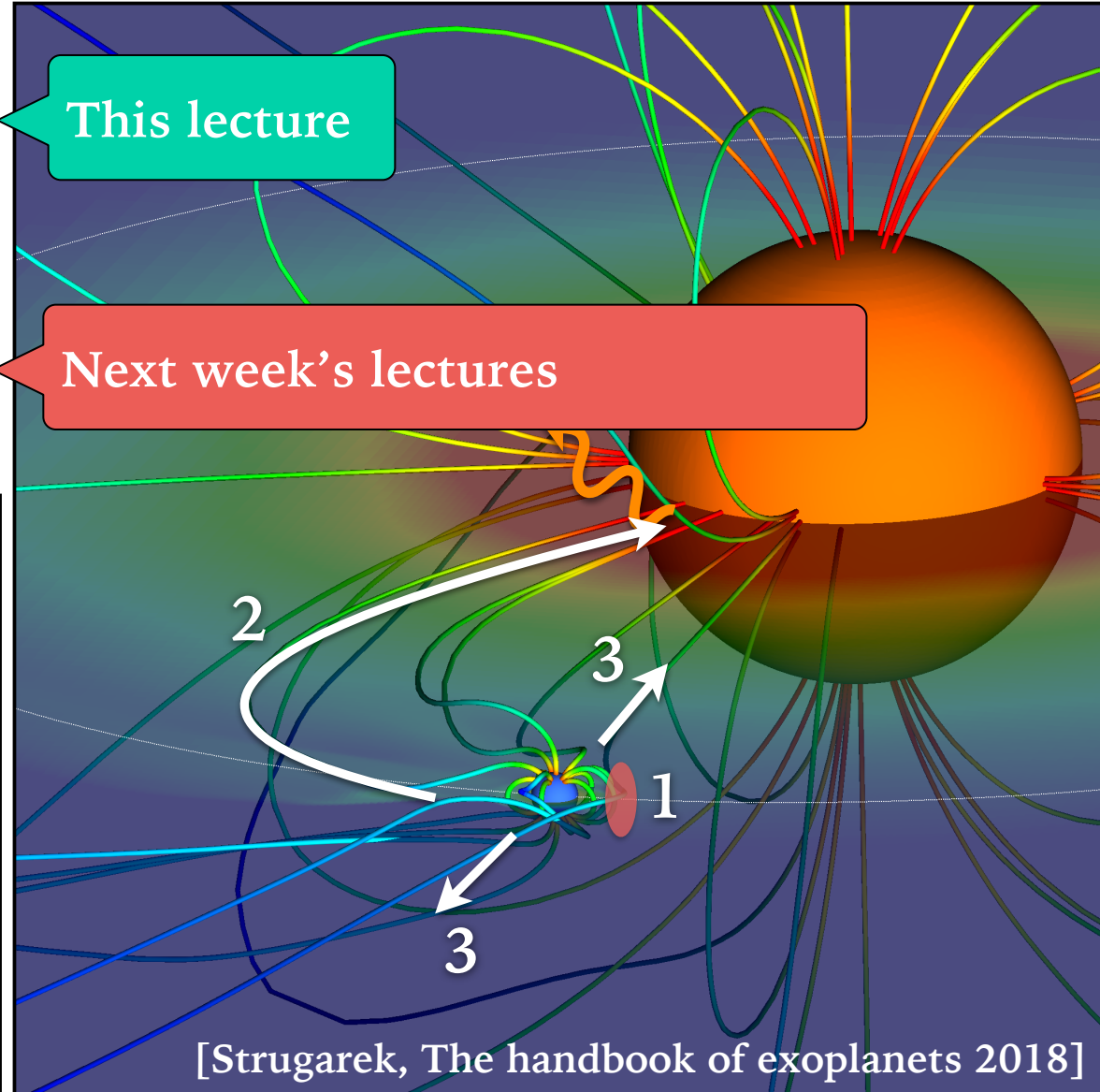
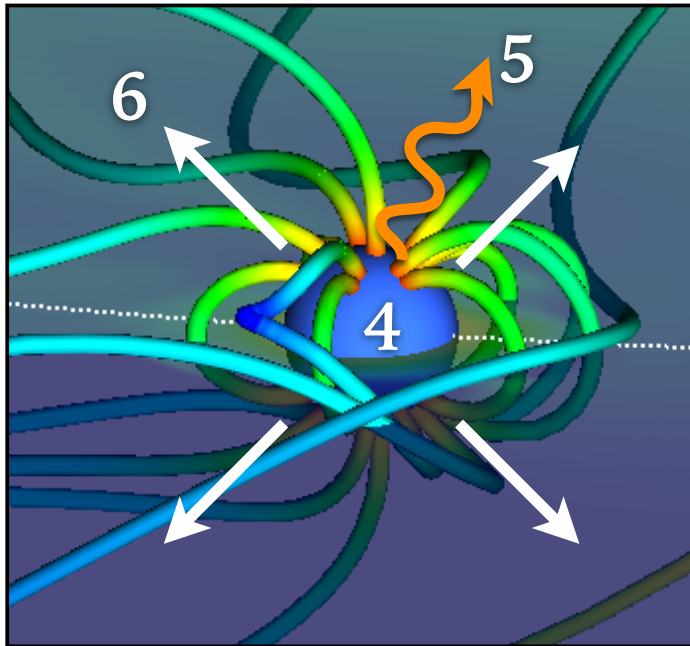
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[Strugarek, The handbook of exoplanets 2018]

# Major effects of star-planet magnetic interactions

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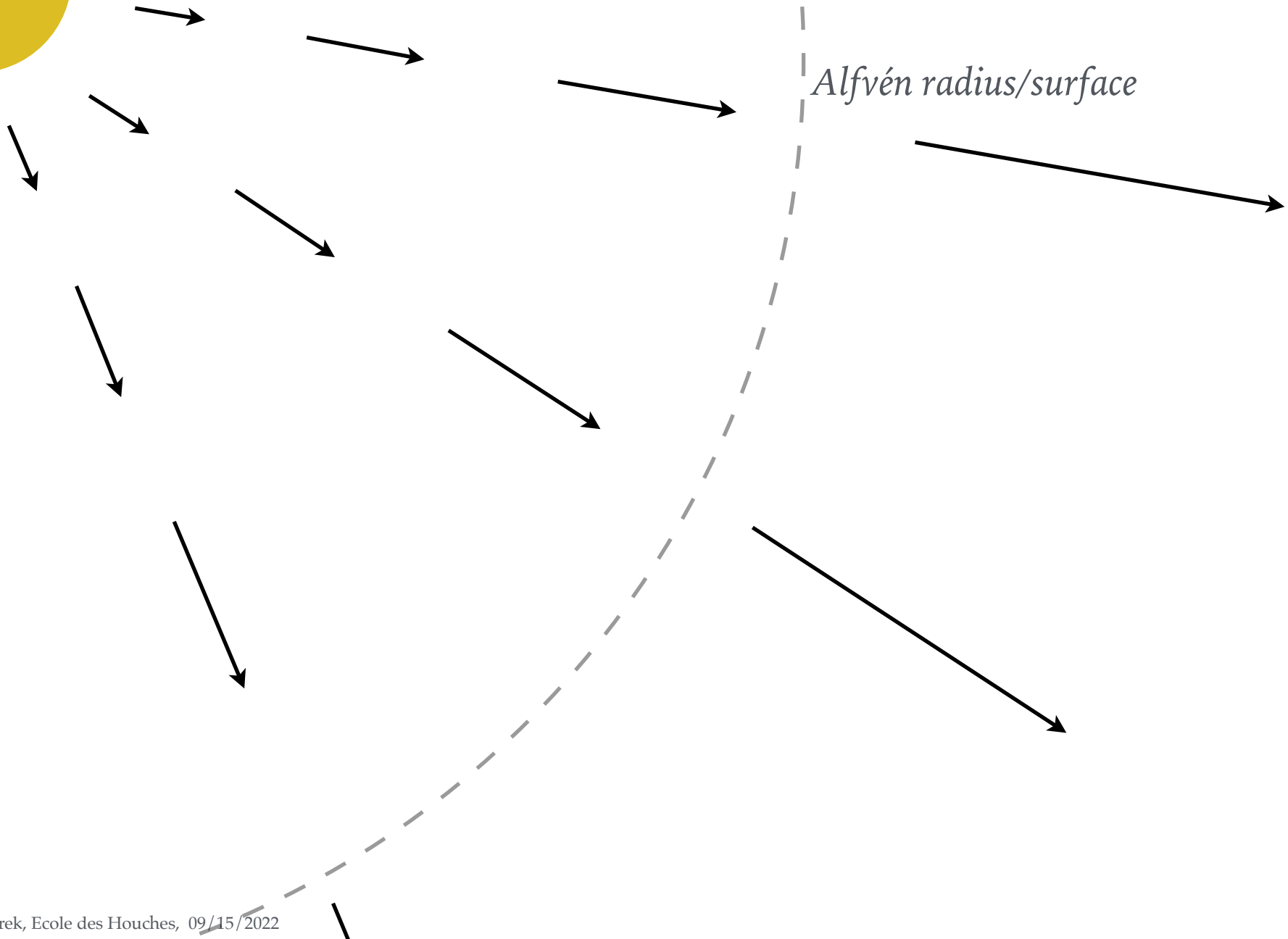
[Strugarek, The handbook of exoplanets 2018]



Star

# Two magnetic interaction regimes

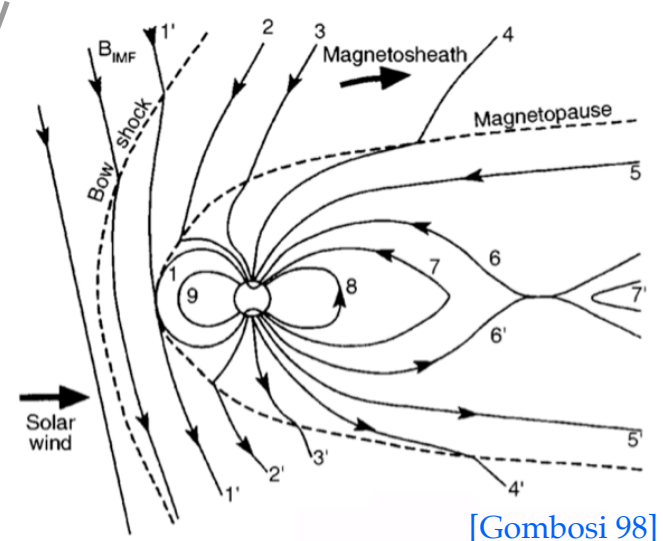
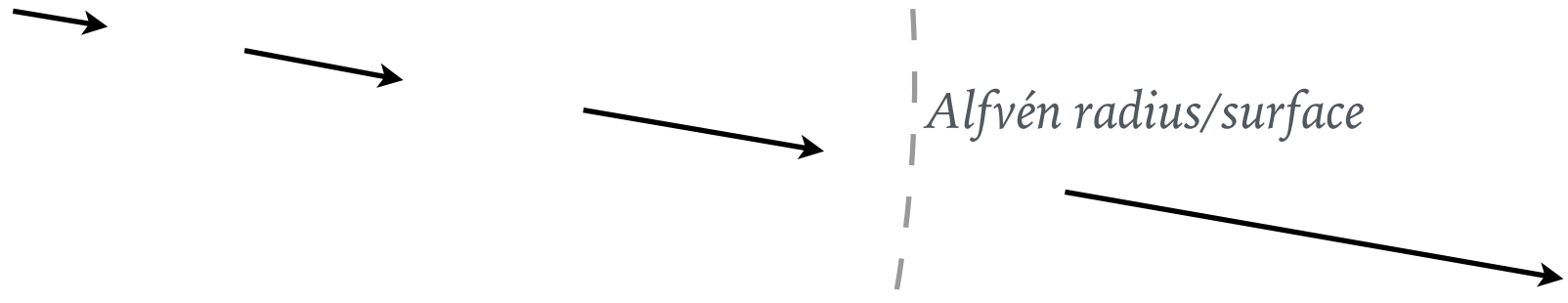
*Alfvén radius/surface*





Star

# Two magnetic interaction regimes

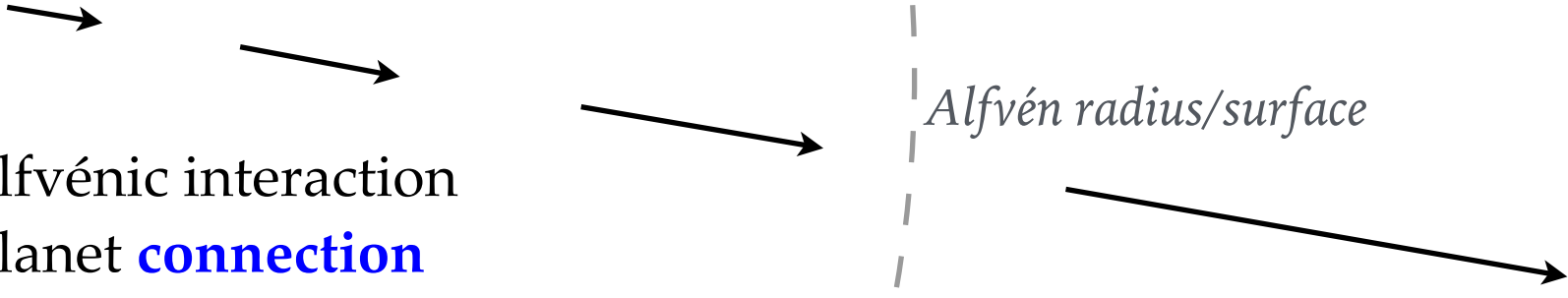


[Gombosi 98]

**Super-alfvénic interaction**  
**Shocks**

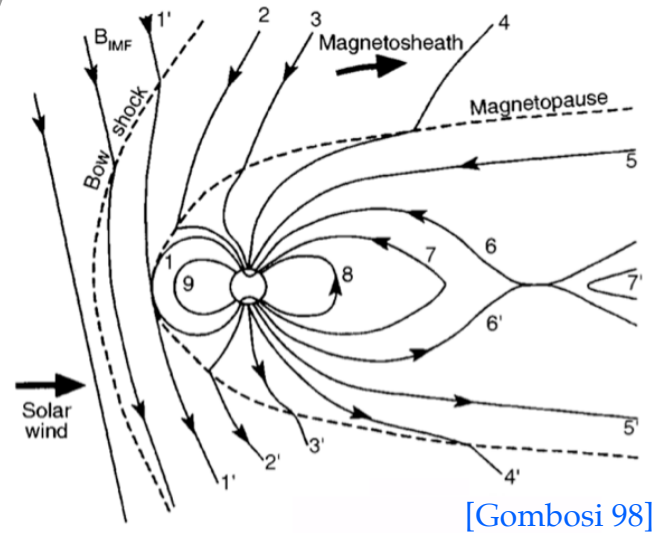
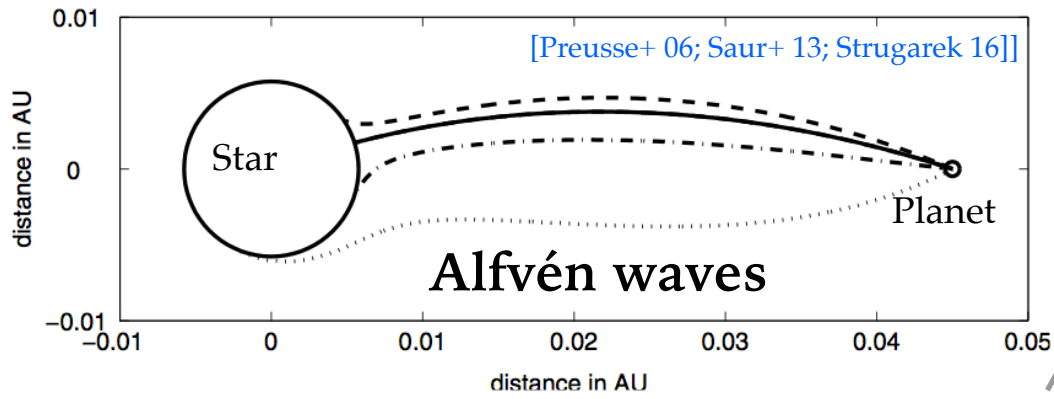


# Two magnetic interaction regimes



**Sub**-alfvénic interaction  
Star-planet **connection**

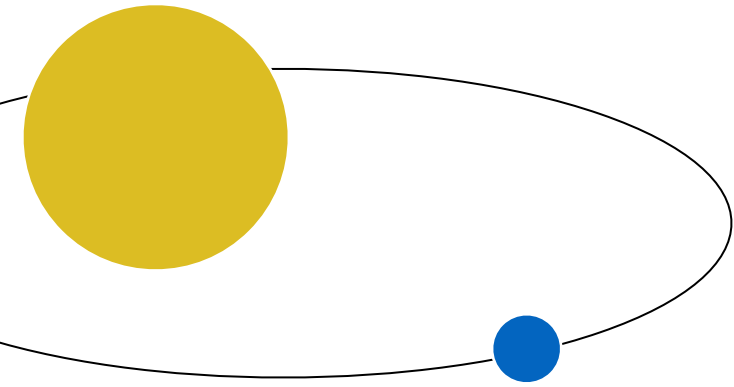
*Alfvén radius/surface*



[Gombosi 98]

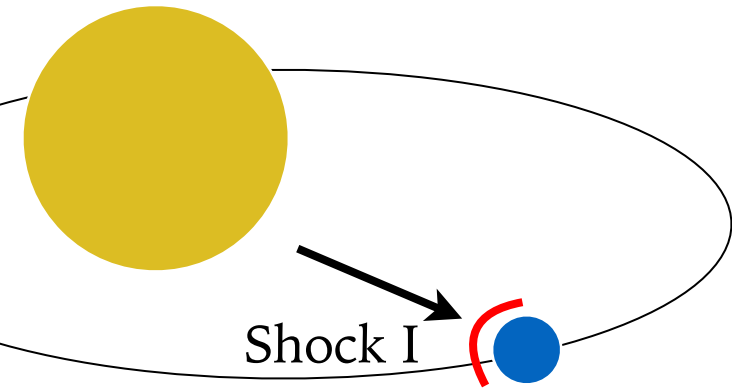
**Super**-alfvénic interaction  
**Shocks**

# What are the necessary criteria for shocks?





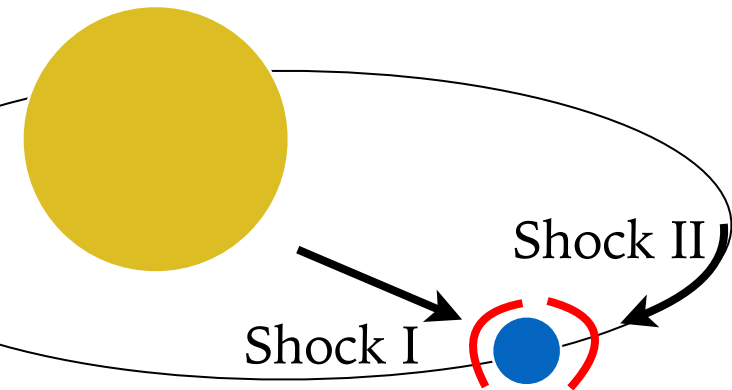
# What are the necessary criteria for shocks?



## Shock I

*Radial wind velocity*  $> v_A$

# What are the necessary criteria for shocks?



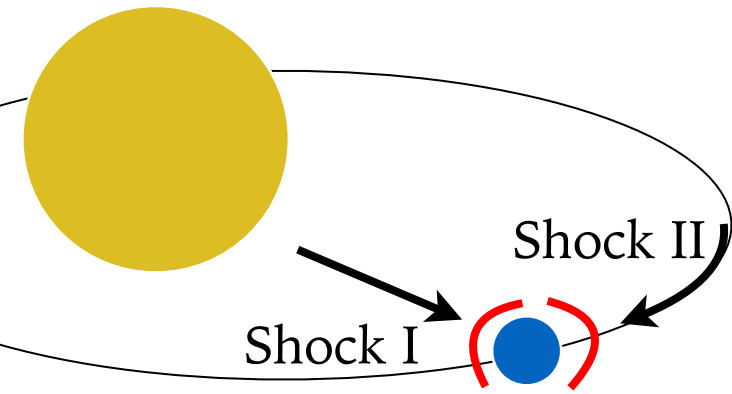
## Shock I

*Radial wind velocity*  $> v_A$

## Shock II

*|Keplerian motion - rotation|*  $> v_A$

# What are the necessary criteria for shocks?



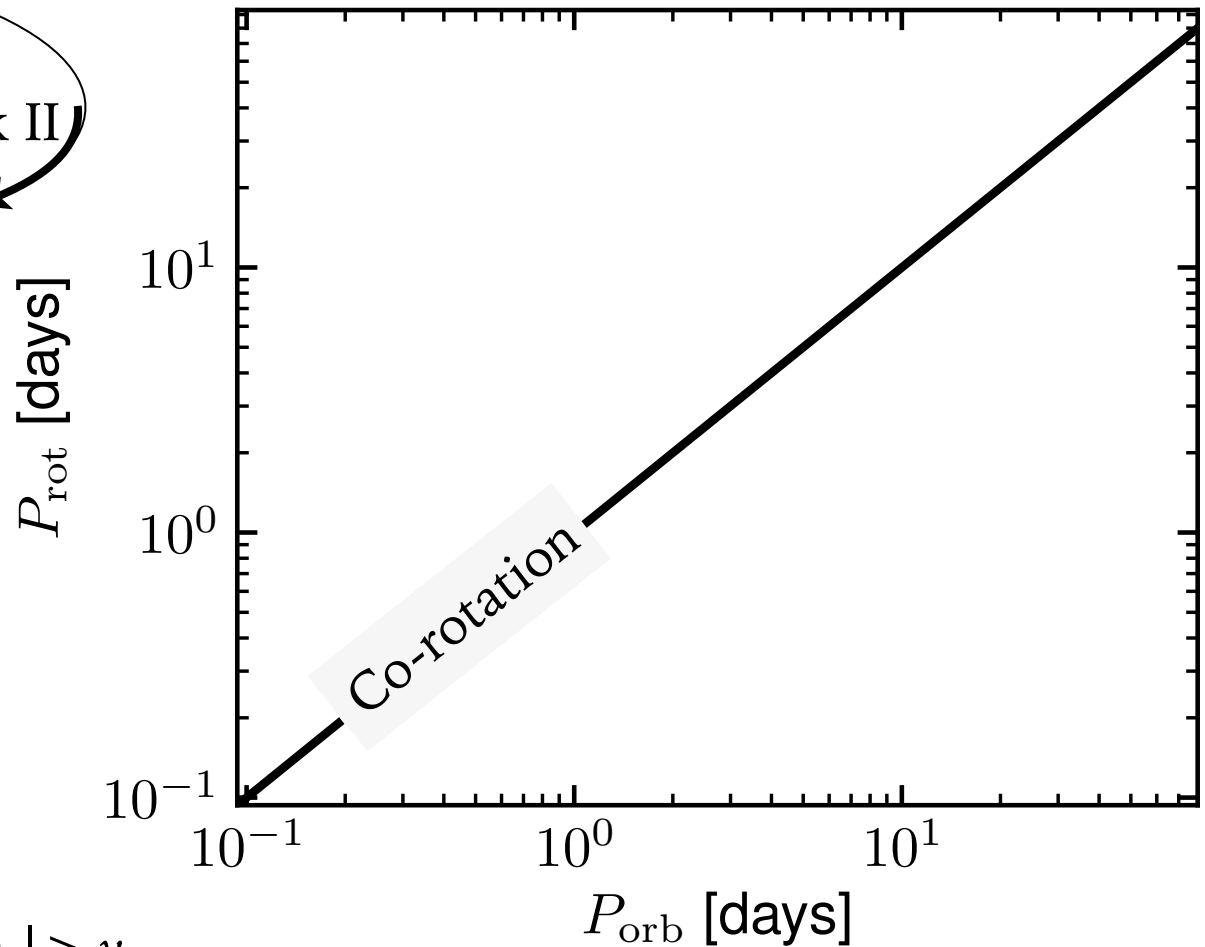
## Shock I

*Radial wind velocity*  $> v_A$

## Shock II

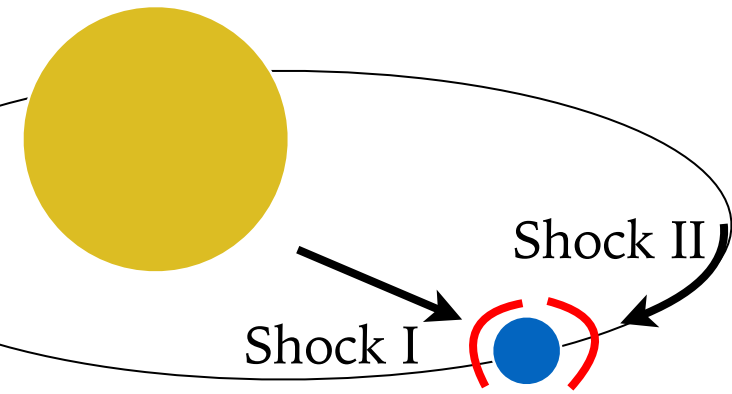
$| \textit{Keplerian motion} - \textit{rotation} | > v_A$

K-type star (e.g. HD 189733)



(Weber-Davis wind model)

# What are the necessary criteria for shocks?



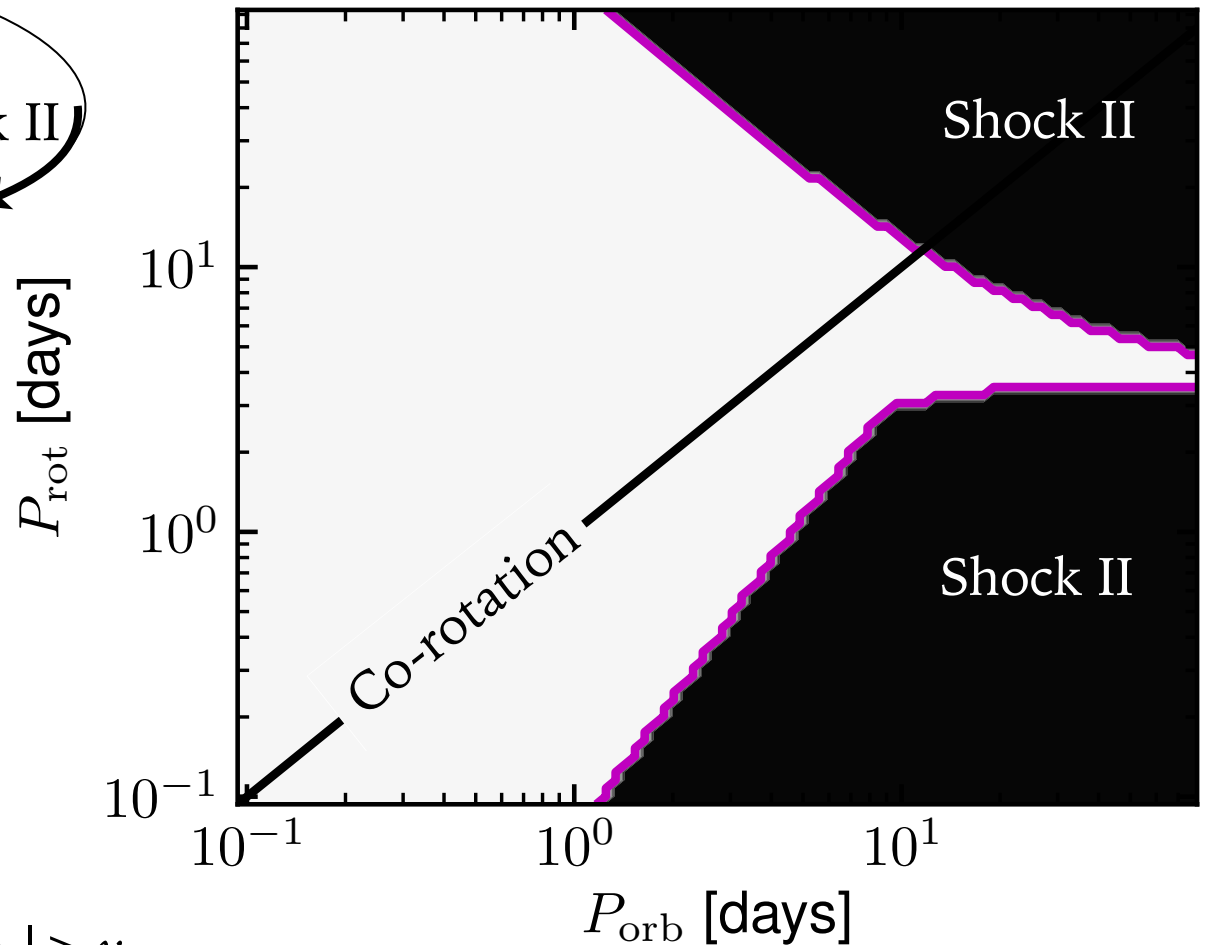
## Shock I

*Radial wind velocity*  $> v_A$

## Shock II

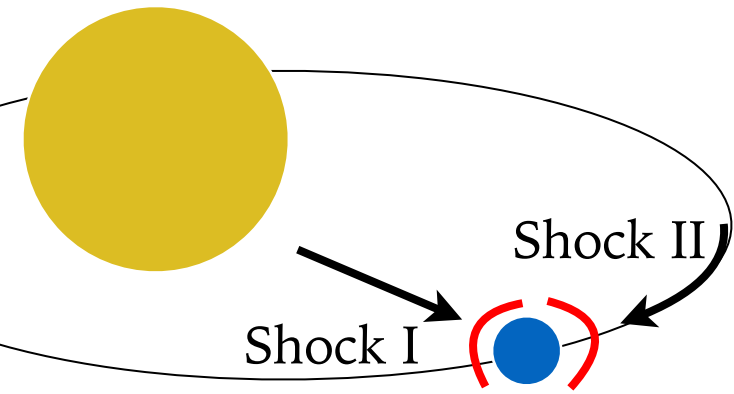
$| \textit{Keplerian motion} - \textit{rotation} | > v_A$

K-type star (e.g. HD 189733)



(Weber-Davis wind model)

# What are the necessary criteria for shocks?



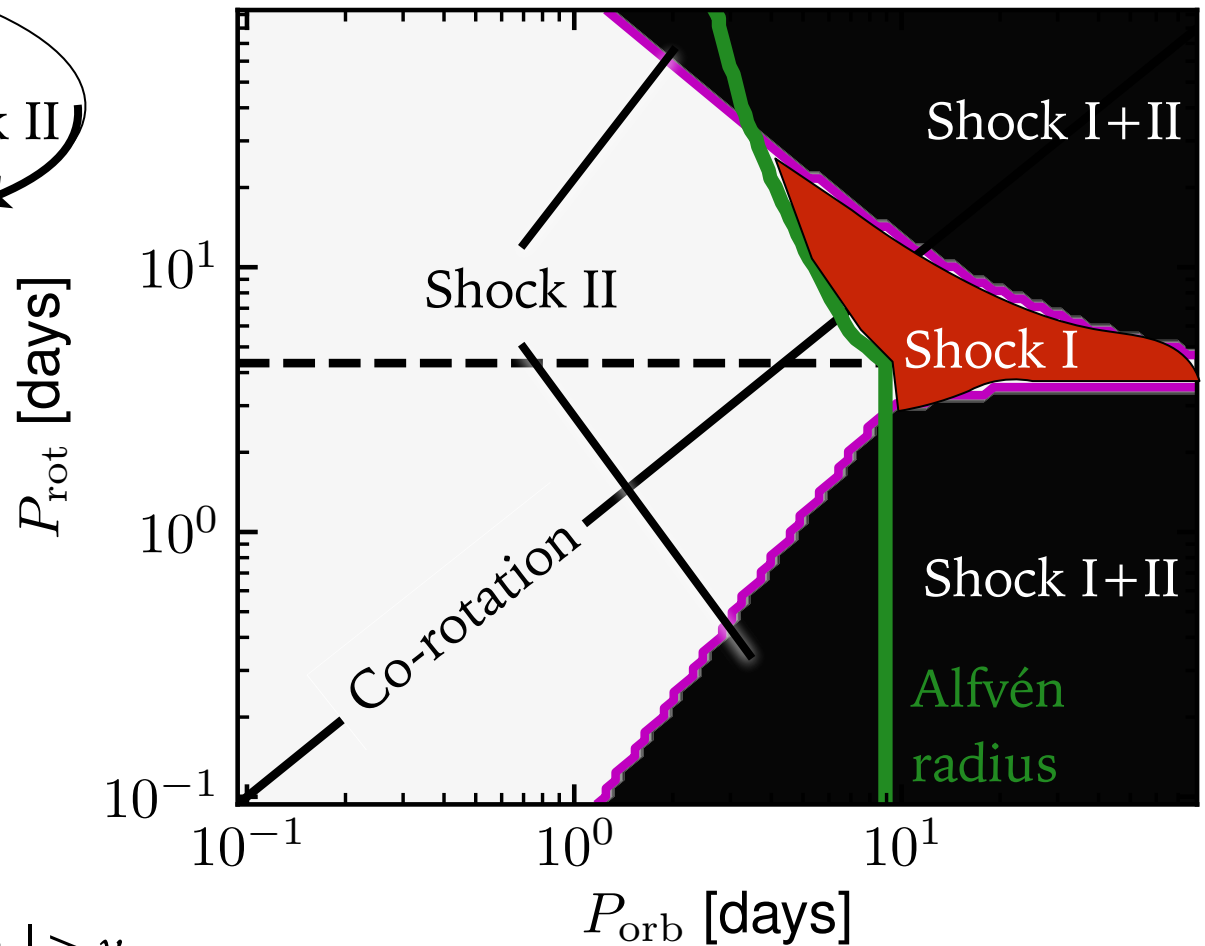
## Shock I

*Radial wind velocity*  $> v_A$

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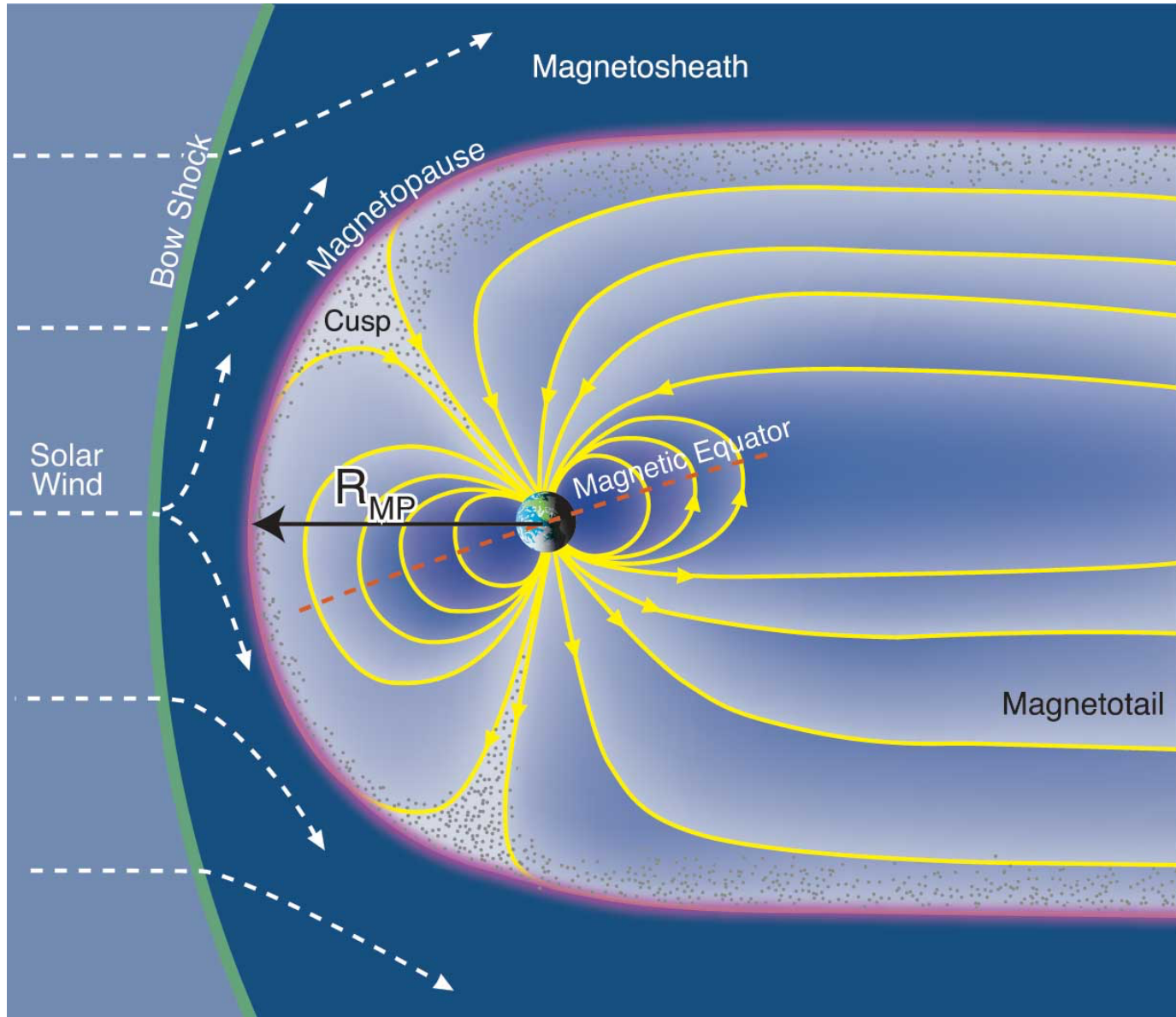


(Weber-Davis wind model)

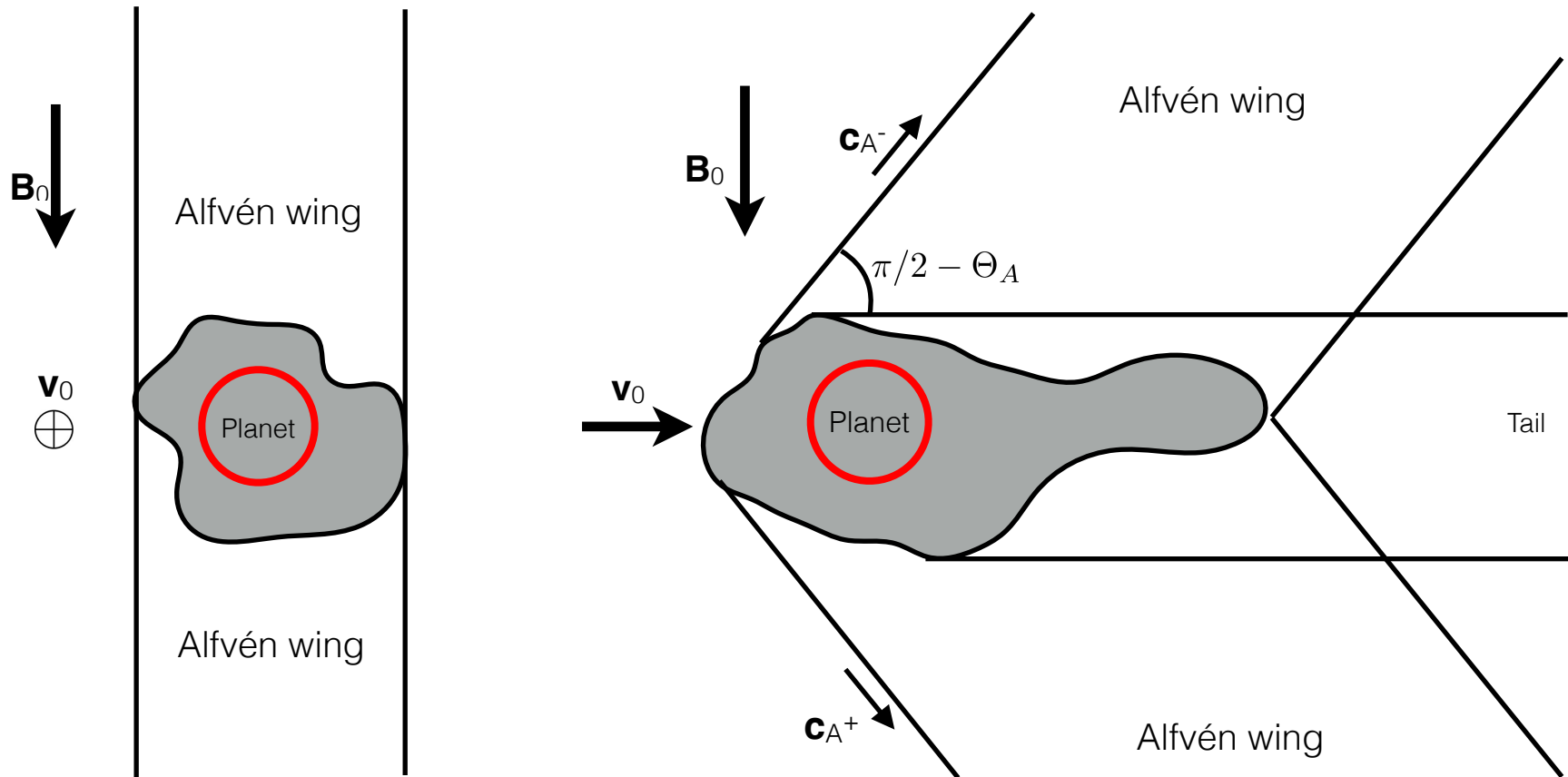
# Nomenclature of (magnetized) interactions

	Flow		
Obstacle		Weakly/non magnetized	Strongly magnetized
Weakly/non magnetized		<i>Venus, Mars</i>	Unipolar interaction <i>(Io, hot exoplanets?)</i>
Strongly magnetized		<i>Earth, Jupiter</i>	Dipolar interaction <i>(hot exoplanets?)</i>
		Super-Alfvénic	Sub-Alfvénic

# Super-alfvénic regime: the case of the Earth



# The concept of Alfvén wings

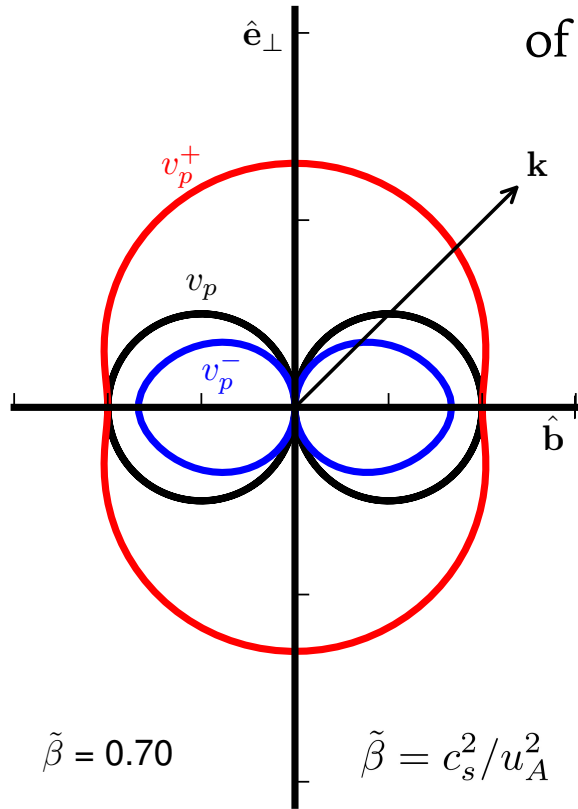


Alfvén wings are the **superposition** of Alfvén waves. In general, all types of magneto-sonic waves are triggered by the interaction between the obstacle and its environment.



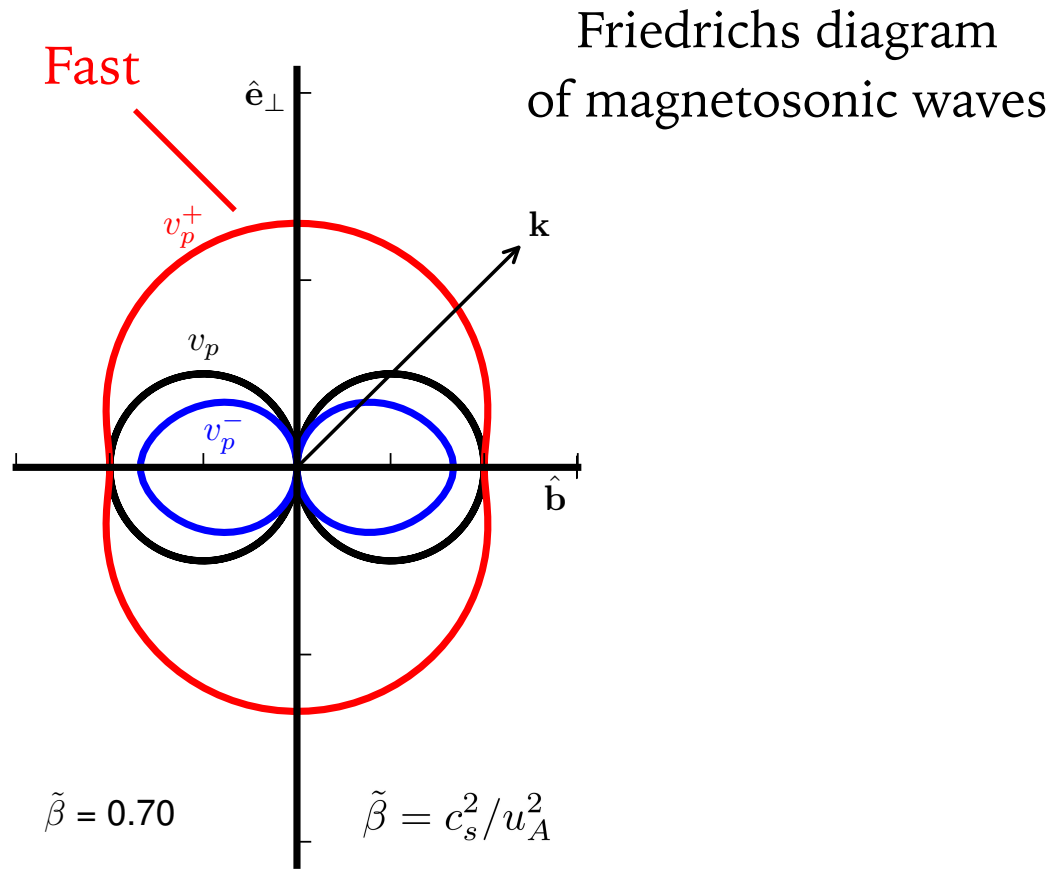
# Side note on magneto-sonic waves (I)

Friedrichs diagram  
of magnetosonic waves



**Phase velocity** of magnetosonic waves:  
no preferred direction

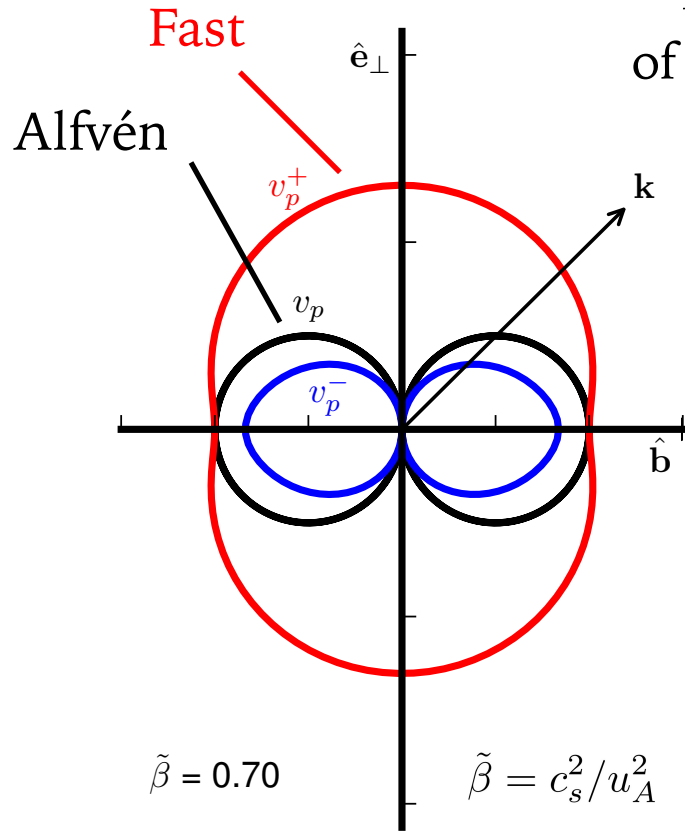
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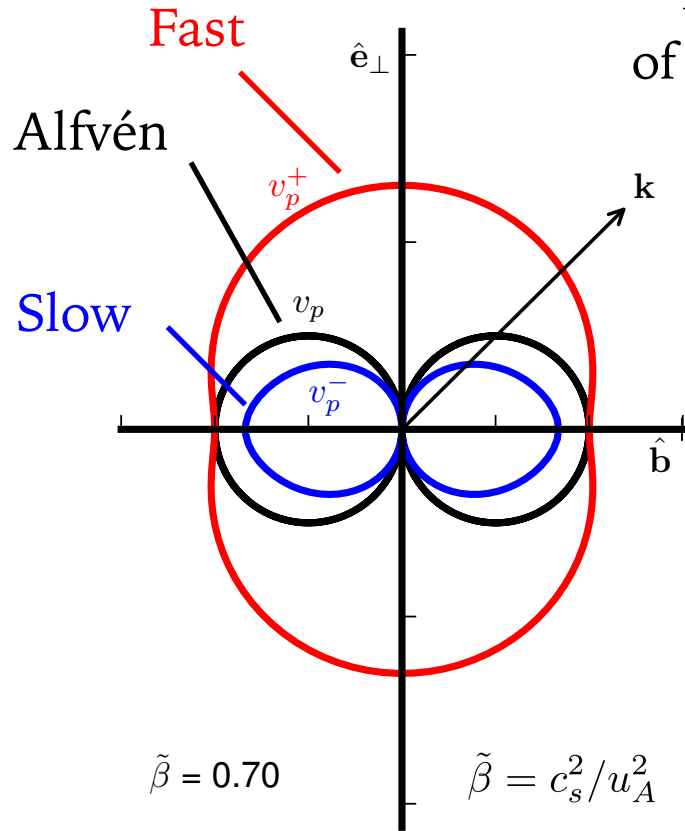
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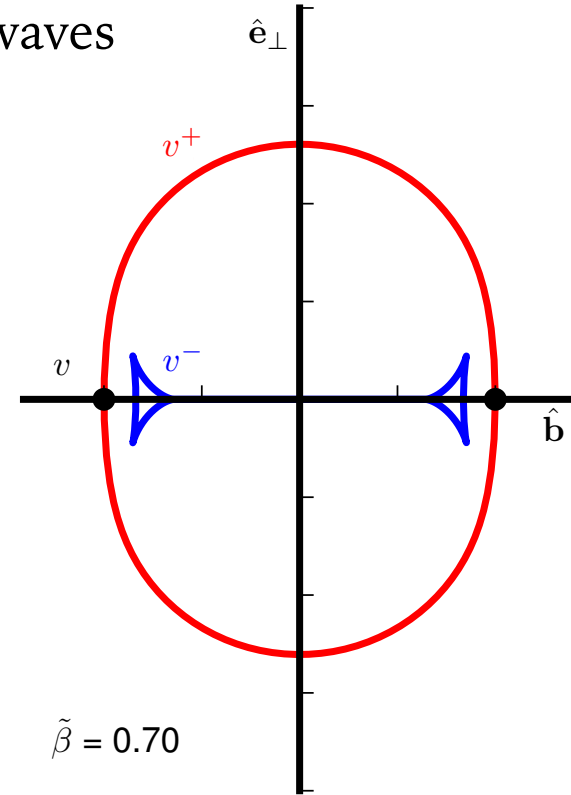
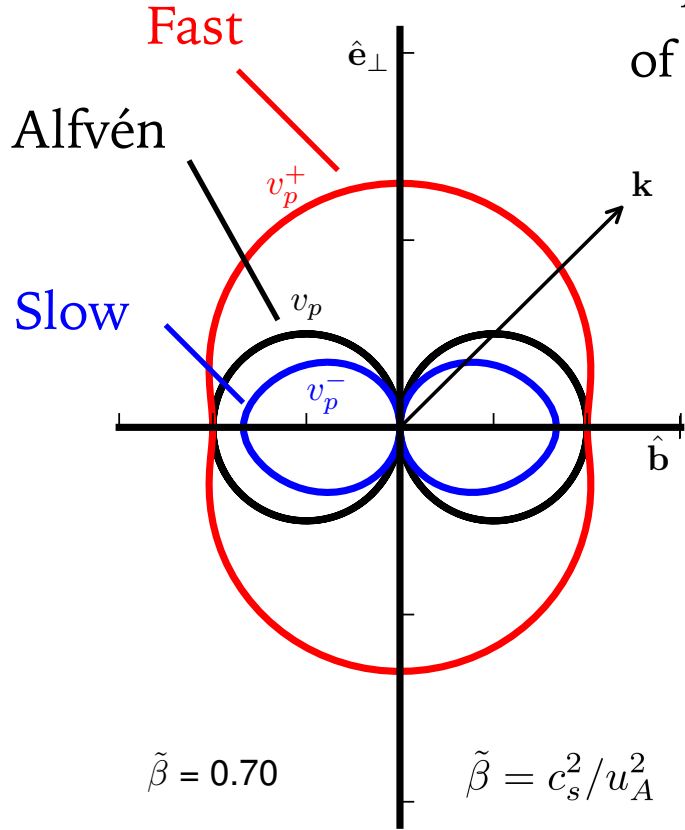
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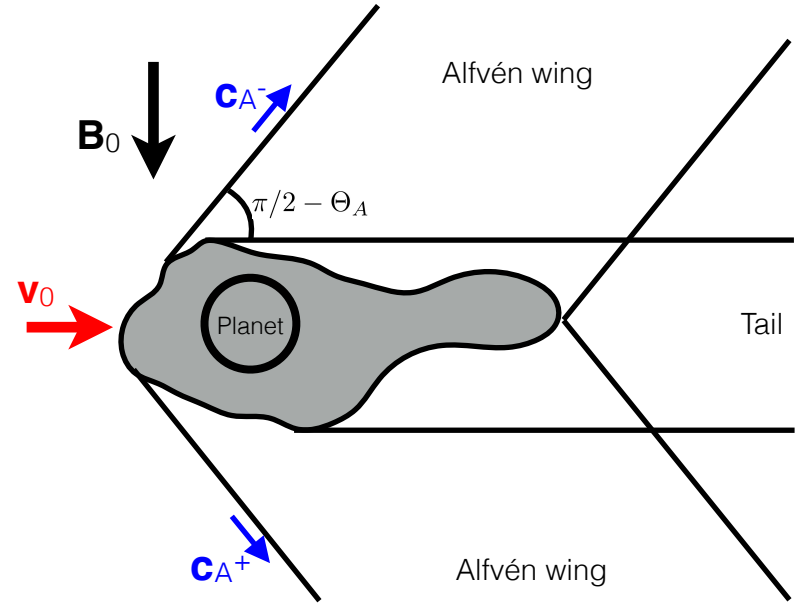
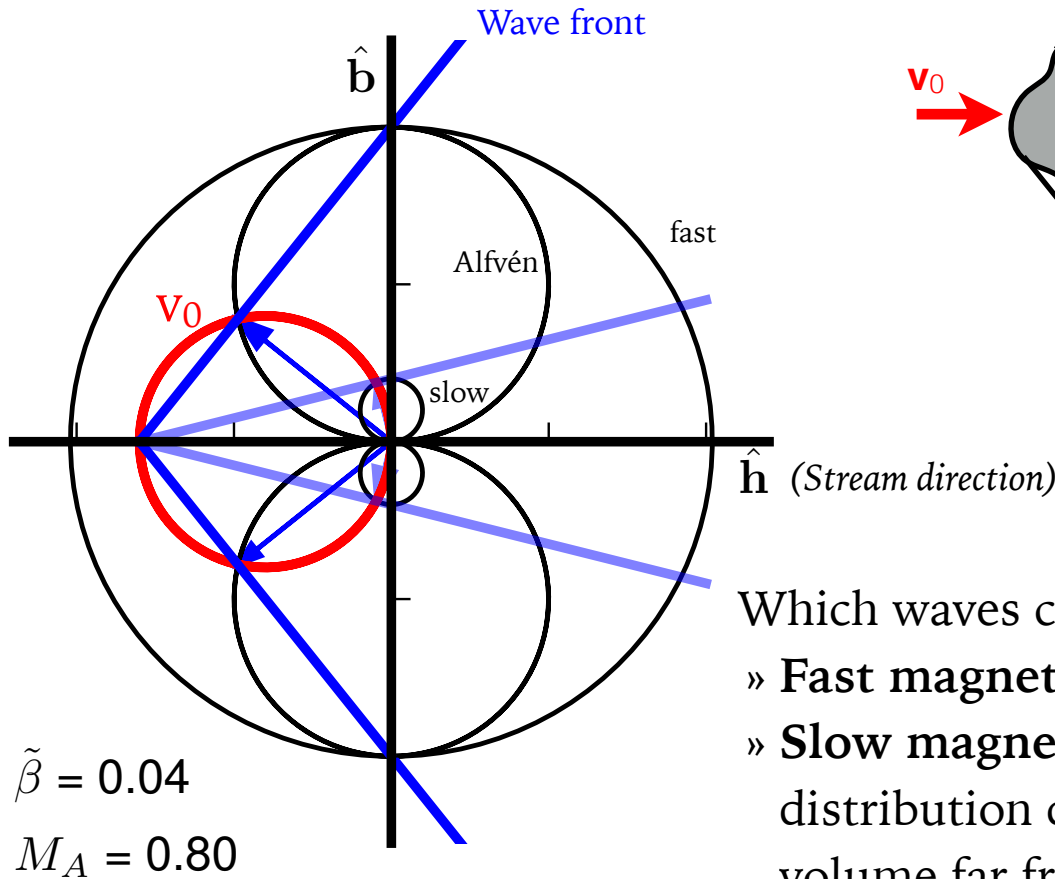


**Phase velocity** of magnetosonic waves:  
no preferred direction

**Group velocity** of magnetosonic waves:  
the pure Alfvén wave (black dots) is degenerate,  
energy is always transported along  $\mathbf{b}$

# Side note on magneto-sonic waves (II)

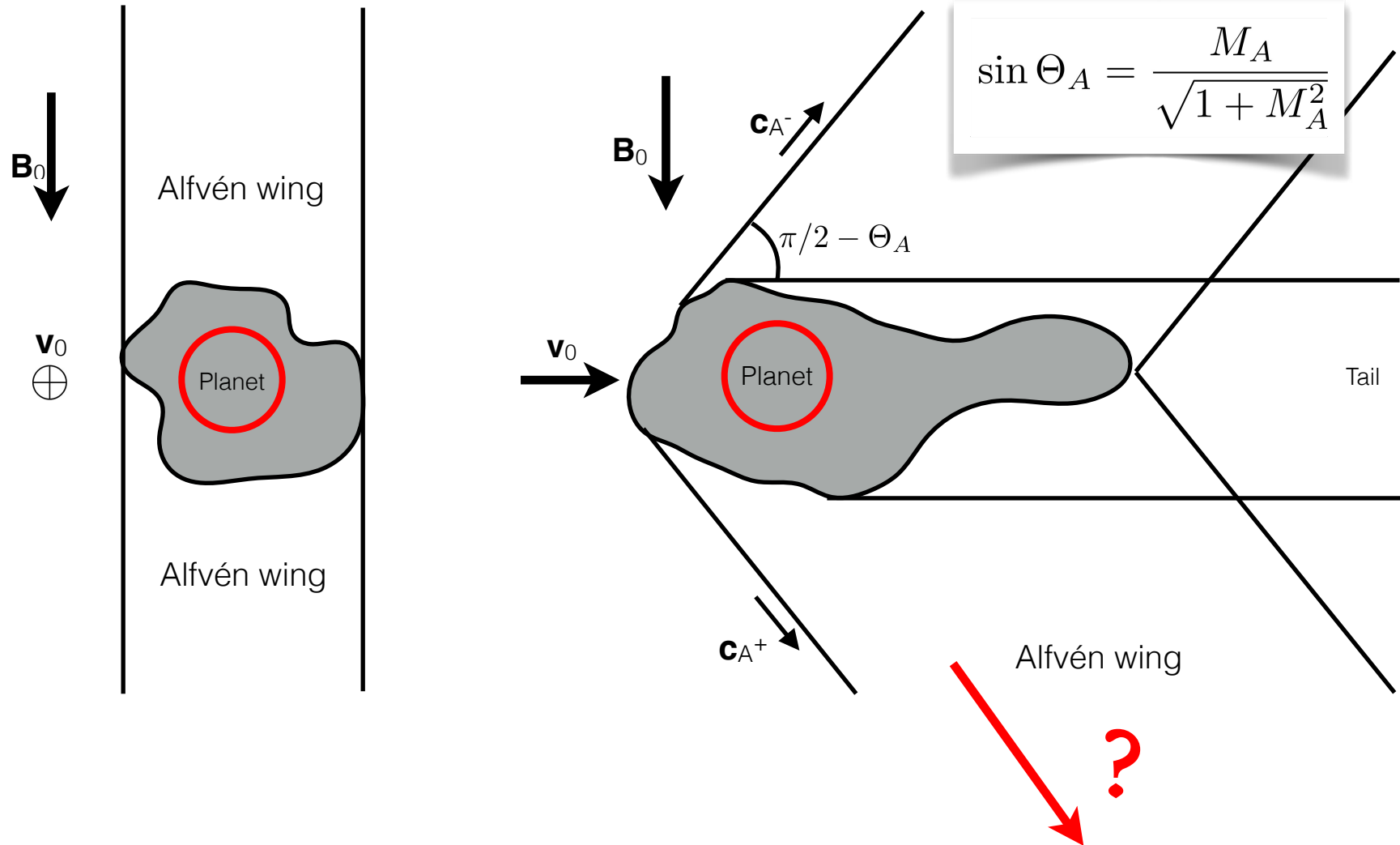
Friedrichs diagram  
in the frame where the obstacle is at rest



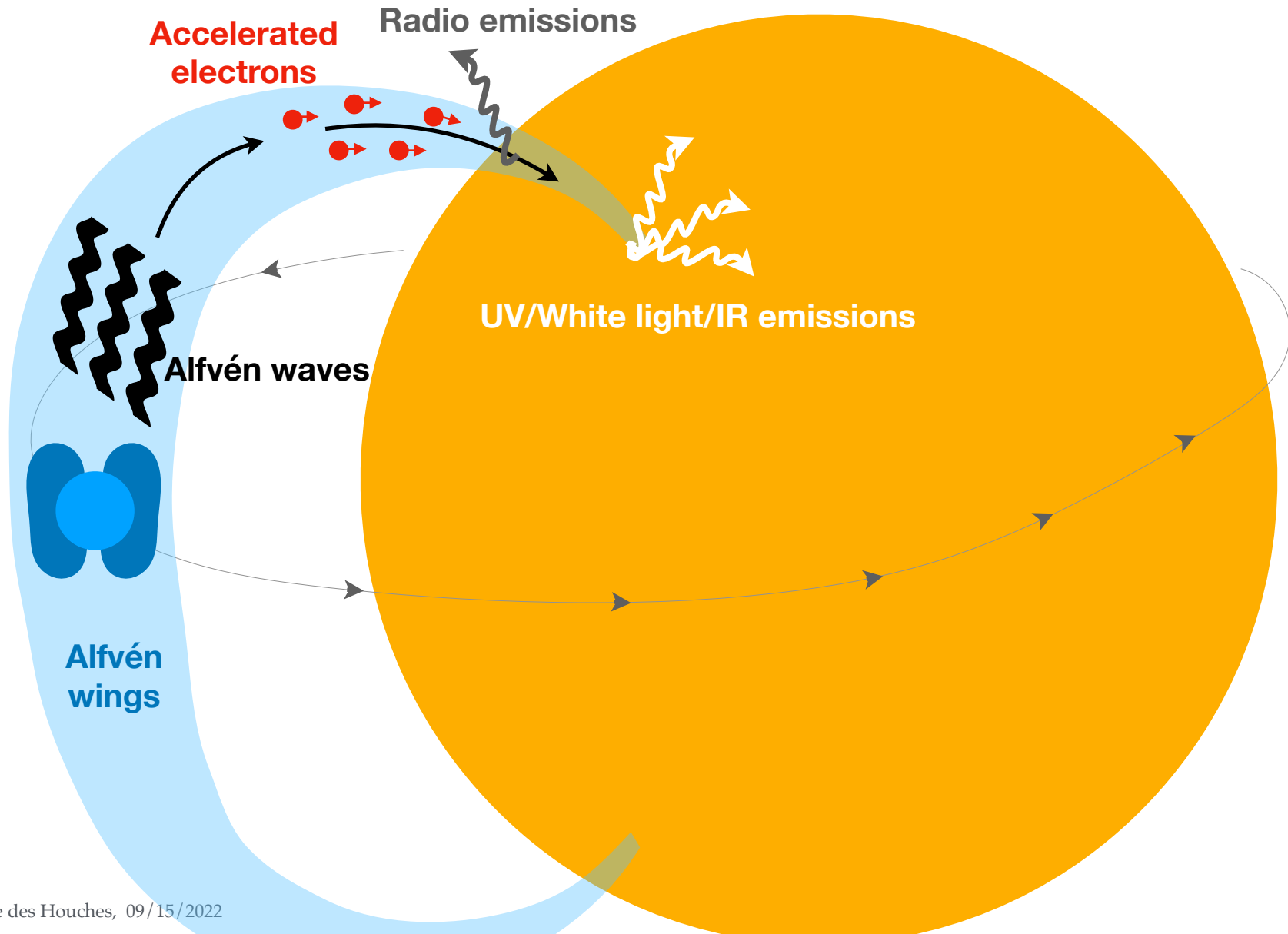
Which waves can actually propagate?

- » **Fast magnetosonic:** no intersection ( $M_A < 1$ )
- » **Slow magnetosonic:** the group velocity distribution dilutes the wave energy over a large volume far from the obstacle
- » **Alfvén:** propagate along  $\mathbf{c}_A^\pm = \mathbf{v}_0 \pm \mathbf{u}_A$

# The concept of Alfvén wings

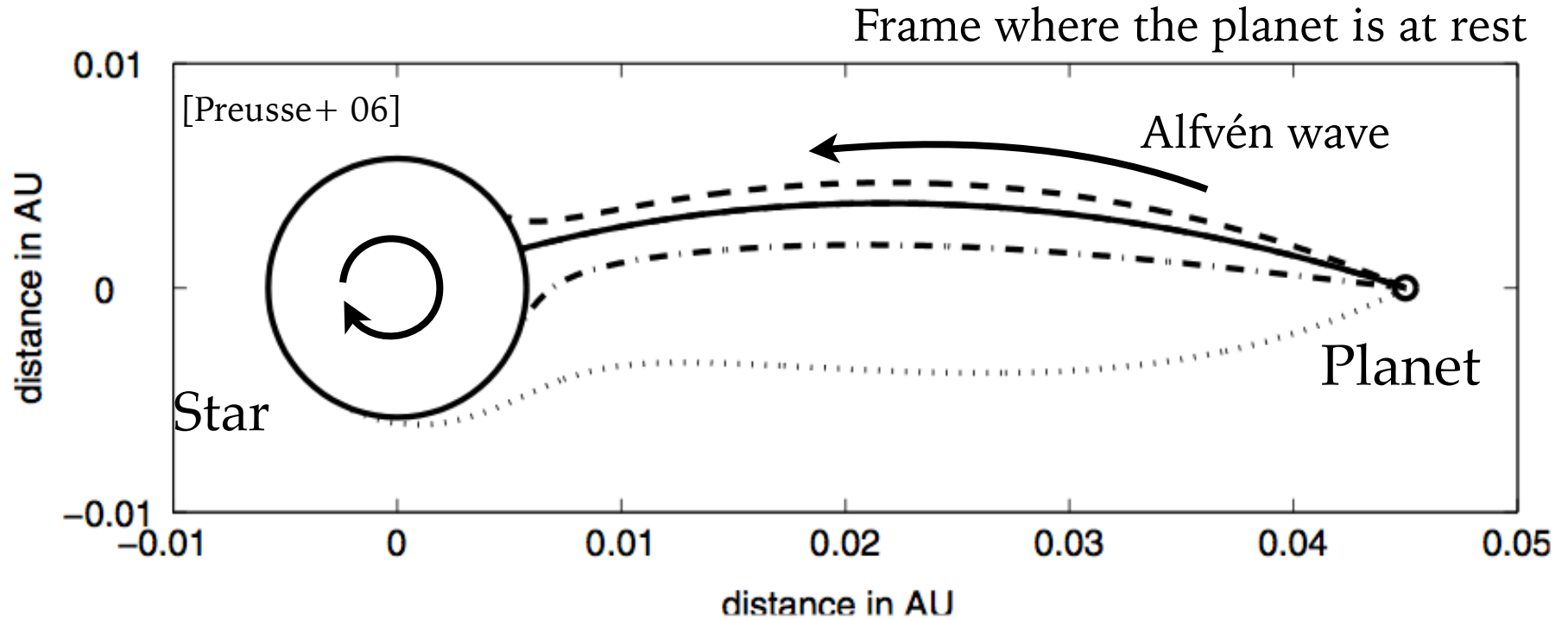


# The concept of Alfvén wings: a global picture

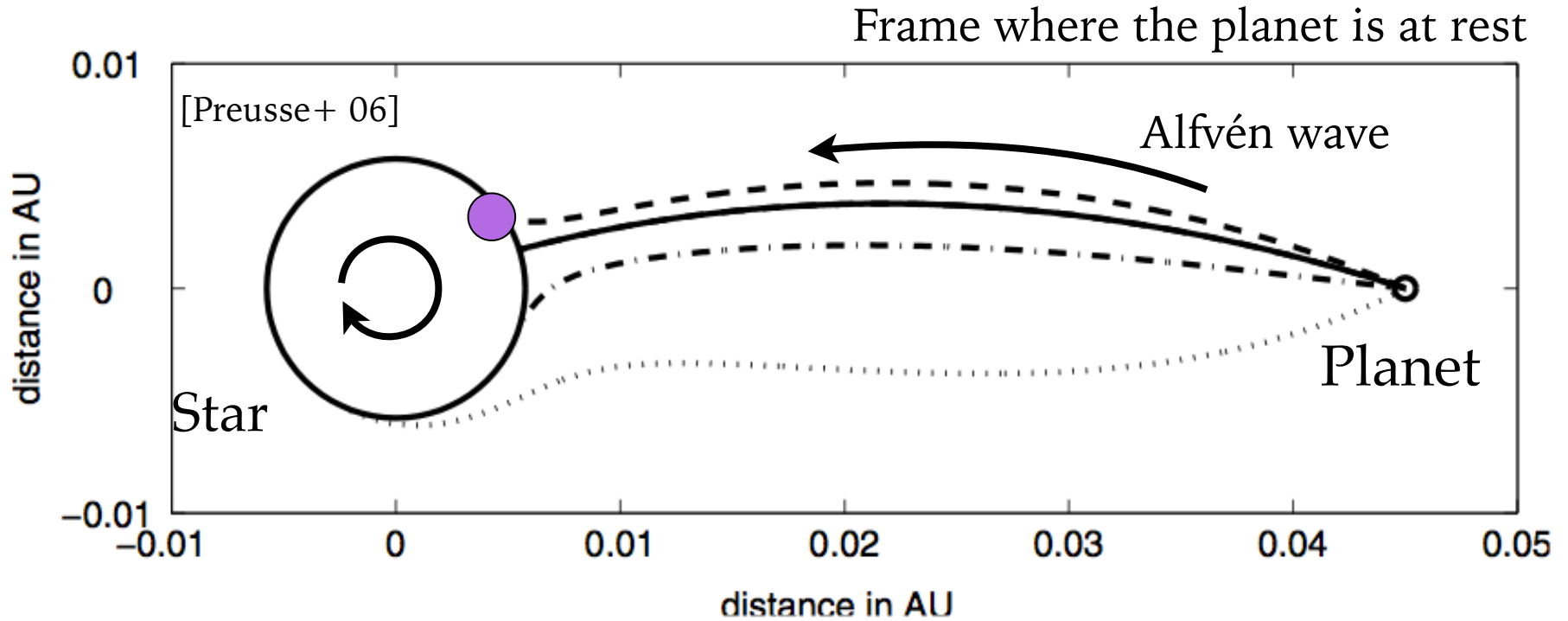




# Sub-alfvénic regime: magnetic connectivity

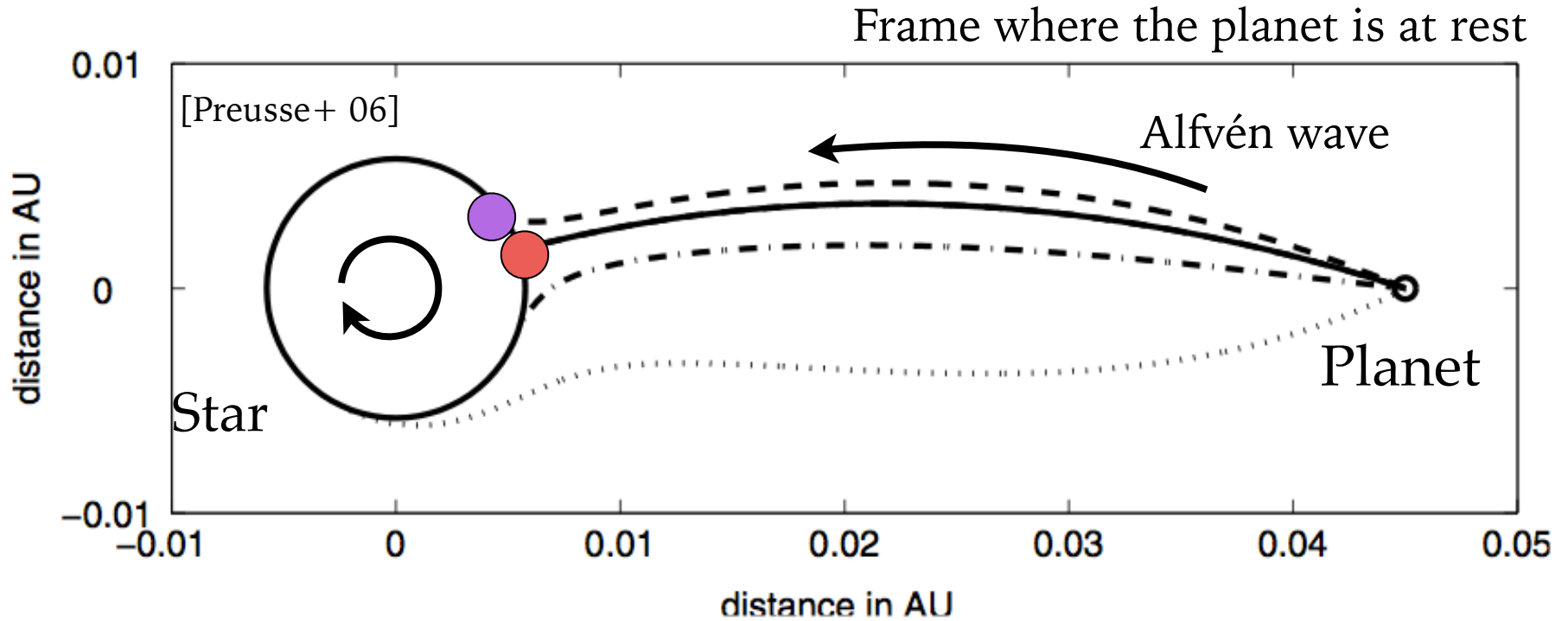


# Sub-alfvénic regime: magnetic connectivity



●  $P_{\text{orb}} = 12$  days

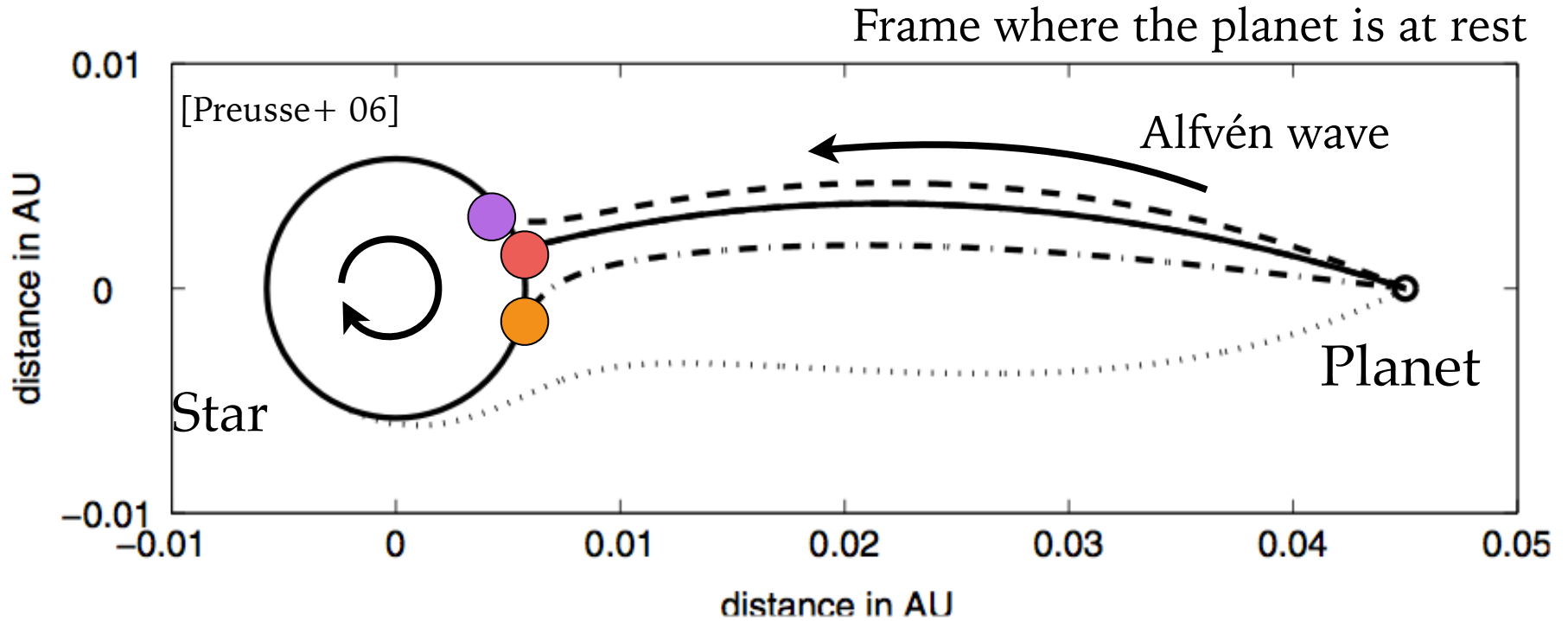
# Sub-alfvénic regime: magnetic connectivity



●  $P_{\text{orb}} = 12$  days

●  $P_{\text{orb}} = 9$  days

# Sub-alfvénic regime: magnetic connectivity

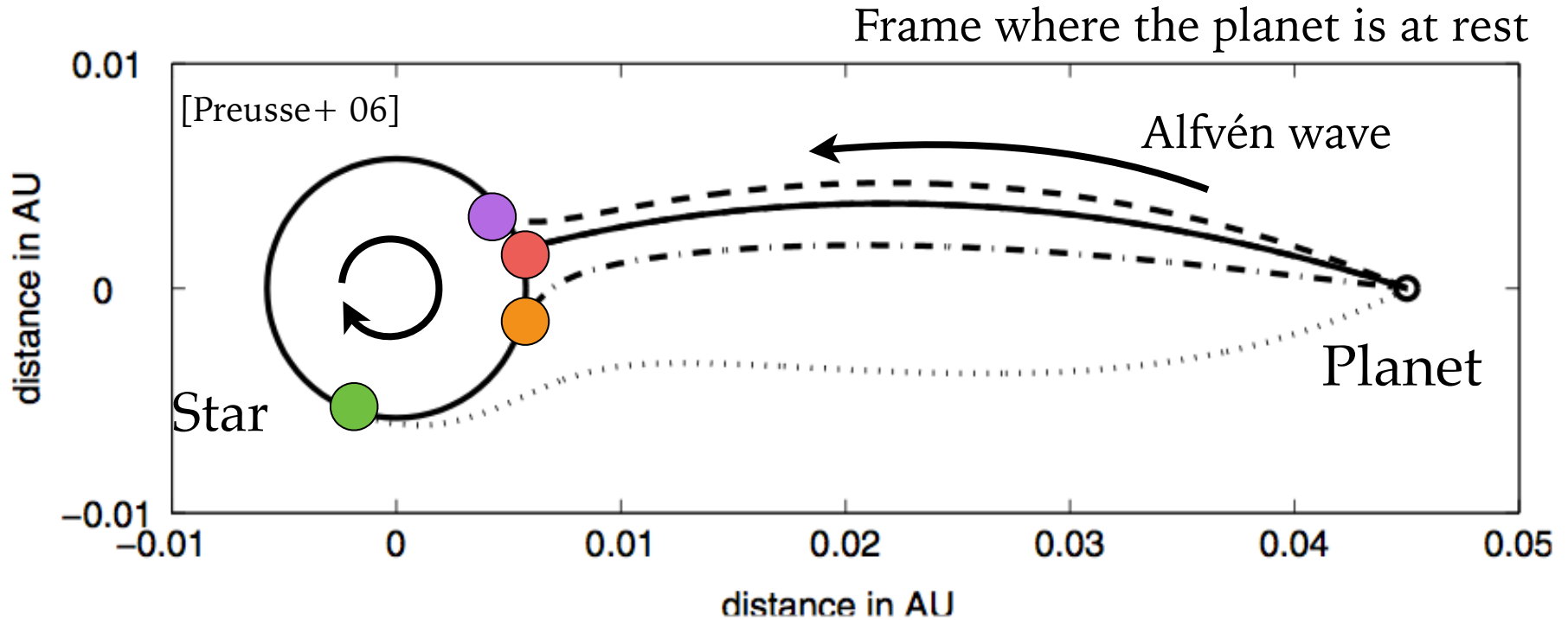


●  $P_{\text{orb}} = 12$  days

●  $P_{\text{orb}} = 9$  days

●  $P_{\text{orb}} = 6$  days

# Sub-alfvénic regime: magnetic connectivity



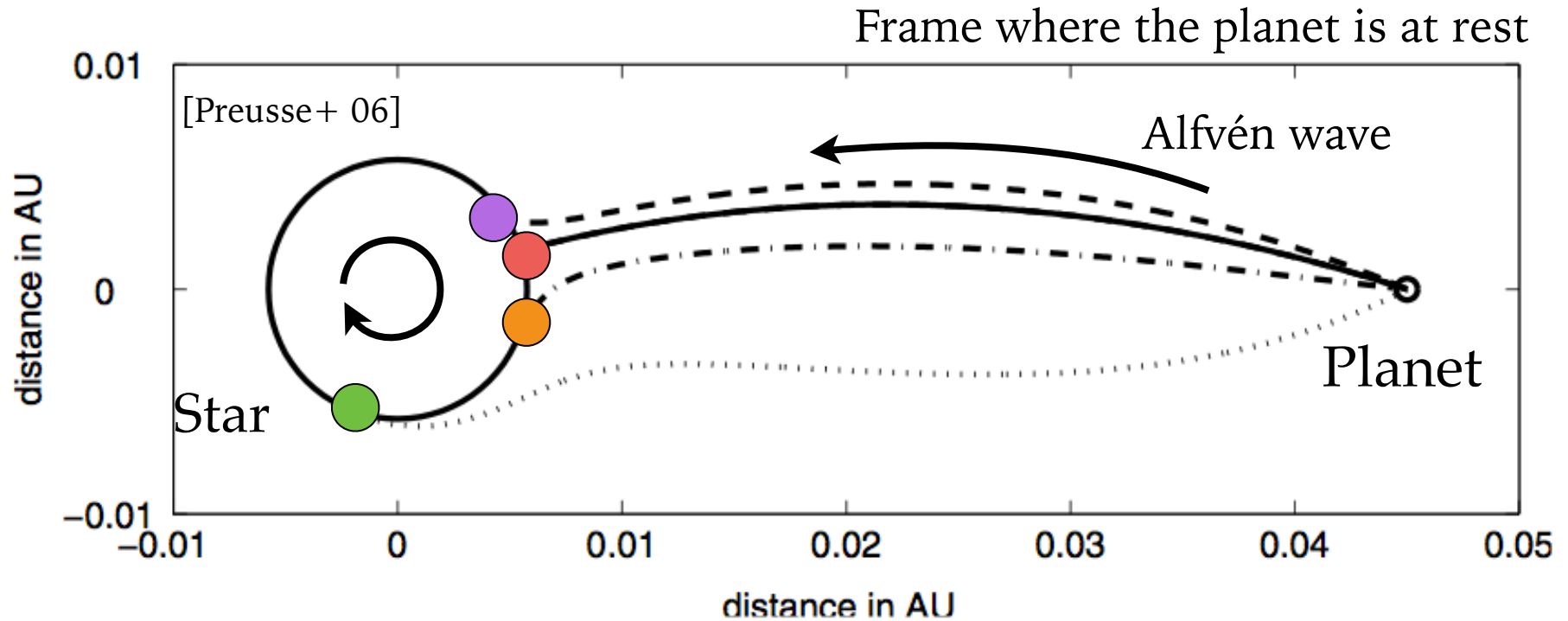
●  $P_{\text{orb}} = 12$  days

●  $P_{\text{orb}} = 9$  days

●  $P_{\text{orb}} = 6$  days

●  $P_{\text{orb}} = 3$  days

# Sub-alfvénic regime: magnetic connectivity



●  $P_{\text{orb}} = 12$  days

●  $P_{\text{orb}} = 9$  days

●  $P_{\text{orb}} = 6$  days

●  $P_{\text{orb}} = 3$  days

» The impact point can be dephased ( $\Delta\phi$ )

»  $\Delta\phi$  strongly depends on the **star**

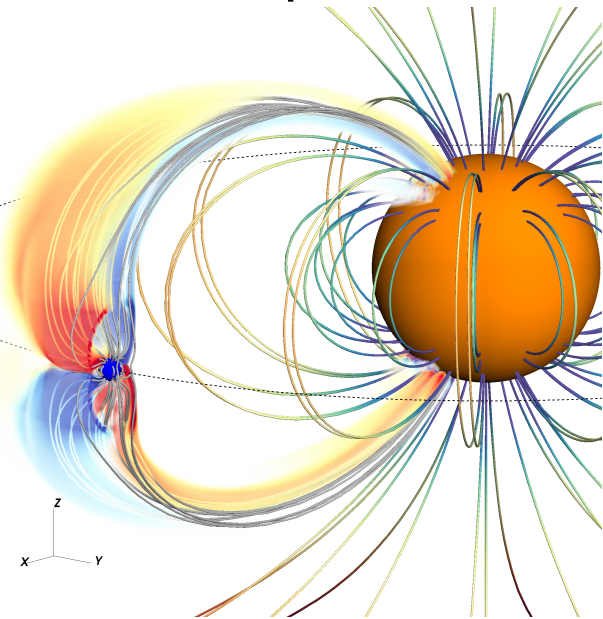
»  $B_{\star}$  (amplitude, topology)

»  $T_{\text{corona}}$  & stellar wind

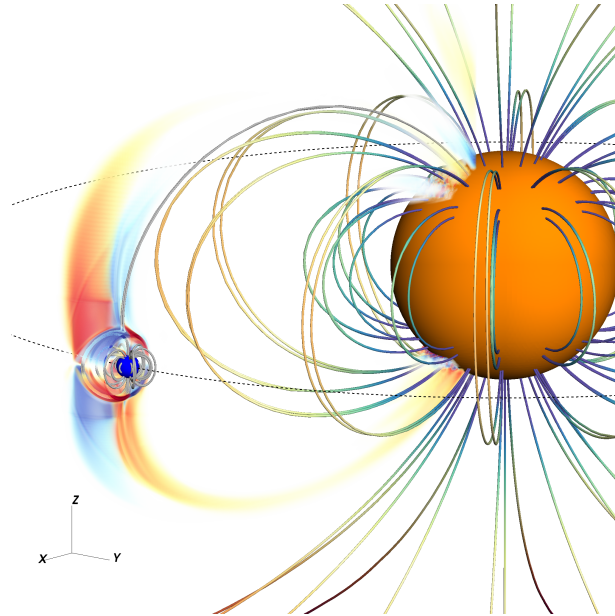
» Rotation

# Towards complex magnetic topologies

Stellar dipole

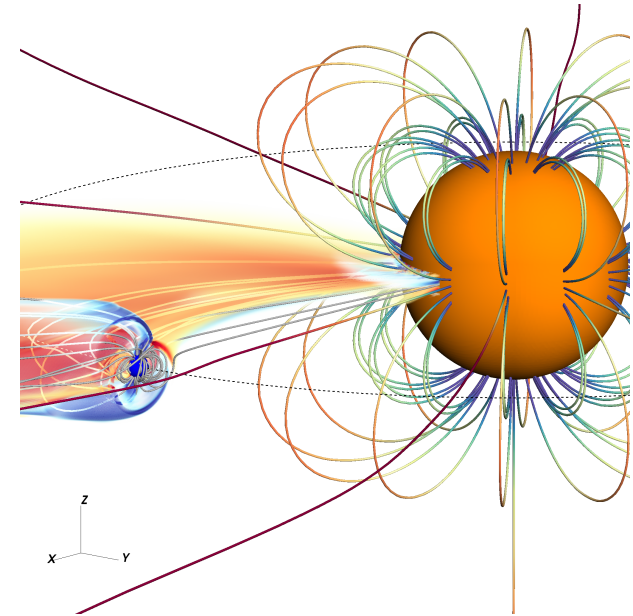


Two strong wings



Two weak wings

Stellar quadrupole

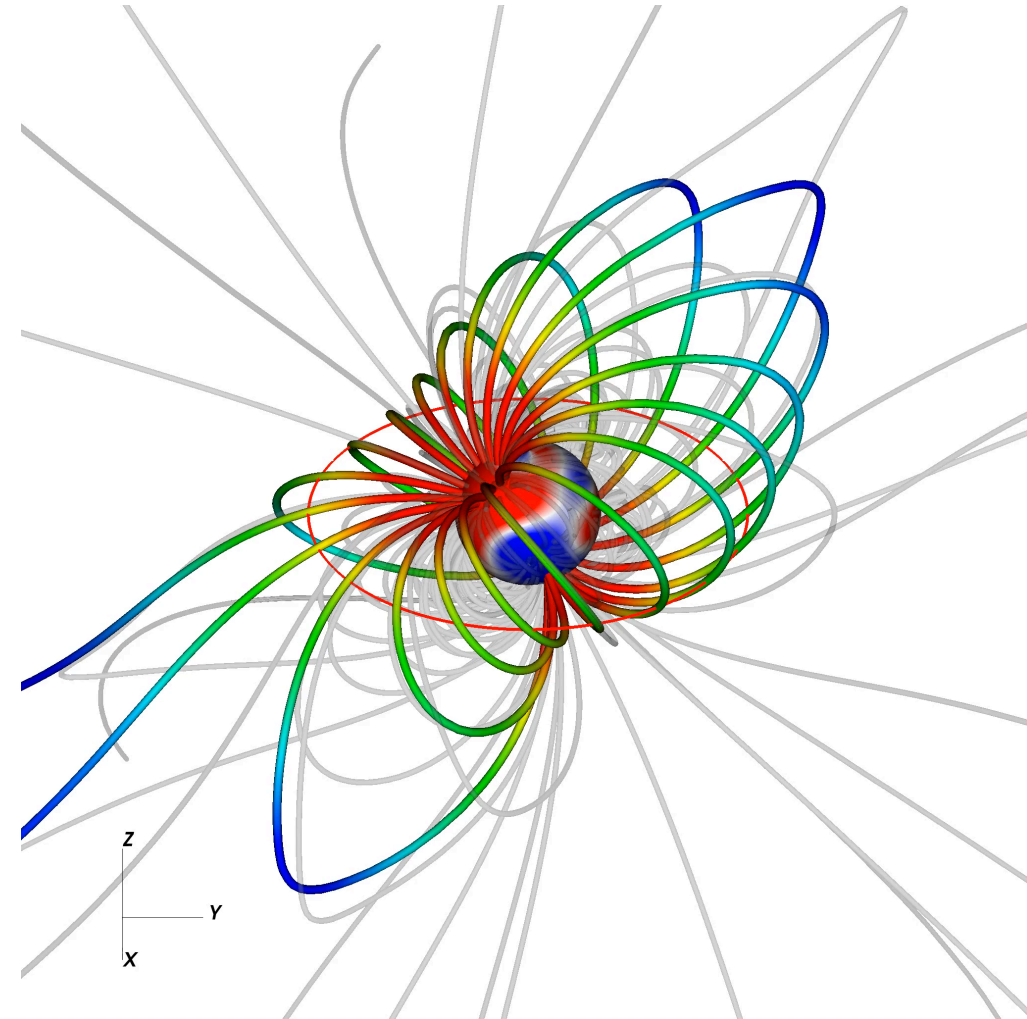


One strong wing

The interaction footpoint is at a given (latitude, longitude) position

The interaction footpoint is a longitudinal band at the equator

# Towards realistic complex magnetic topologies: Kepler-78



## Kepler-78 system

### *Stellar characteristics*

$T_{\text{eff}}$ [K]	$5089 \pm 50$
$M_{\star}$ [ $M_{\odot}$ ]	$0.81 \pm 0.08$
$R_{\star}$ [ $R_{\odot}$ ]	$0.74 +0.1, -0.08$
$P_{\text{rot}}$ [days]	12.5

---

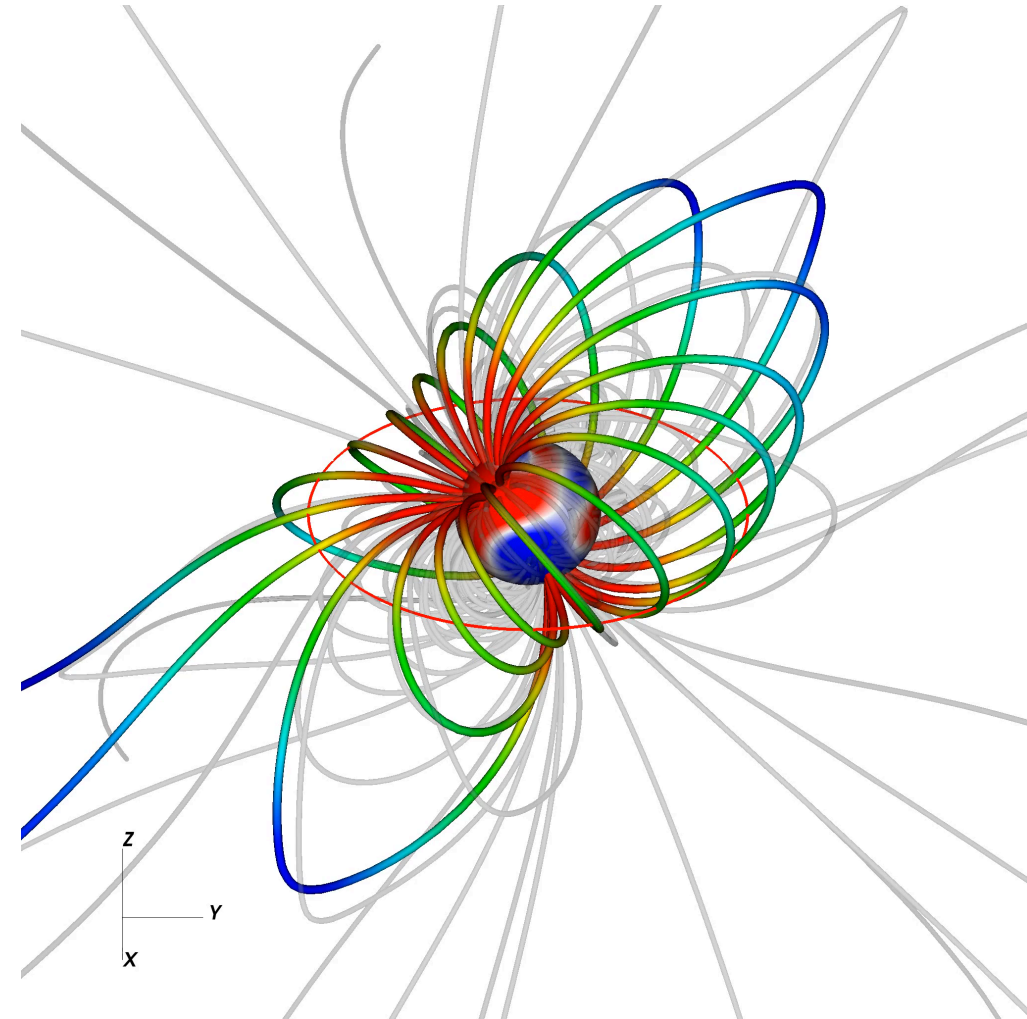
### *Planet characteristics*

$R_p$ [ $R_{\oplus}$ ]	$1.16 +0.19, -0.14$
$M_p^1$ [ $M_{\oplus}$ ]	$1.86 \pm 0.25$
$R_{\text{orb}}$ [ $R_{\star}$ ]	$3.0 +0.5, -1.0$

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# Towards realistic complex magnetic topologies: Kepler-78



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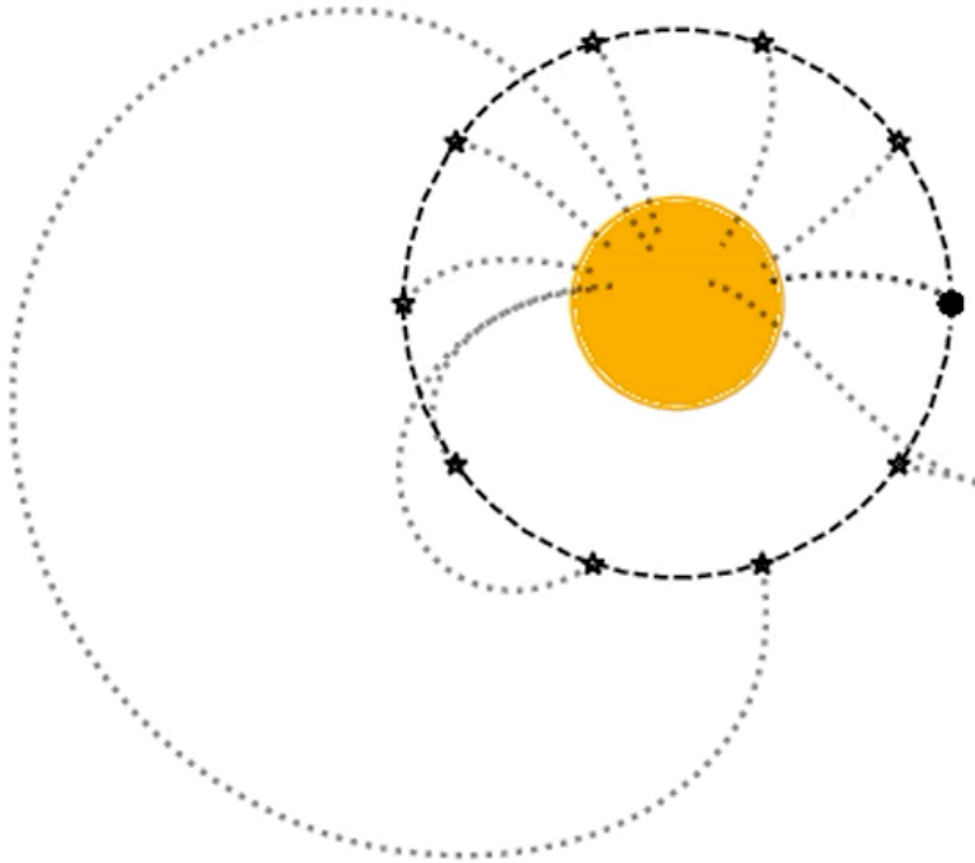
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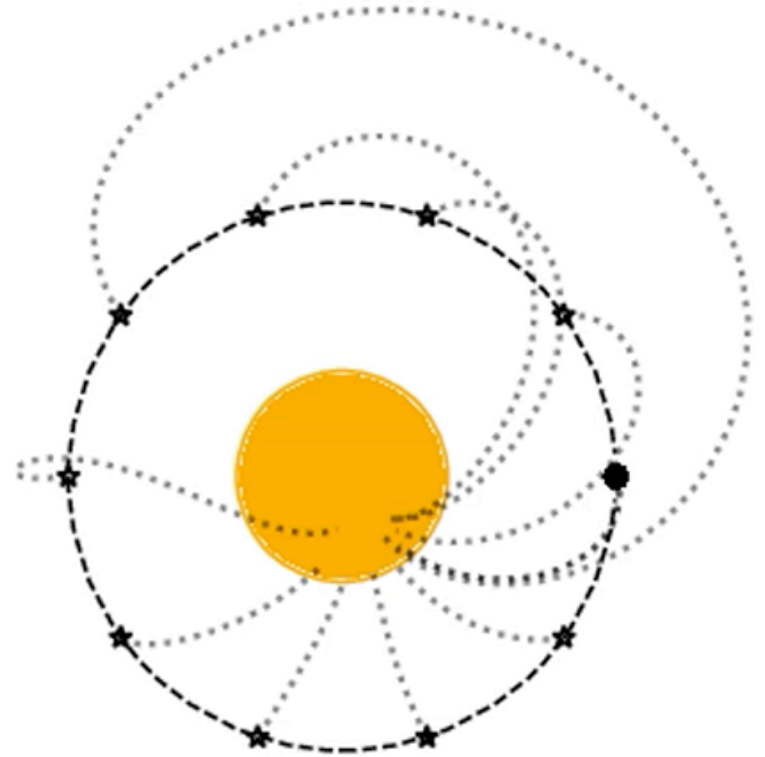
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# Towards realistic complex magnetic topologies: Kepler-78

Alfvén wing (+)



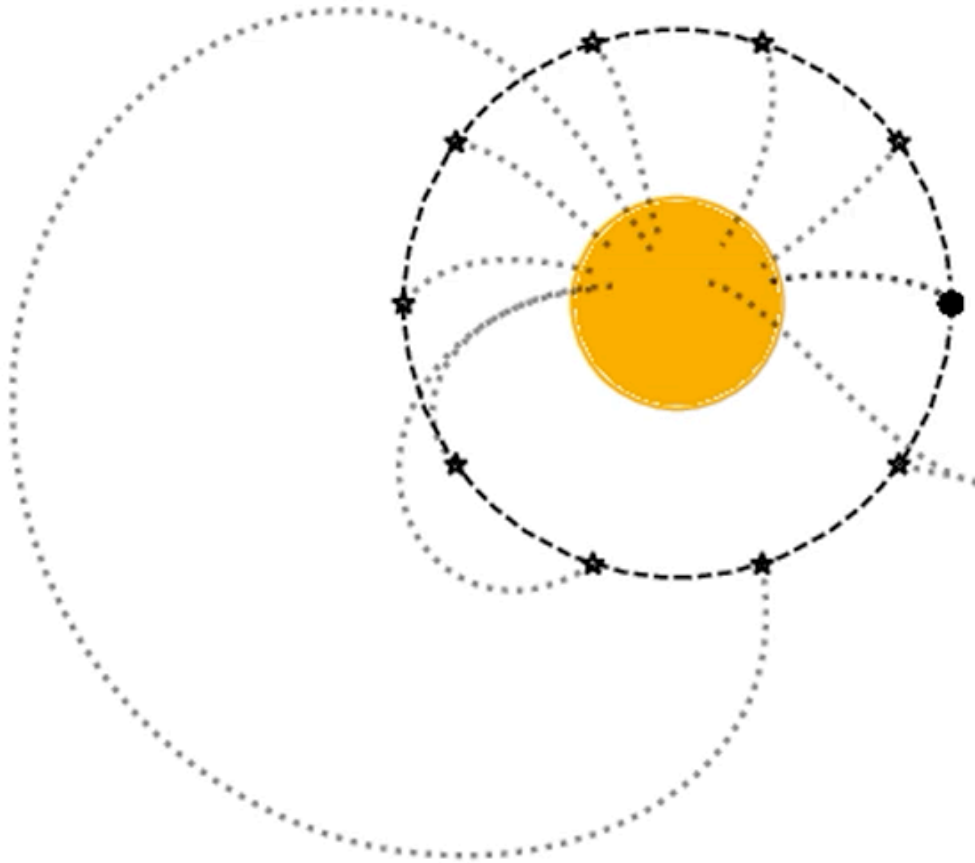
Alfvén wing (-)



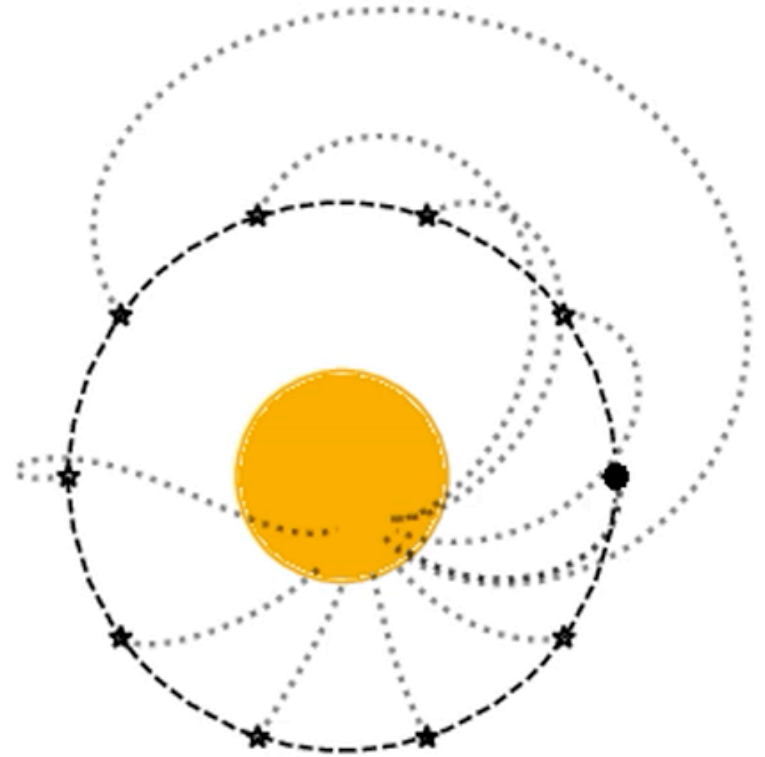
View from top

# Towards realistic complex magnetic topologies: Kepler-78

Alfvén wing (+)



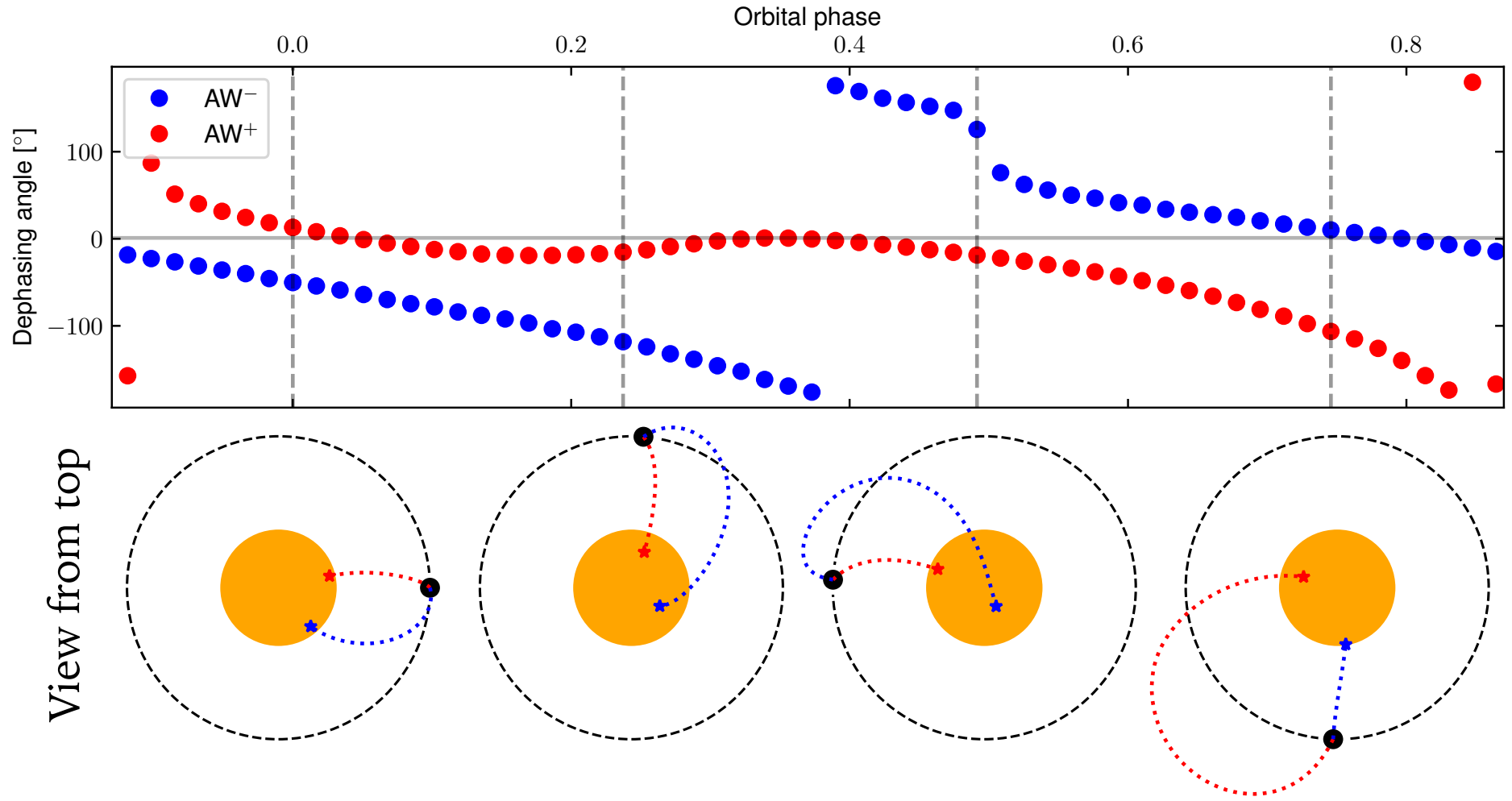
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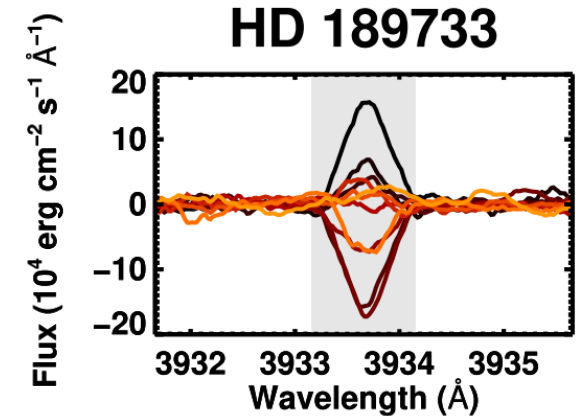
View from top

[Strugarek+ 19]

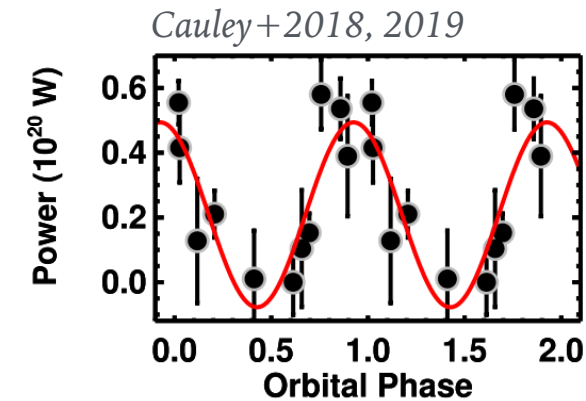
# Towards realistic complex magnetic topologies: Kepler-78



# Possible star-planet interaction in HD 189733



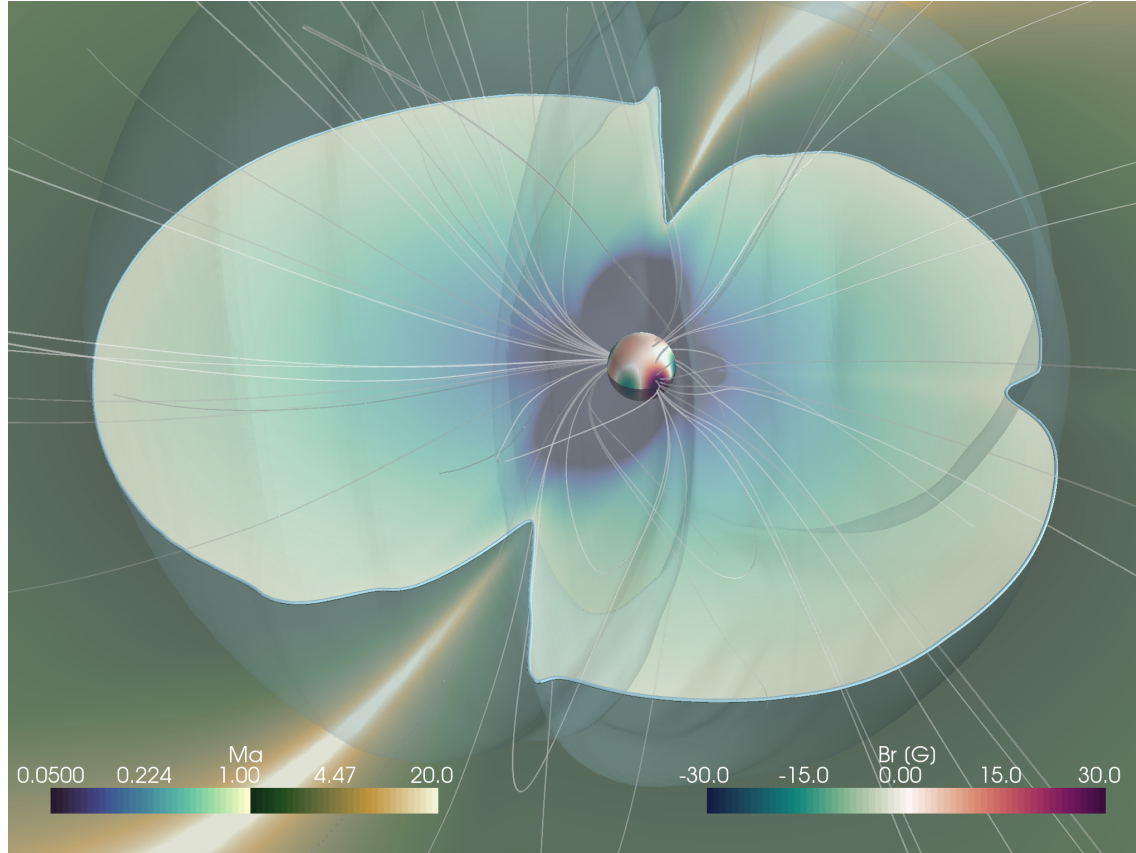
A signal in the Ca II H&K bands of the host star was observed to be correlated with the orbital phase of the hot exoplanet at **1 epoch out of 6 epochs studied** (Cauley+ 2018)



Applying star-planet interaction scaling law (stretch-and-break mechanism), Cauley+ 2019 deduced a field strength of about 80 G

Can we test this?

# Detailed modelling of HD 189733, August 2013



The hot Jupiter HD 189733b could orbit within the Alfvén surface

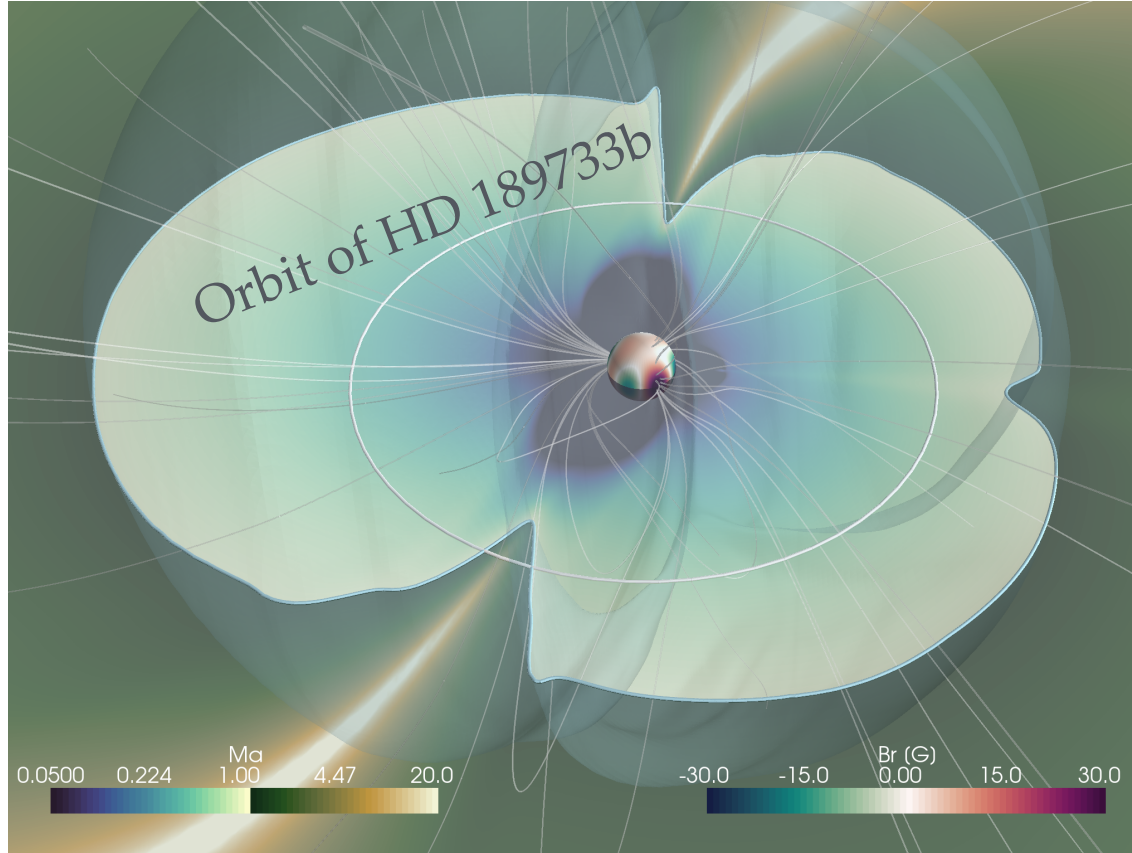
But the connectivity is modulated by the ‘complex’ magnetic topology of the star

[Strugarek+ 22]



Simulation data available on Galactica database

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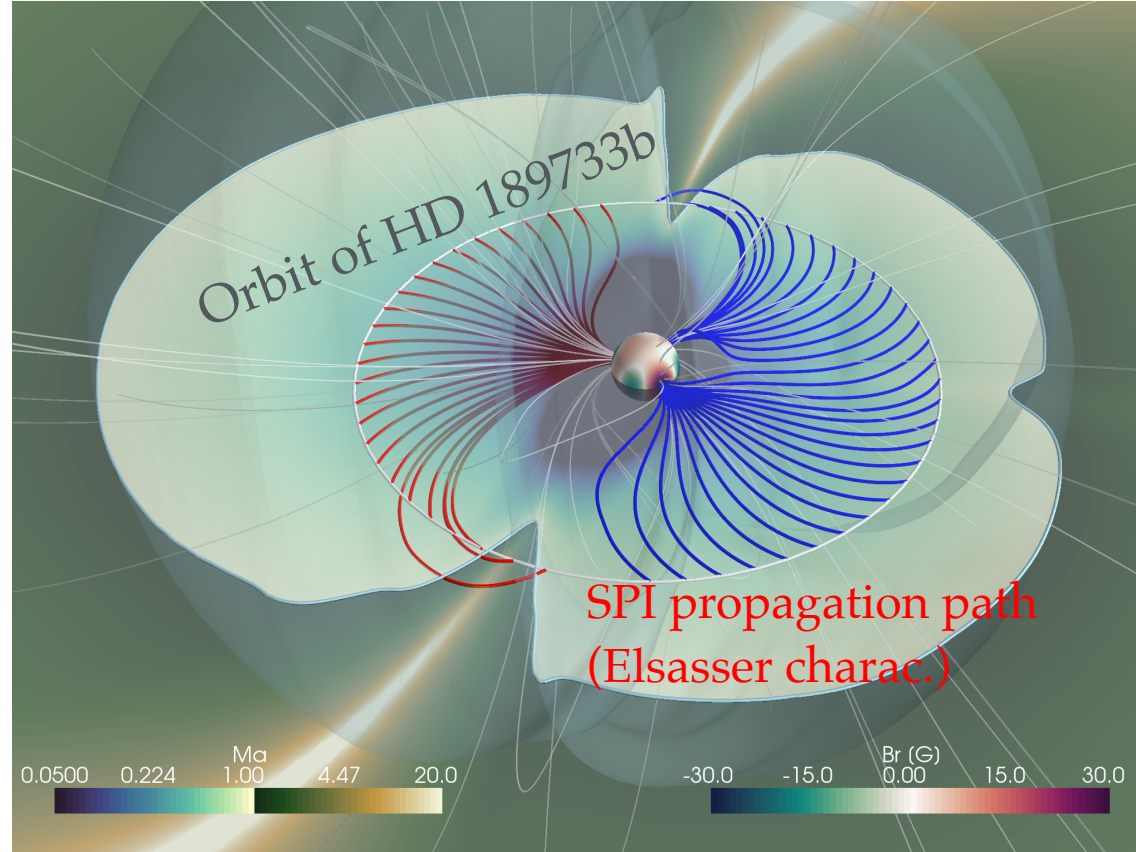
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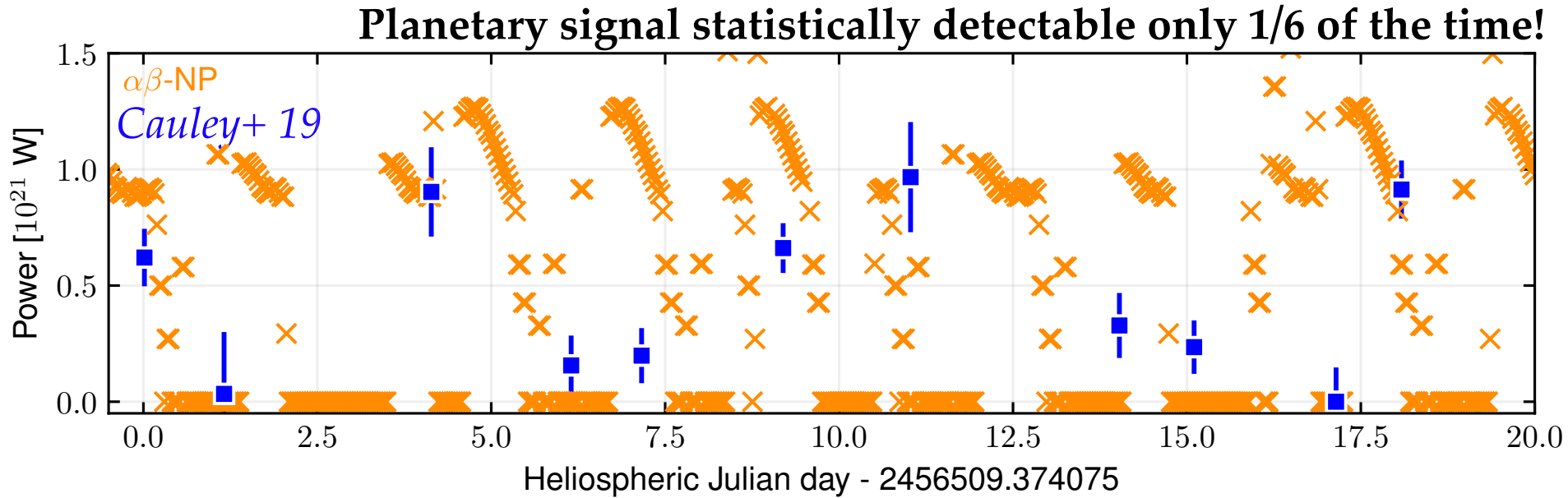
[Strugarek+ 22]



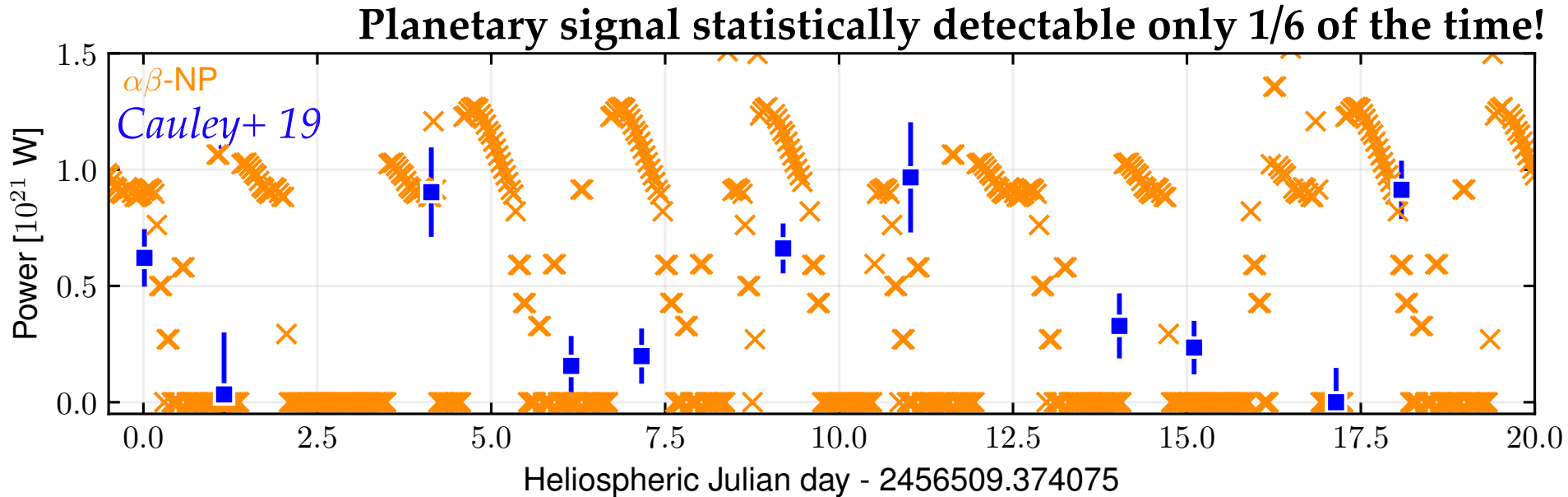
Simulation data available on Galactica database



# Detailed modelling of HD 189733

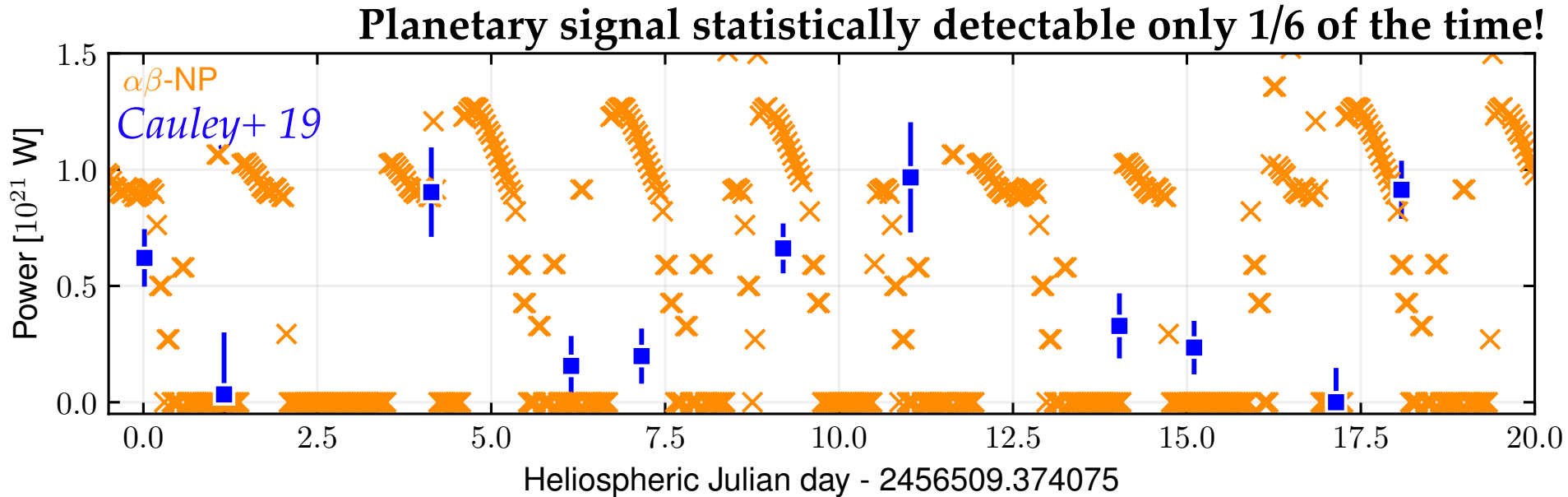


# Detailed modelling of HD 189733



Only the stretch-and-break mechanism provides enough power to explain the level of stellar chromospheric emission

# Detailed modelling of HD 189733



Only the stretch-and-break mechanism provides enough power to explain the level of stellar chromospheric emission

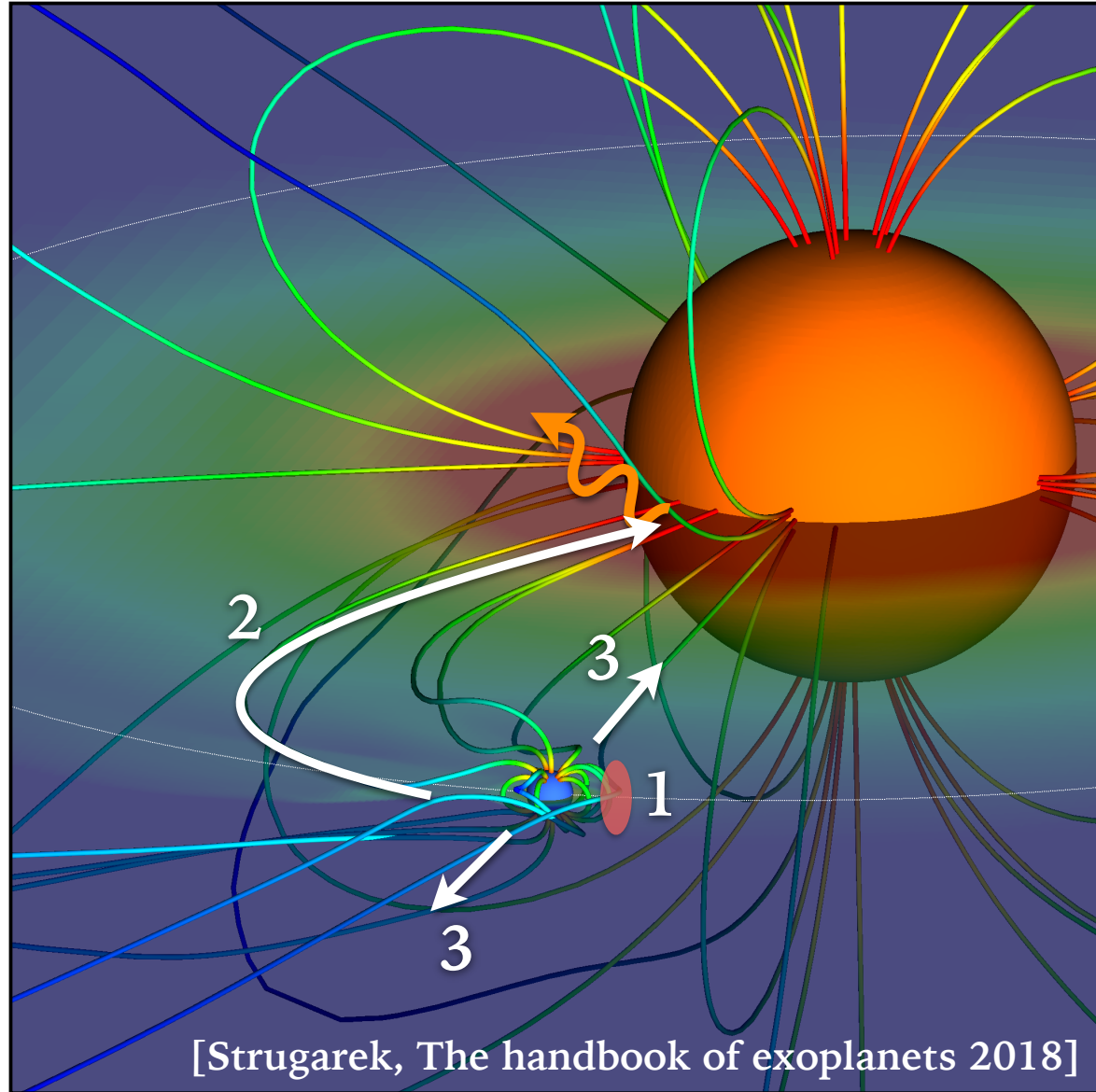
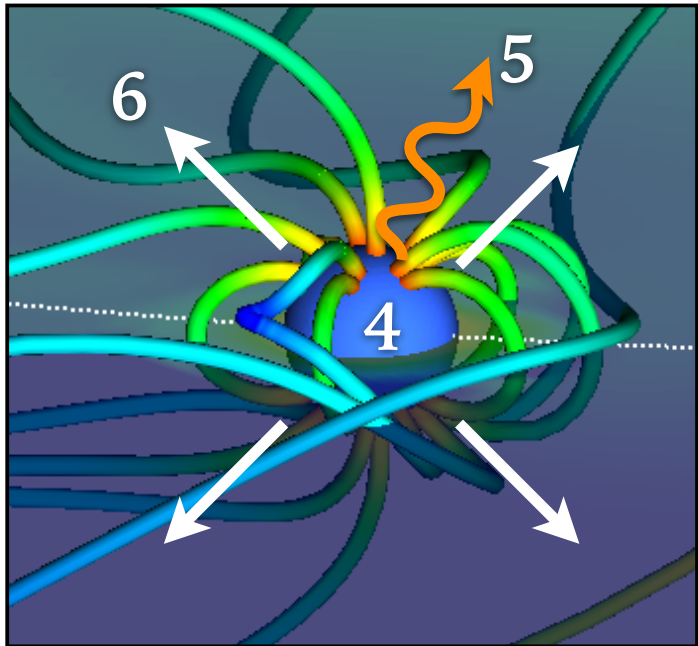
Denser observational campaigns are needed, and could lead to the confirmation of the star-planet magnetic interaction interpretation

# A few points to remember about star-planet magnetic interactions

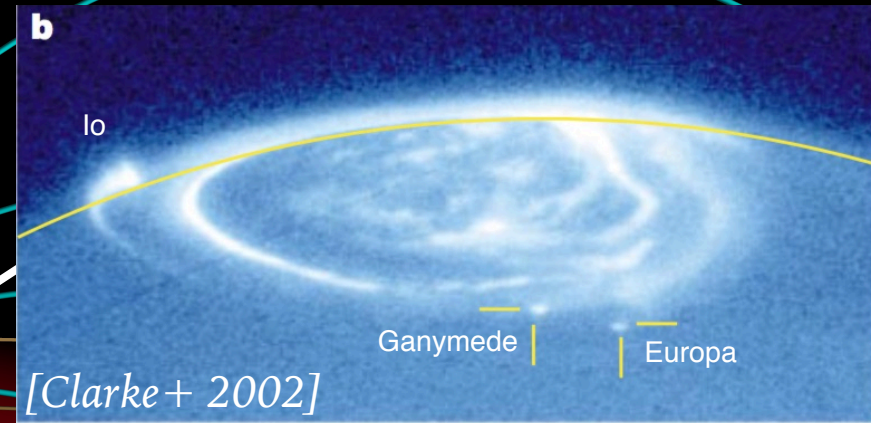
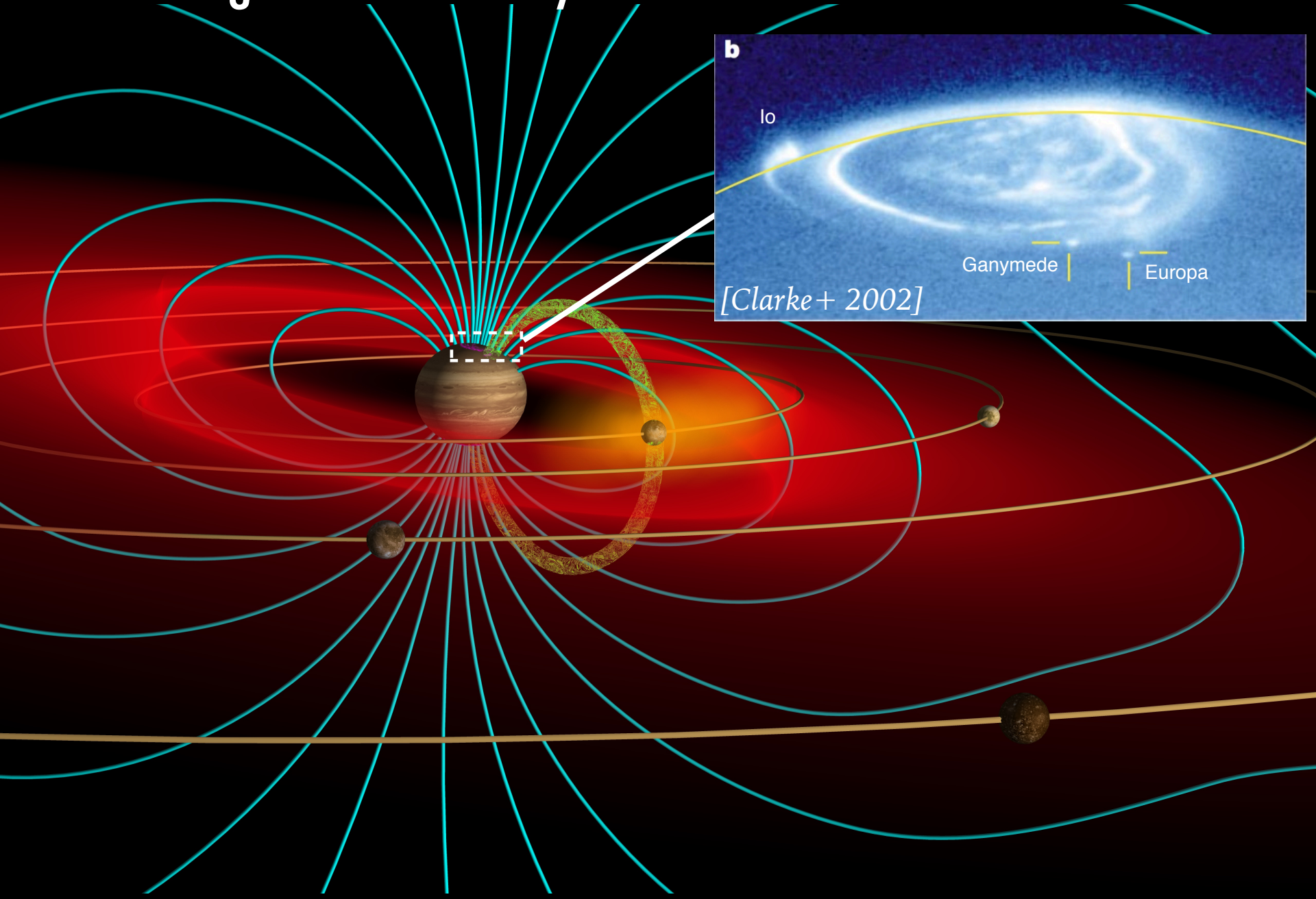
- » Any tracer can be **strongly dephased** from the orbital phase
- » It is in general not necessarily directly correlated by the orbital motion : its temporal variability originates from the convolution between the
  - » Orbital period (**observed**)
  - » Alfvén wave travel time (**requires a wind/corona model**)
  - » Stellar rotation period (**observed**)
- » It is generally variable on **short timescales** (inhomogeneities along the orbital path) and **long timescales** (stellar magnetic cycle/variability)
- » It is **systematically awaited** for planets close to their host

# Major effects of star-planet magnetic interactions

1. Shocks & geometry
2. Energy channeling
3. Planet migration
4. Planet heating
5. Planet emissions
6. Atmospheric escape



# Alfvén wings in the solar system : the Jovian case



# Energetics of magnetic interaction

Kinetic energy flux

$$P_c = \left( \frac{1}{2} \rho v^2 \right) v A_{\text{obst}}$$

Poynting flux intercepted by obstacle

$$P_m = \frac{B_{\perp}^2}{2\mu_0} v A_{\text{obst}}$$

# Energetics of magnetic interaction

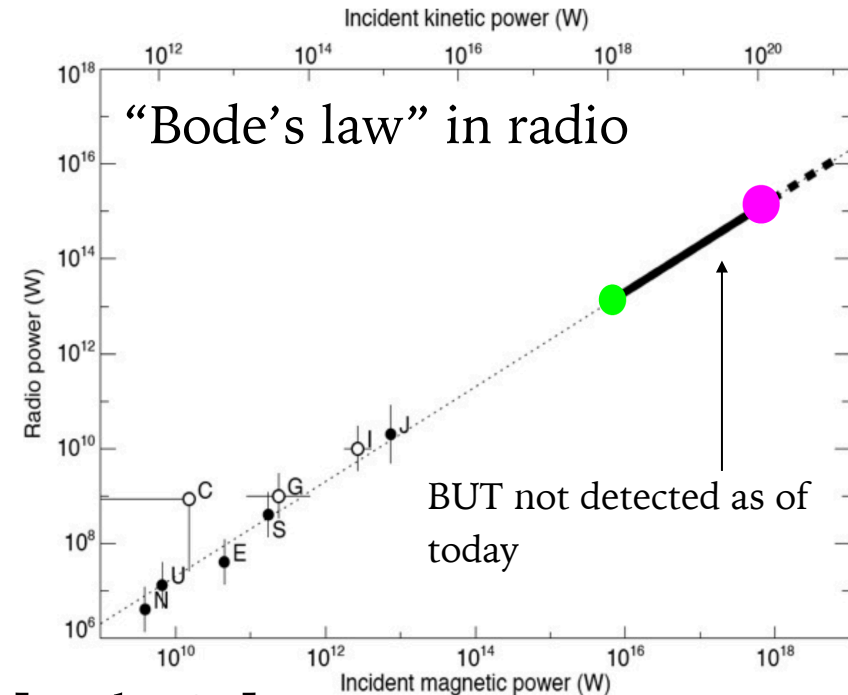
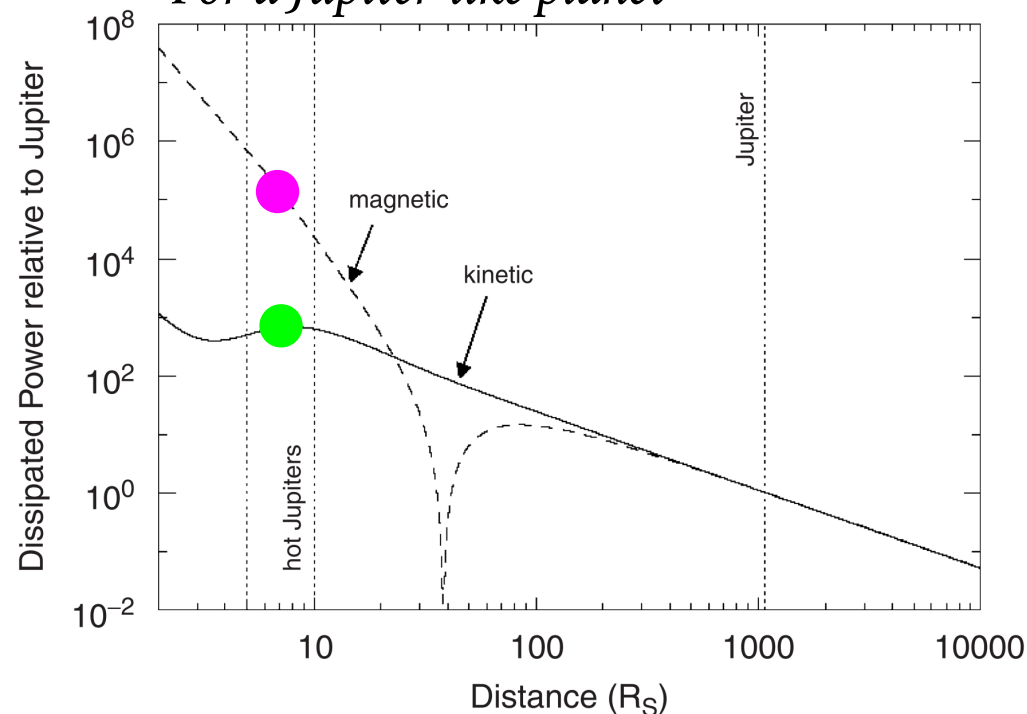
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*For a Jupiter-like planet*



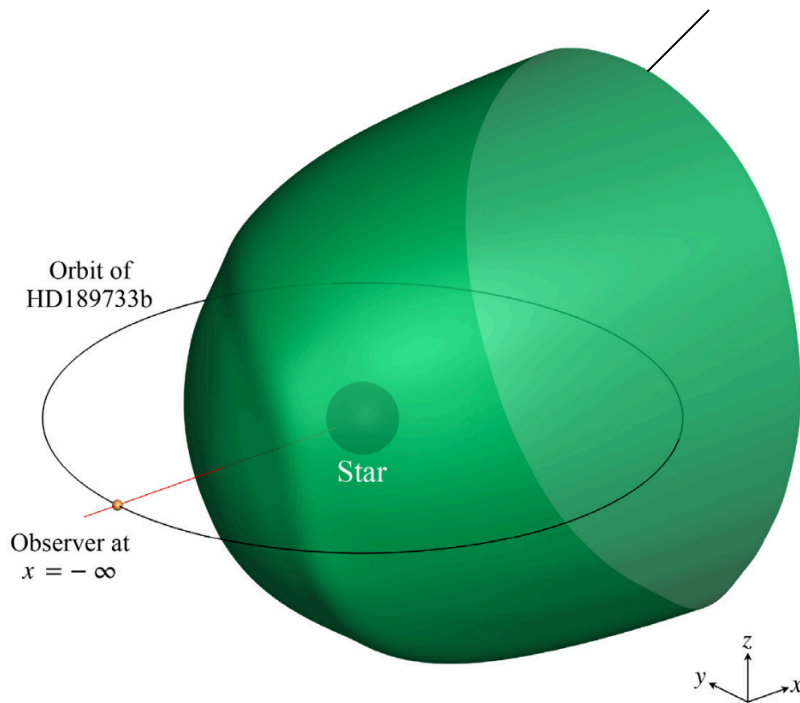
[Zarka 07]



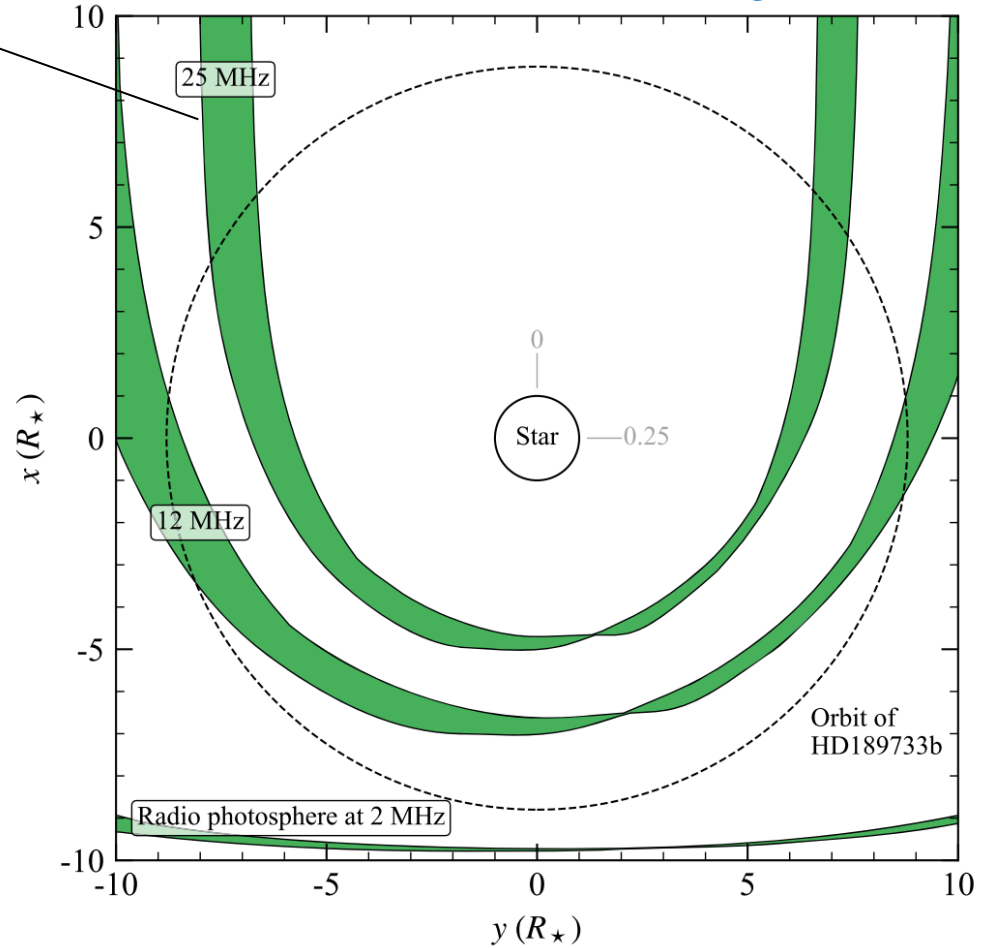
# Why is it so hard to detect radio emissions from exoplanets?

Radio 'photosphere' @ 25 MHz

[Kavanagh + 2019]



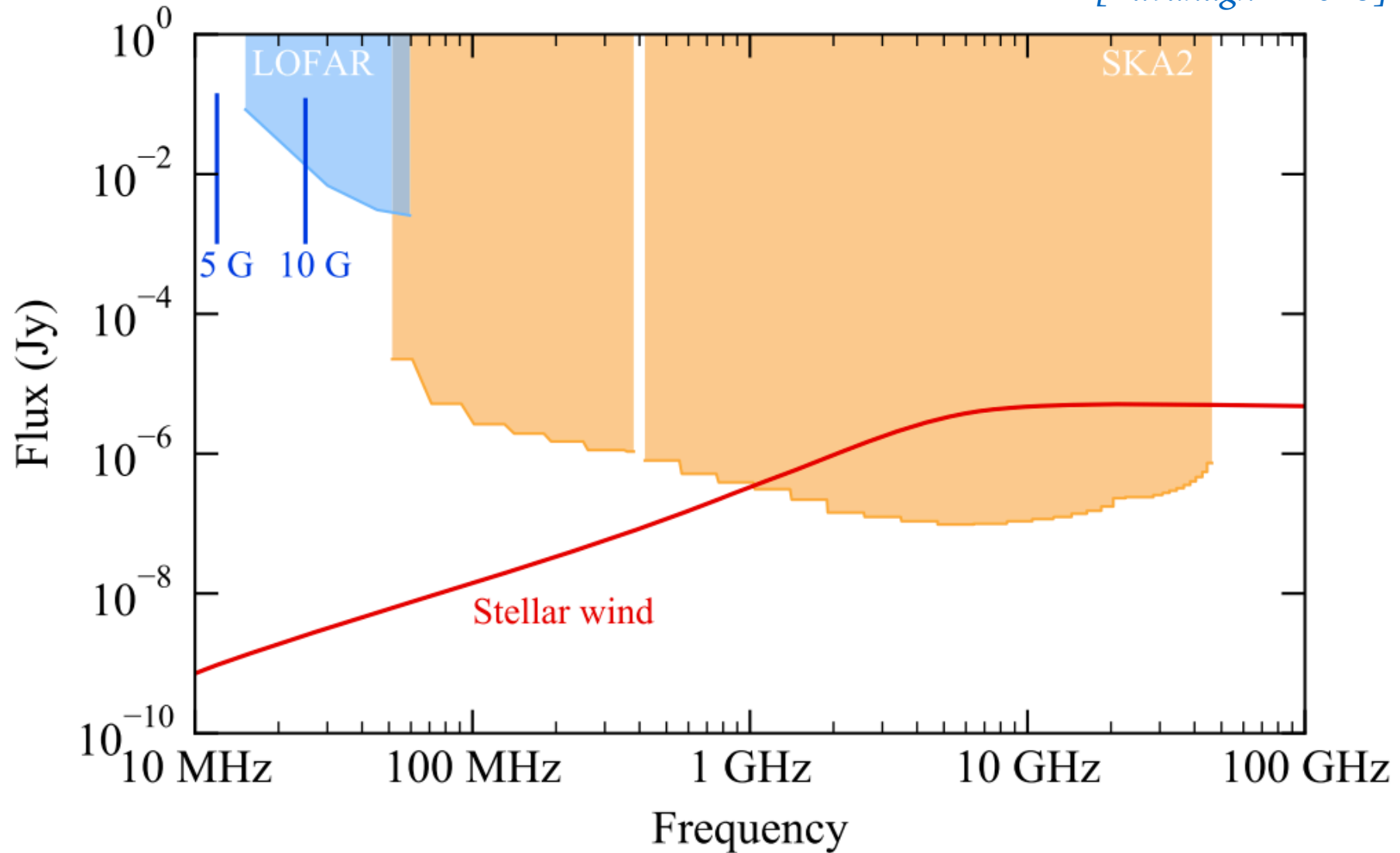
HD 189733



[see also Nichols & Milan 2016]

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# Energetics of magnetic interaction

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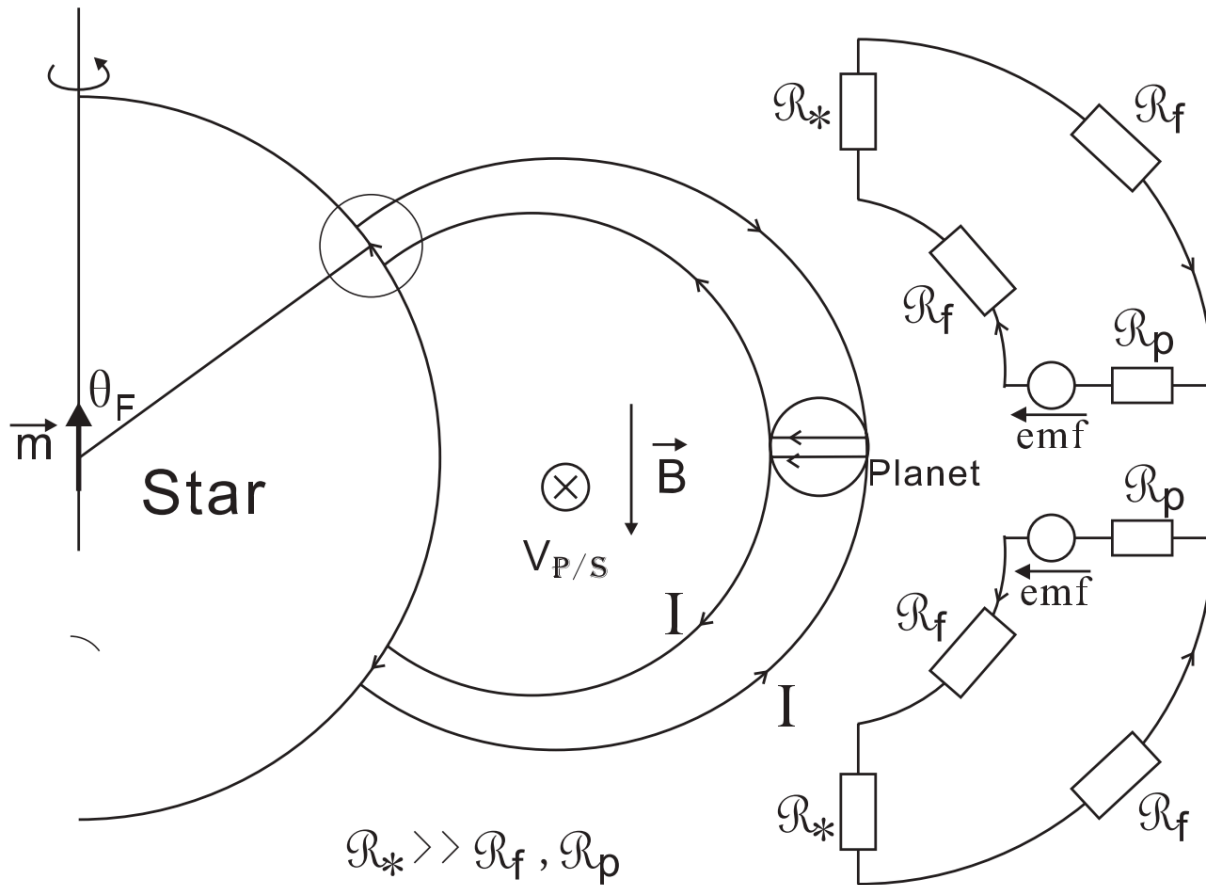
Poynting flux intercepted by obstacle

$$P_m = \frac{B_{\perp}^2}{2\mu_0} v A_{\text{obst}}$$

**Several non-trivial questions need now to be answered:**

- » What are the shape and size ( $A_{\text{obst}}$ ) of the obstacle?
- » Which fraction of the available power is channeled by the interaction ?
- » How much energy is actually available in the end?

# Obstacle & energy conversion: the unipolar inductor case



## Hypotheses:

- » Negligible planet magnetic field
- » Resistance  $\Sigma$  at the foot of the interaction is **larger** than in the planet (i.e. conductivity is the largest in the planet)

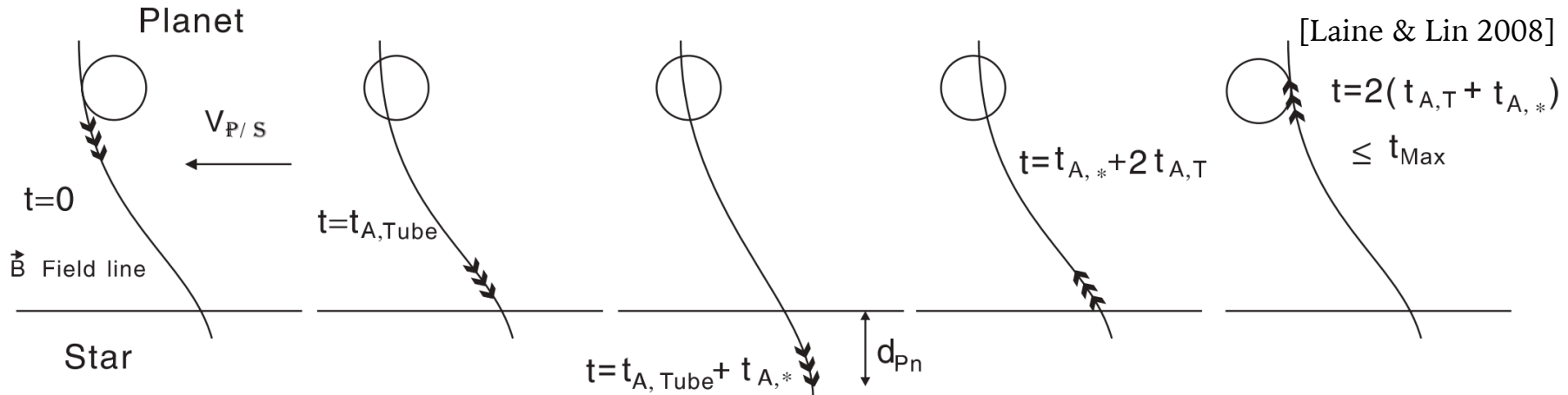
Note: one can derive the net torque applied to the planet's orbit

$$\mathcal{T}_* = 8R_P^2 R_{orb}^2 \sigma B_w^2 \Sigma$$

[Laine & Lin 2008]

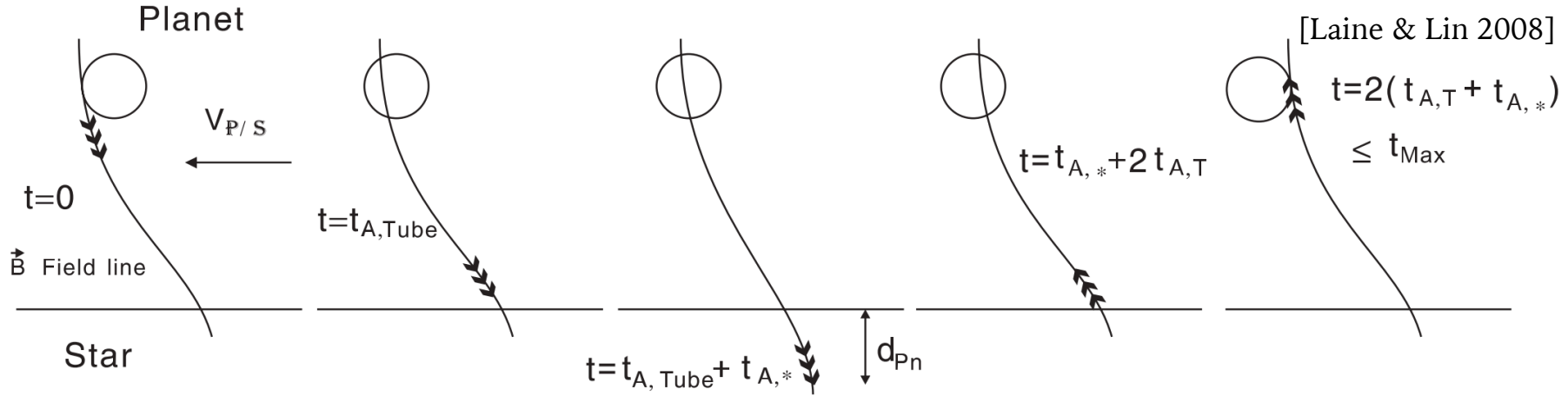
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The unipolar inductor model is valid only if the current systems is closed between the star and the planet:

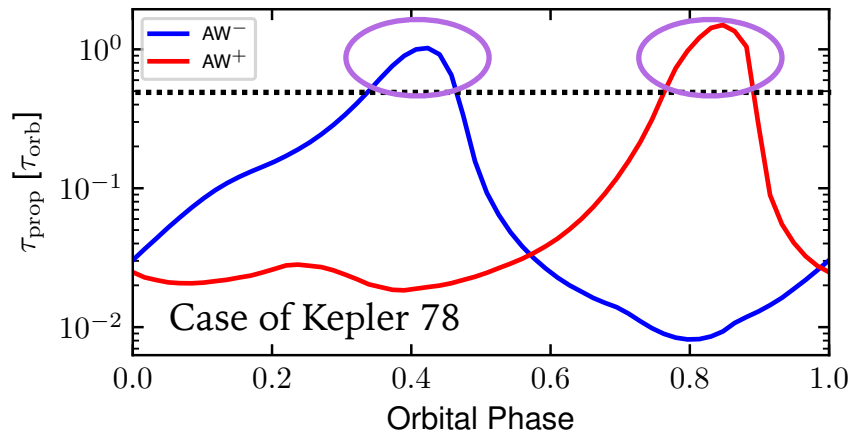


# Obstacle & energy conversion: the unipolar inductor case

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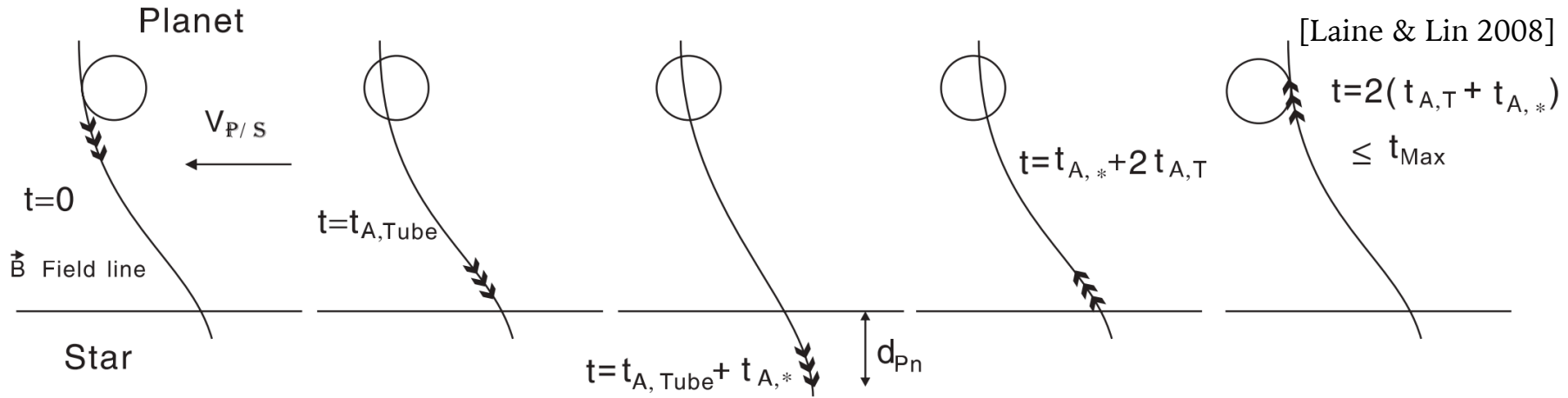


Propagation time from planet to star

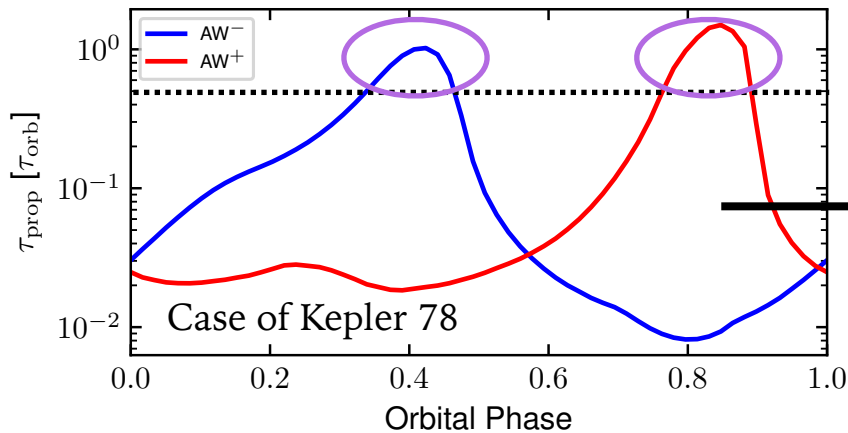


# Obstacle & energy conversion: the unipolar inductor case

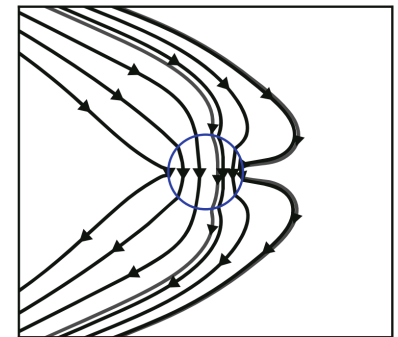
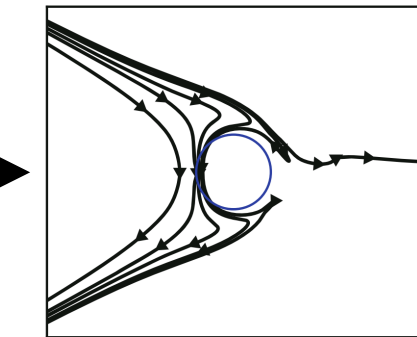
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Propagation time from planet to star

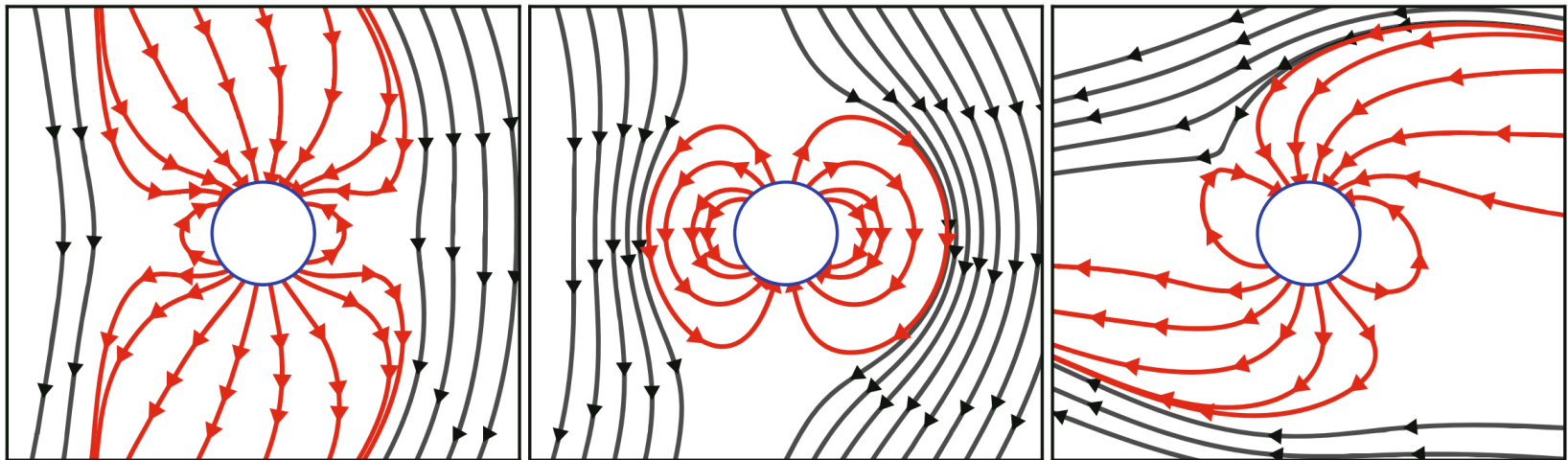


Two possible situations:



# Obstacle & energy conversion: the dipolar case

What if the planetary field cannot be neglected?



Aligned

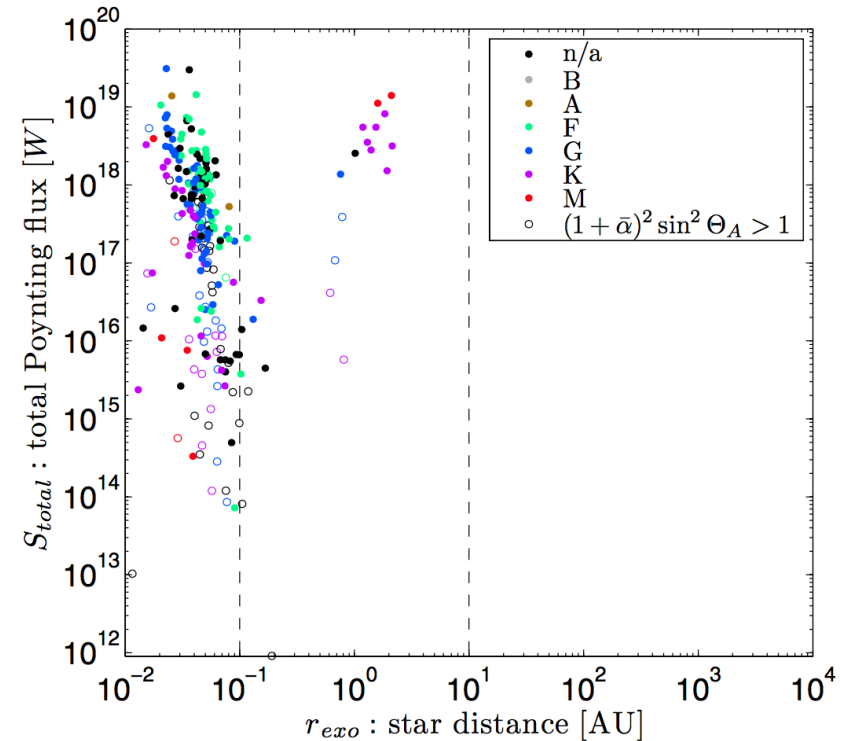
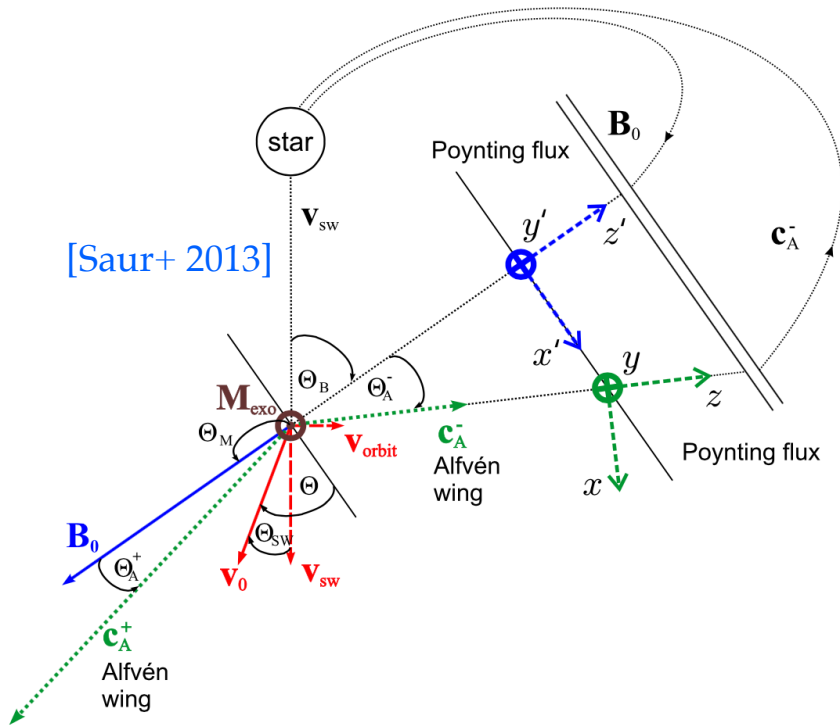
Anti-aligned

Perpendicular





# Obstacle & energy conversion: the dipolar case



The power channeled by the interaction is 
$$P = \left[ M_A \left( \frac{\Sigma_P}{\Sigma_P + 2\Sigma_A} \right)^2 \right] S_w A_{\text{obst}}$$

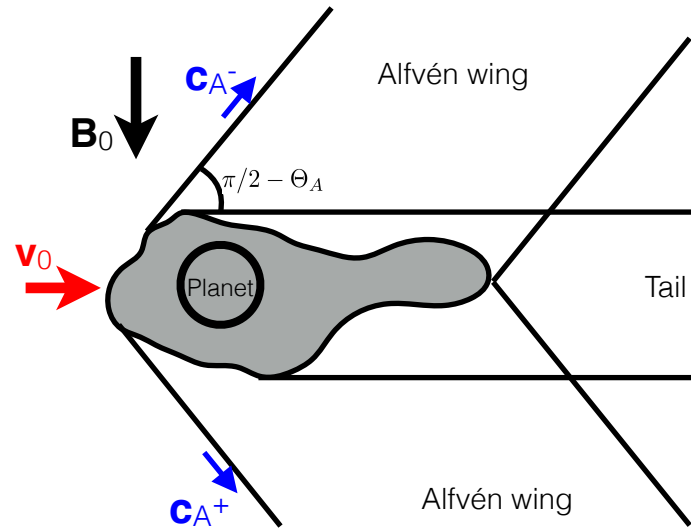
'Alfvén conductance'

$$\Sigma_A = \frac{1}{\mu_0 v_A (1 + M_A^2 - 2M_A \cos \Theta)^{1/2}}$$

$\Sigma_P$ : Pedersen conductance  
(integrated ionospheric conductivity  
 $\perp$  to  $E_{\text{iono}}$  and  $\parallel$  to  $B_{\text{iono}}$ )

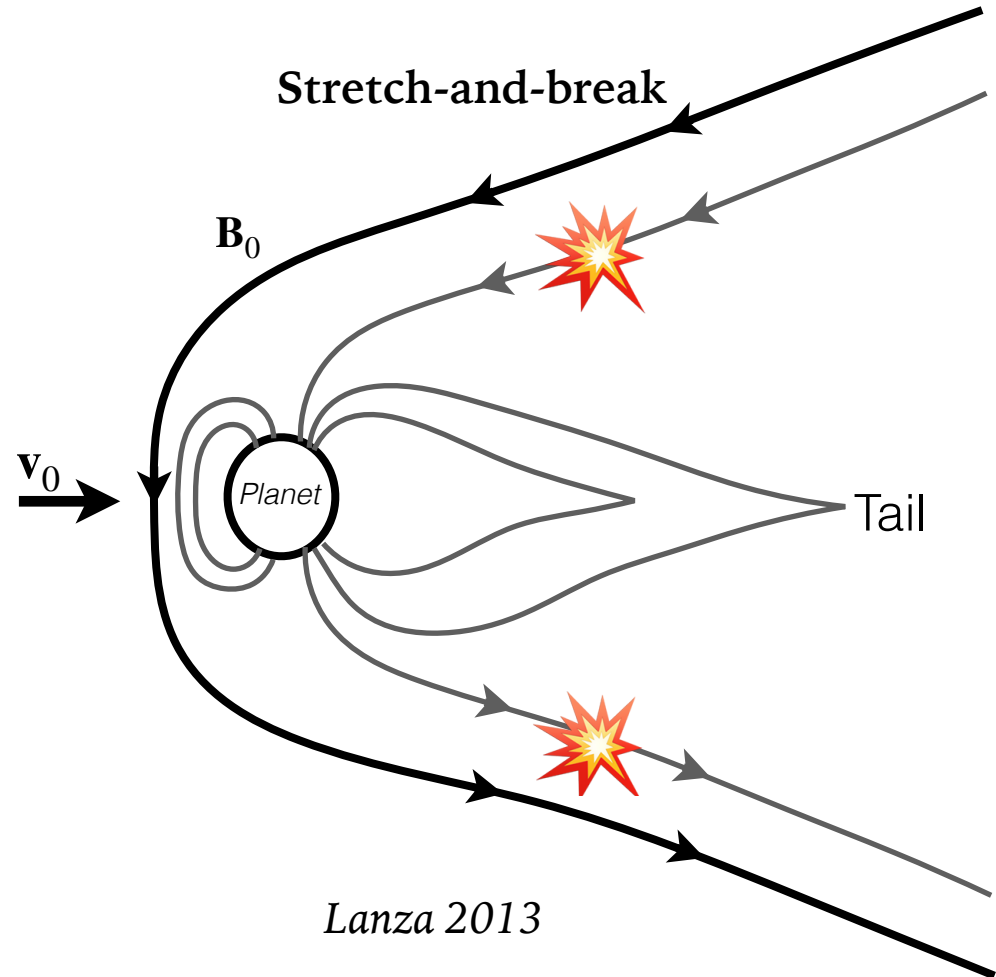
# Alternative theory: the stretch-and-break mechanism

Alfvén wings



*Saur et al. 2013*

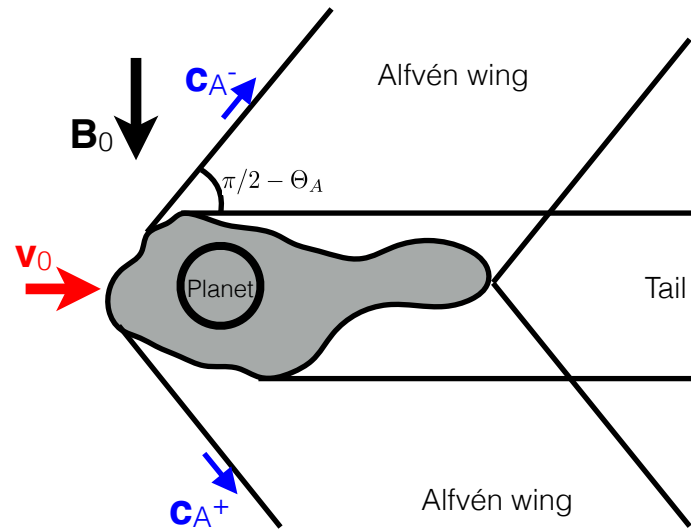
Stretch-and-break



*Lanza 2013*

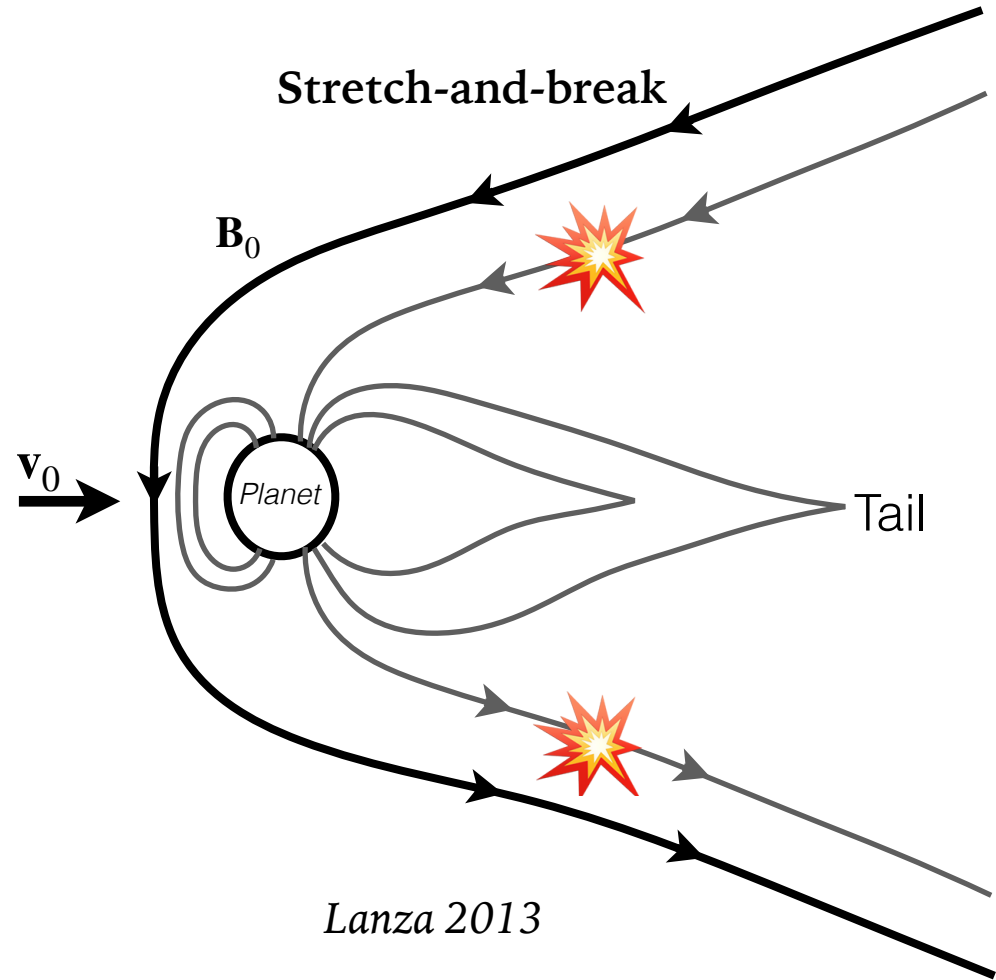
# Alternative theory: the stretch-and-break mechanism

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Stretch-and-break



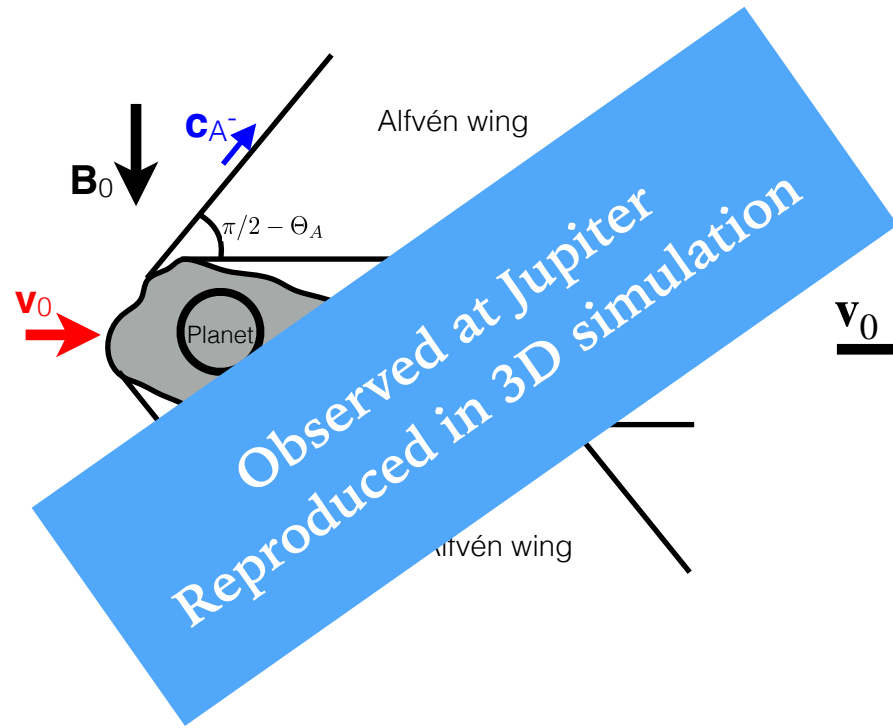
Lanza 2013

$$\text{Poynting flux} \propto v_0 B_w^2 M_a \left( \frac{B_P}{B_w} \right)^{2/3} \ll \text{Poynting flux} \propto v_0 B_w^2 \left( \frac{B_P}{B_w} \right)^2 f_{AP}$$

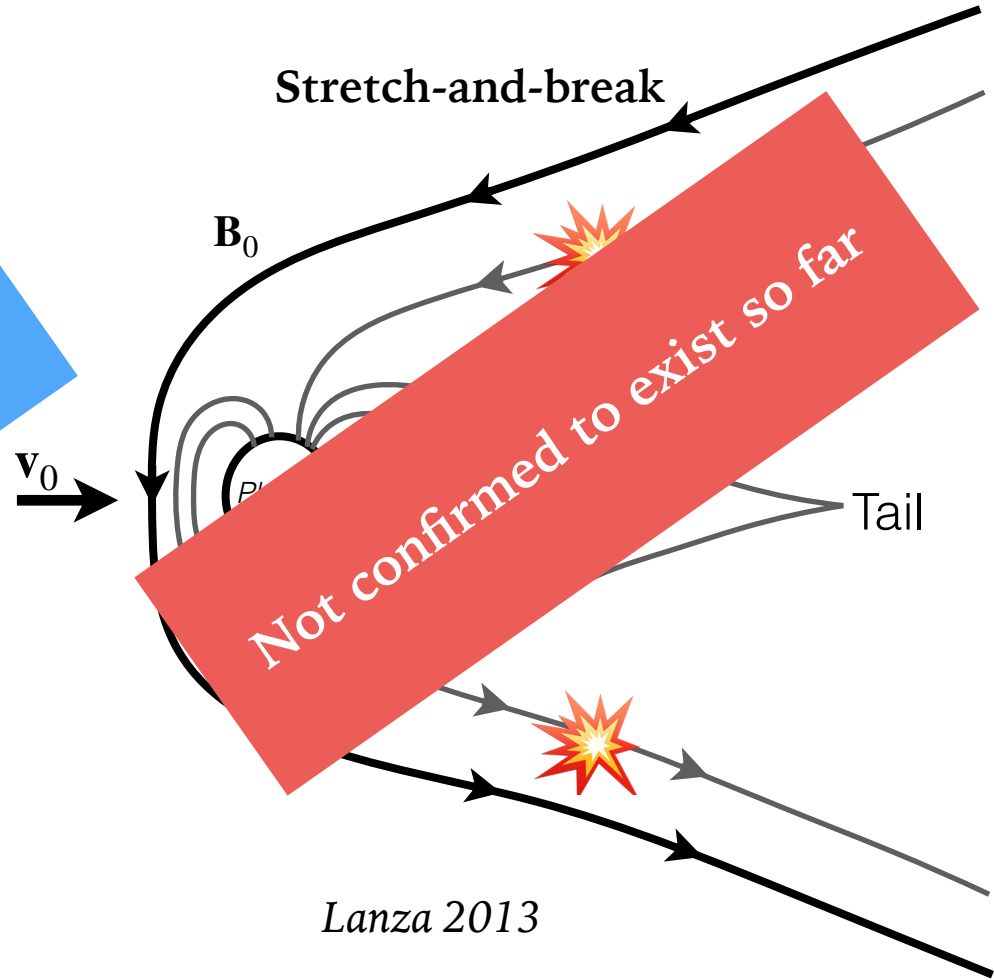
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Lanza 2013

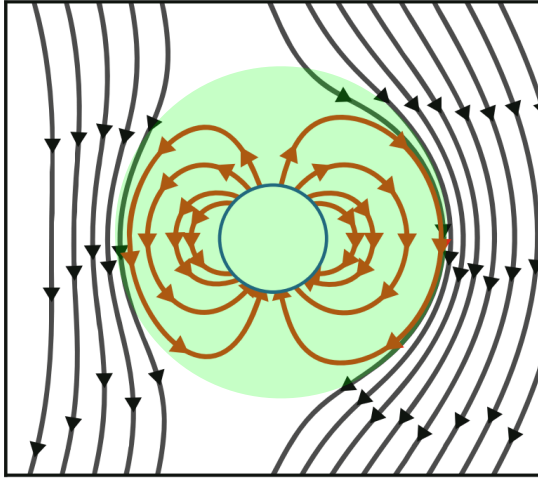
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$\ll$

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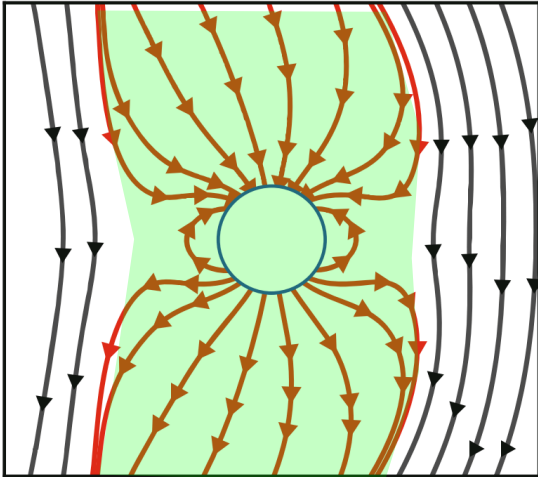
# The obstacle size strongly depends on the topology

Anti-aligned



$$R_{\text{obst}} = R_P \left( \frac{B_P^2}{8\pi P_t} \right)^{1/6}$$

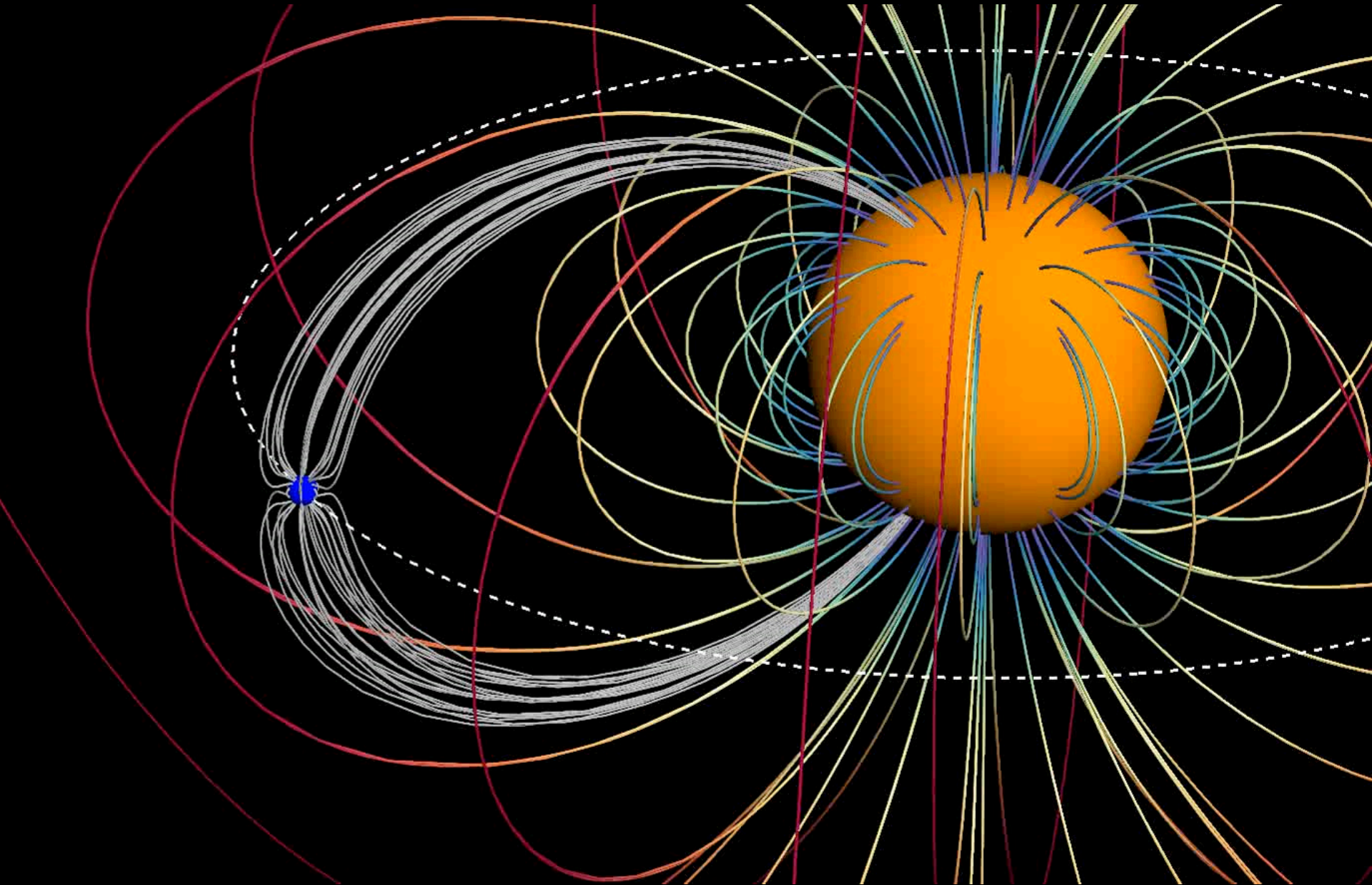
Aligned



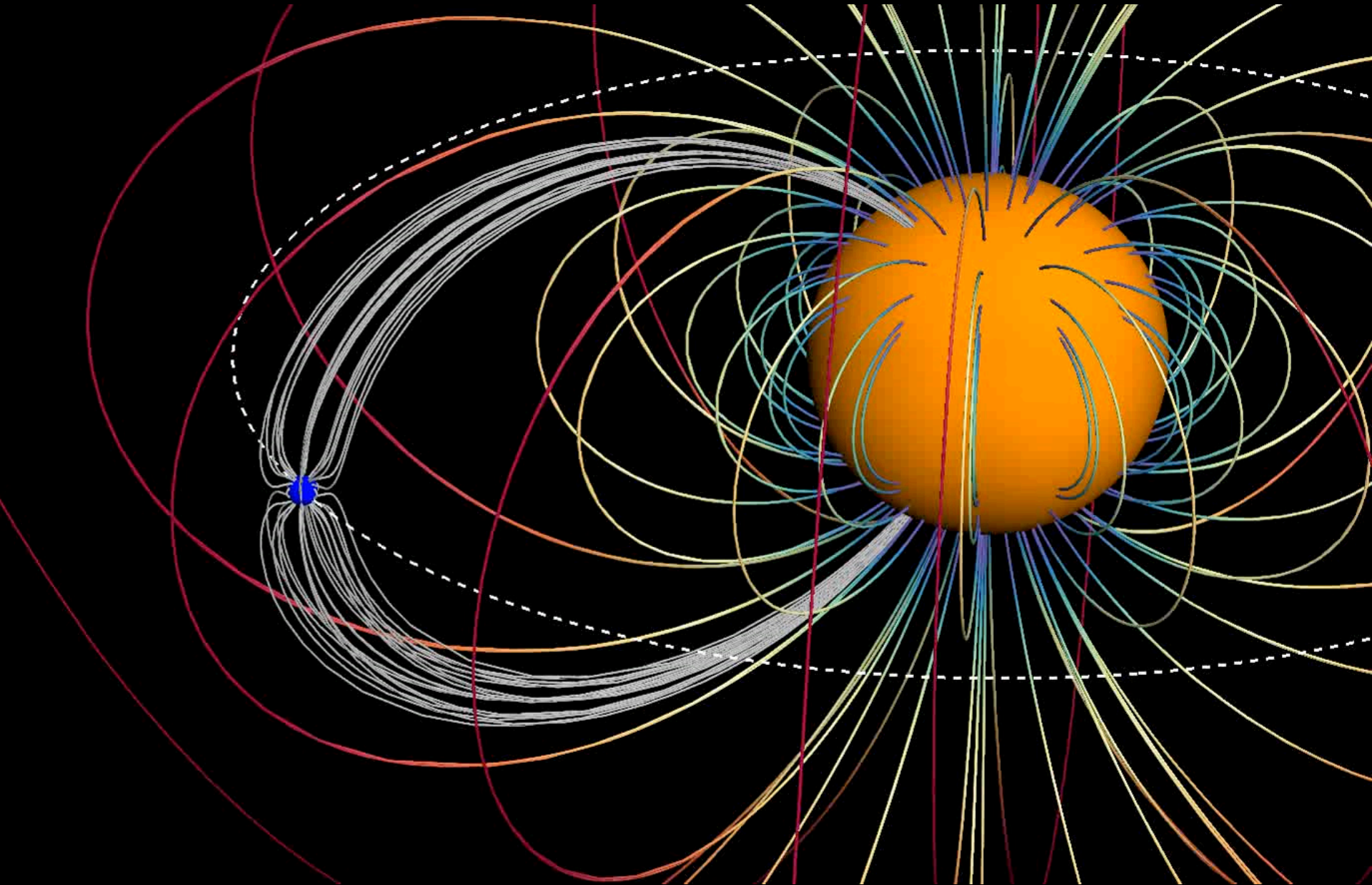
The obstacle can be much bigger, so much more power is accessible to the interaction

▀ possible to constrain this with numerical simulations

# 3D model of magnetic star-planet interactions



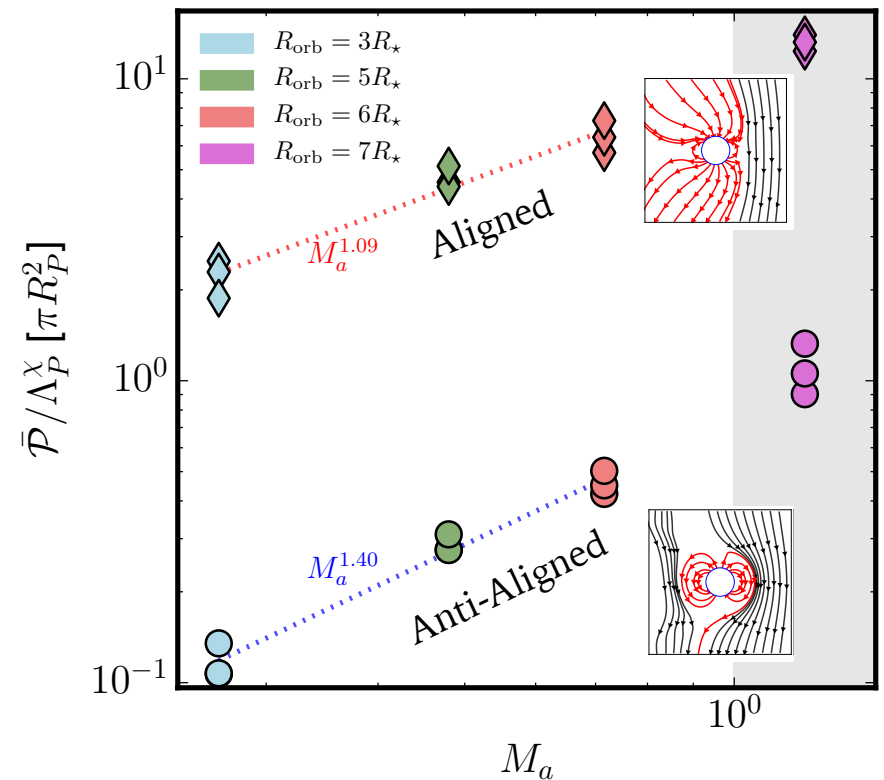
# 3D model of magnetic star-planet interactions





# Poynting flux in Alfvén wings: parametrization

$$\mathcal{P} = A_1 \pi \left( c_d S_w M_a^\xi \bar{\eta}_a^{\nu_3} \right) \cdot \left( R_P^2 \bar{\eta}_P^{\nu_4} \right) \cdot \left( \Lambda_P^\chi \right)$$

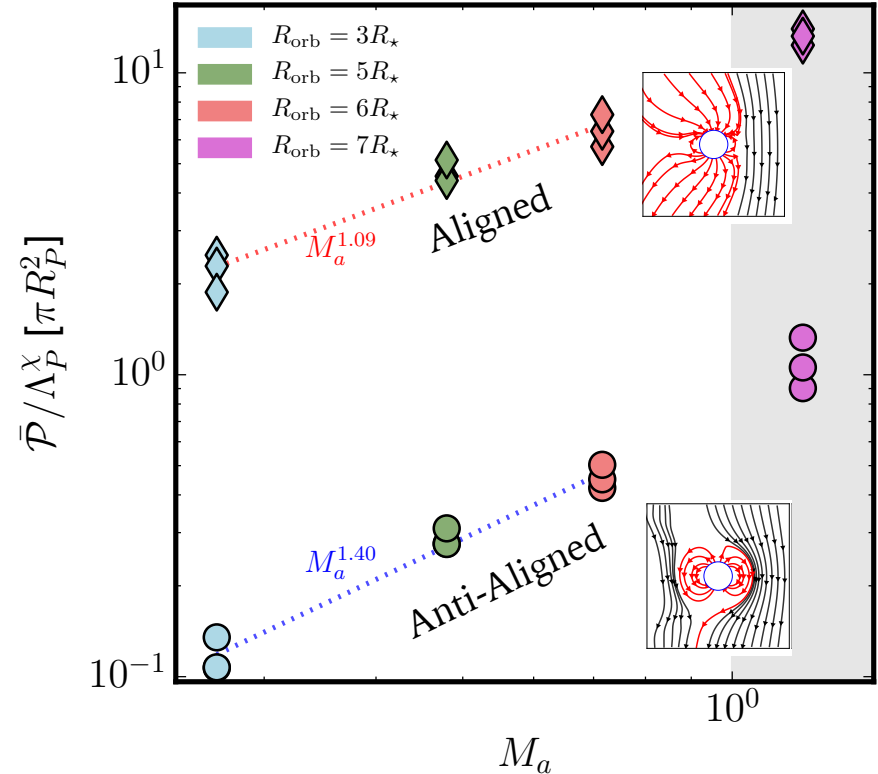


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$$\mathcal{P} = \underbrace{A_1 \pi}_{\text{drag coeff.}} \underbrace{\left( c_d S_w M_a^\xi \bar{\eta}_a^{\nu 3} \right)}_{\text{Alf. Mach}} \cdot \left( R_P^2 \bar{\eta}_P^{\nu 4} \right) \cdot \left( \Lambda_P^\chi \right)$$

*Poynting flux*
*Reconnection eff.*

Wind-dependent parameters  
(magnetic field, density, velocity)

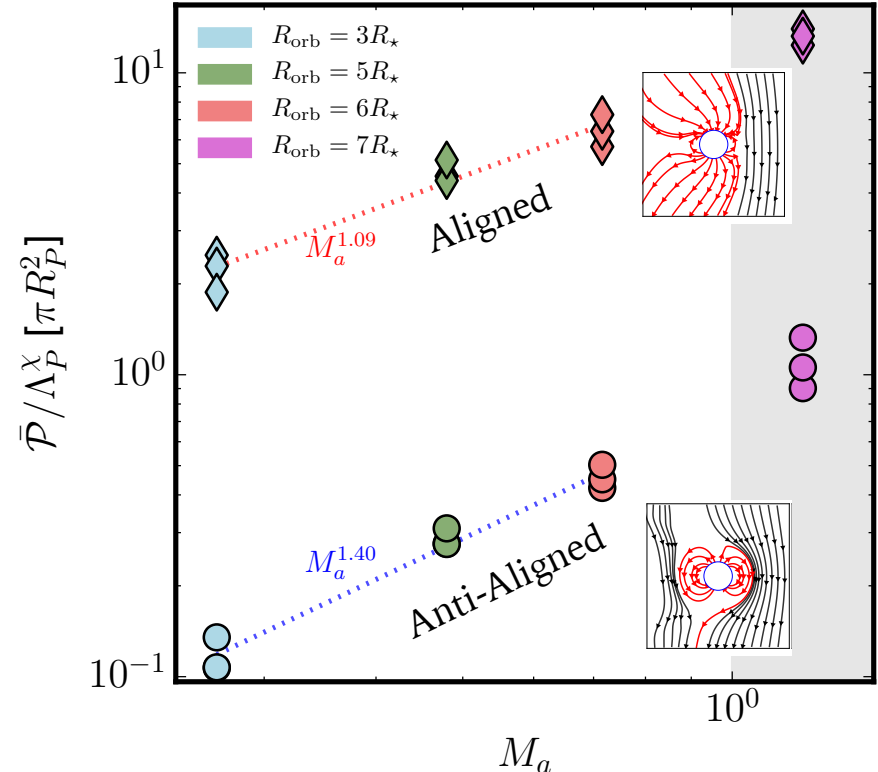


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Planet-dependent parameters  
(radius, composition)



# Poynting flux in Alfvén wings: parametrization

Poynting flux

Reconnection eff.

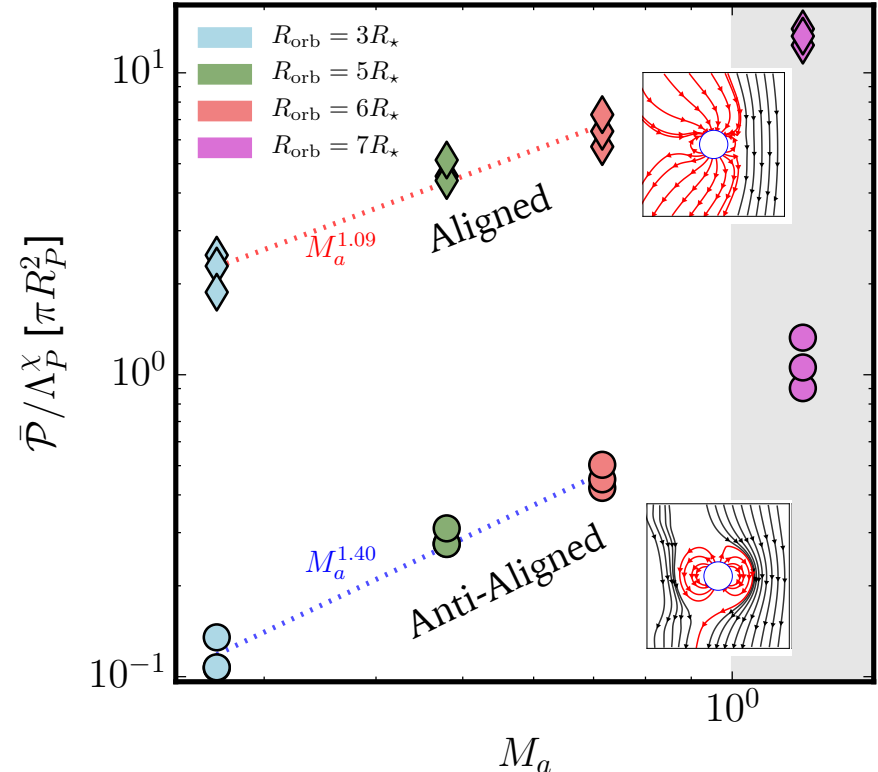
Pedersen diff.

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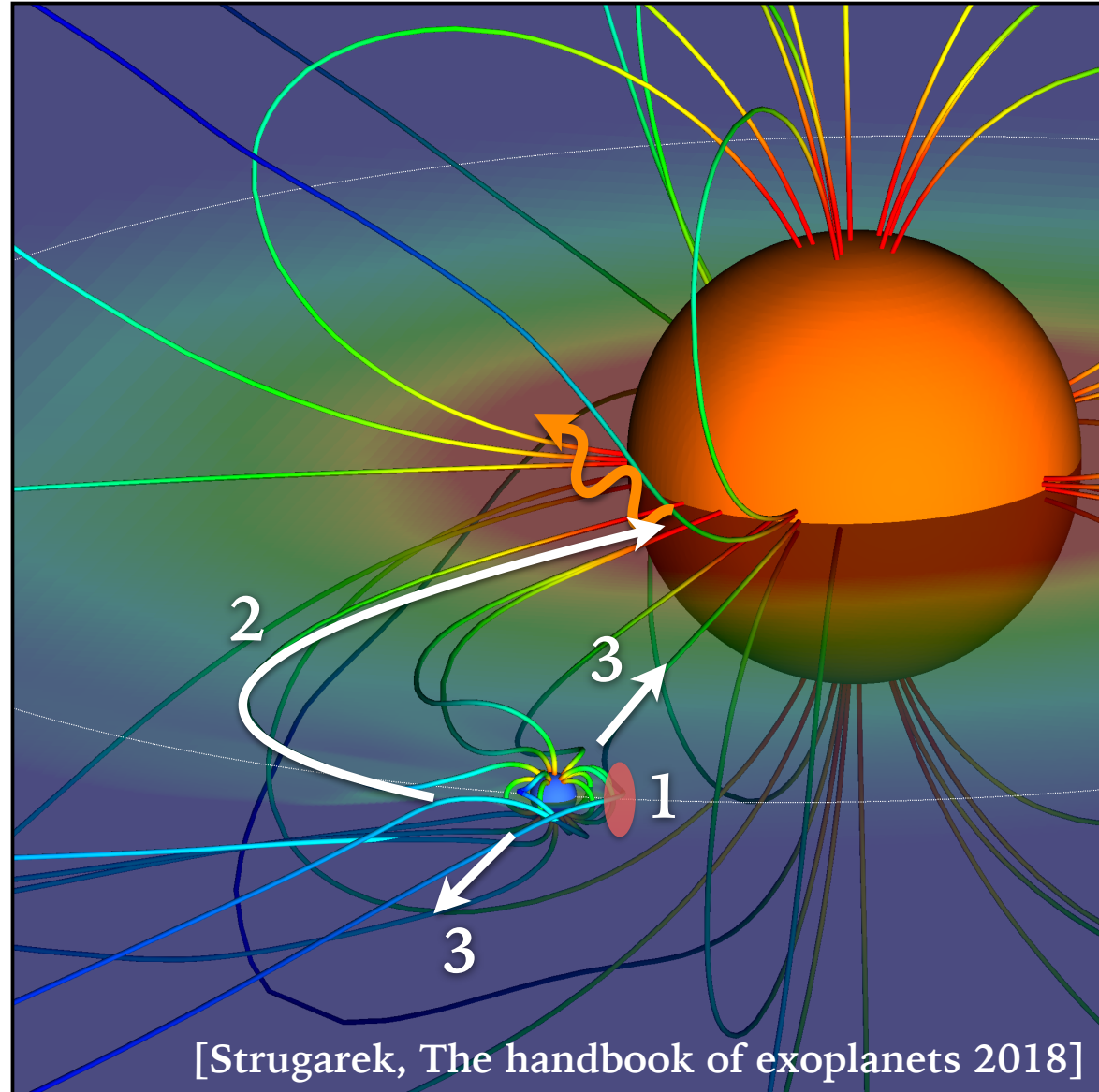
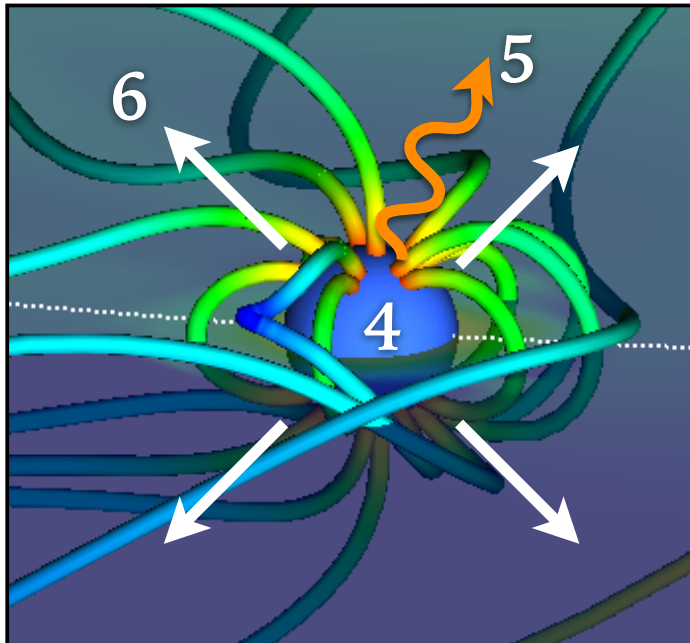
Planet-dependent parameters  
(radius, composition)

Magnetosphere/wind pressure ratio  
(**planetary magnetic field**, wind)



# Major effects of star-planet magnetic interactions

1. Shocks & geometry
2. Energy channeling
- 3. Planet migration**
4. Planet heating
5. Planet emissions
6. Atmospheric escape



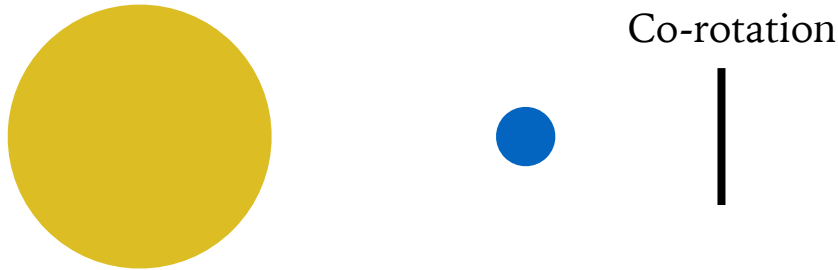
# Planetary migration due to a net torque

The total angular momentum is conserved: a planet migration corresponds to an angular momentum exchange between the orbit of the planet and stellar rotation (in a two-body problem), i.e.  $\dot{J}_\star = -\dot{J}_P$

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$$J_P^{\text{init}} > 3J_\star^{\text{init}}$$

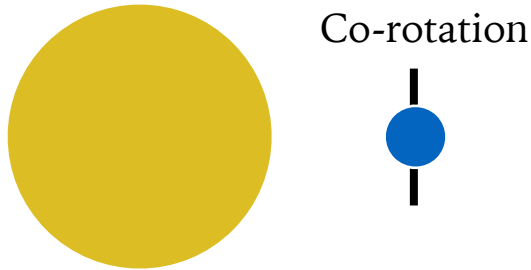


The system evolves towards  
synchronisation

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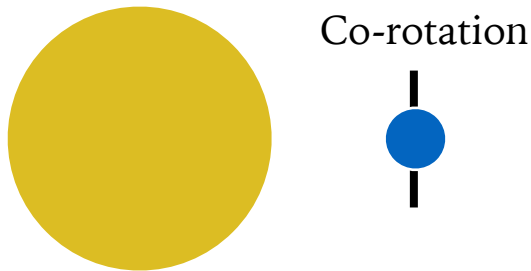
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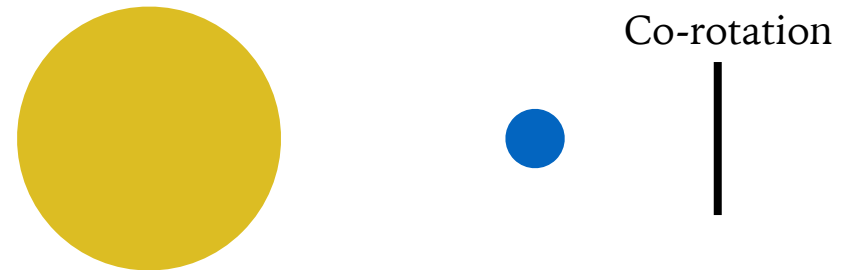
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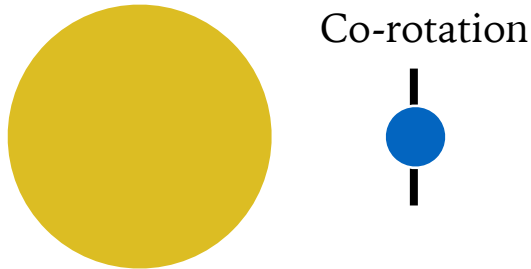


The system evolves until  
the planet is destroyed

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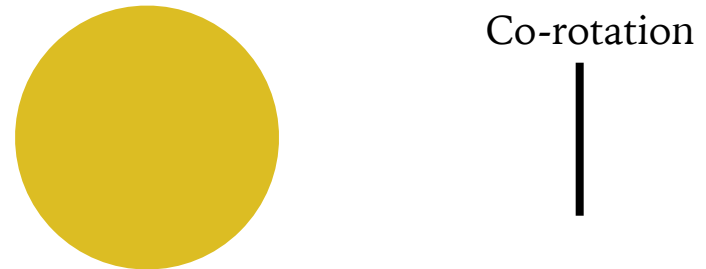
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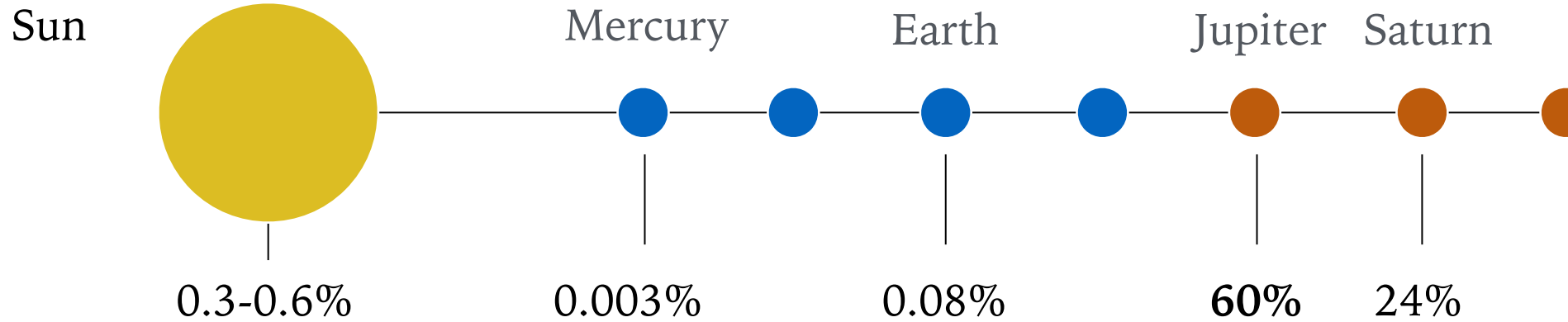
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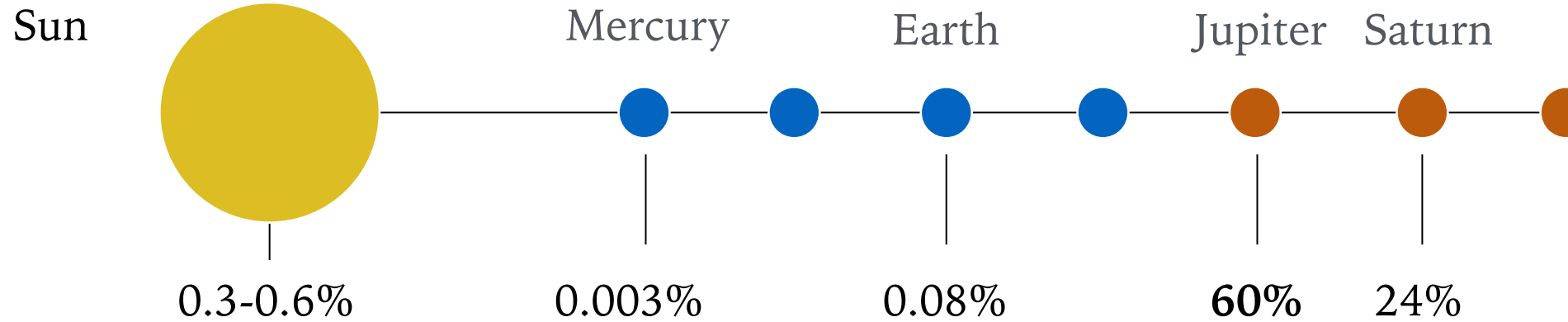
# Angular momentum in orbital systems



$$J_{\star} \simeq \beta M_{\star} R_{\star}^2 \Omega_{\star}$$

$$J_P \simeq M_P \sqrt{GM_{\star} R_{\text{orb}}}$$

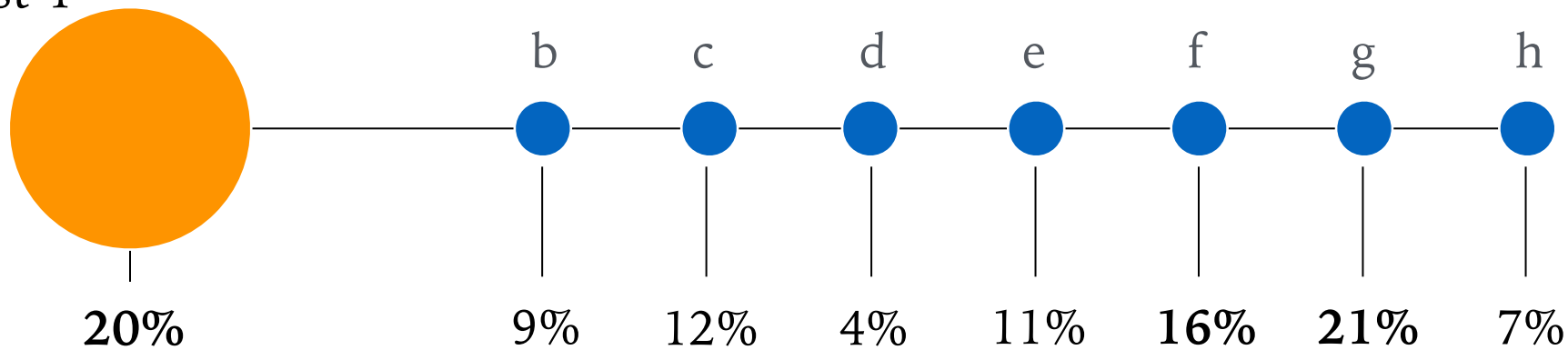
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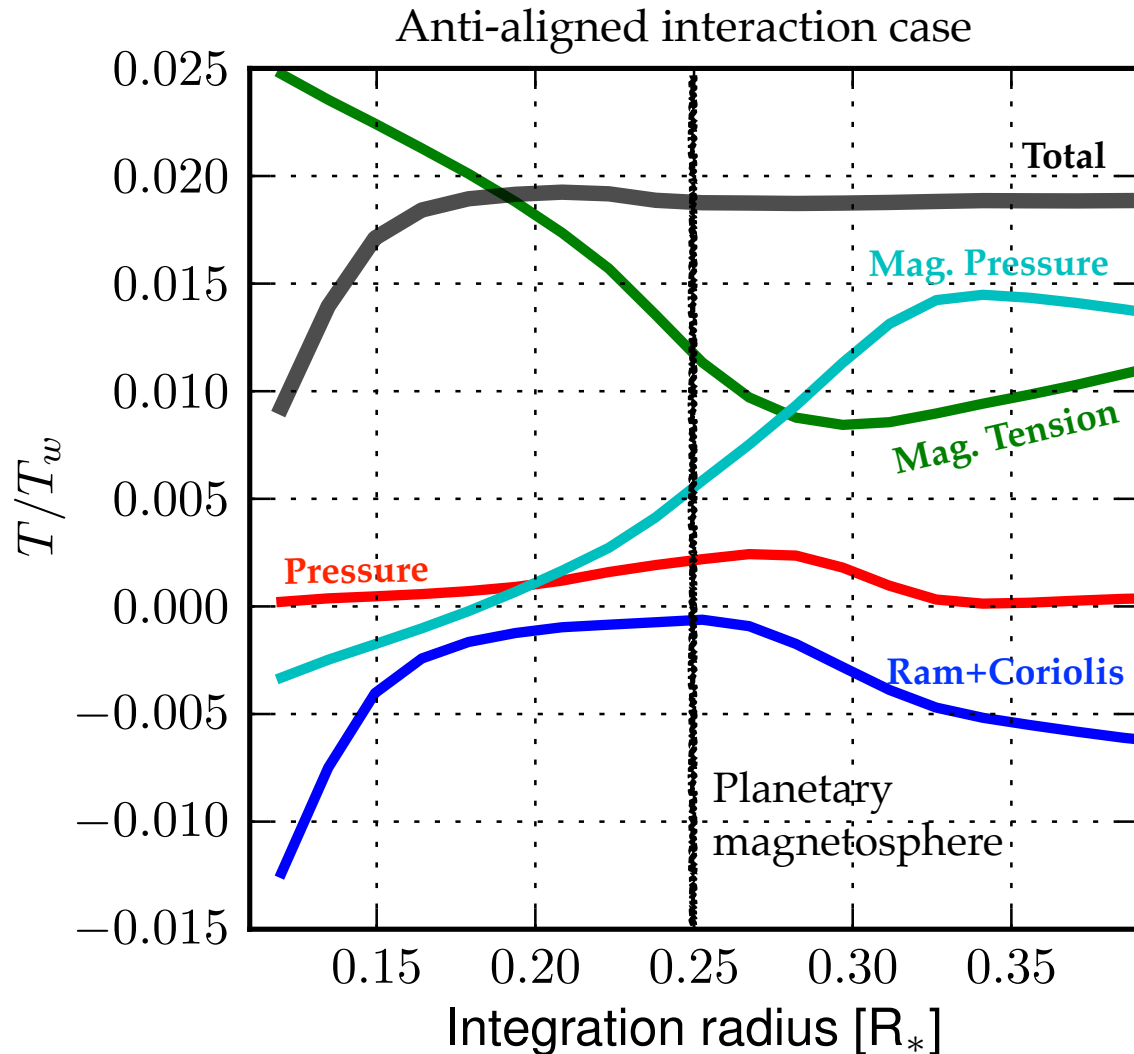
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Trappist-1

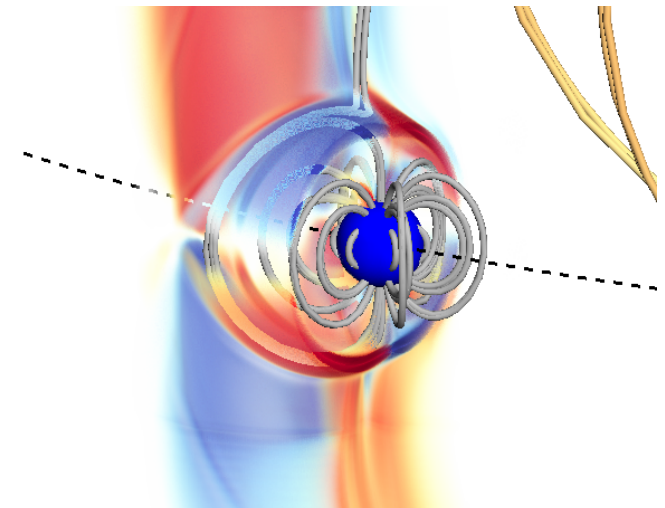


**Can magnetic interaction lead to a net torque?**

# Magnetic torque applied to close-in planets

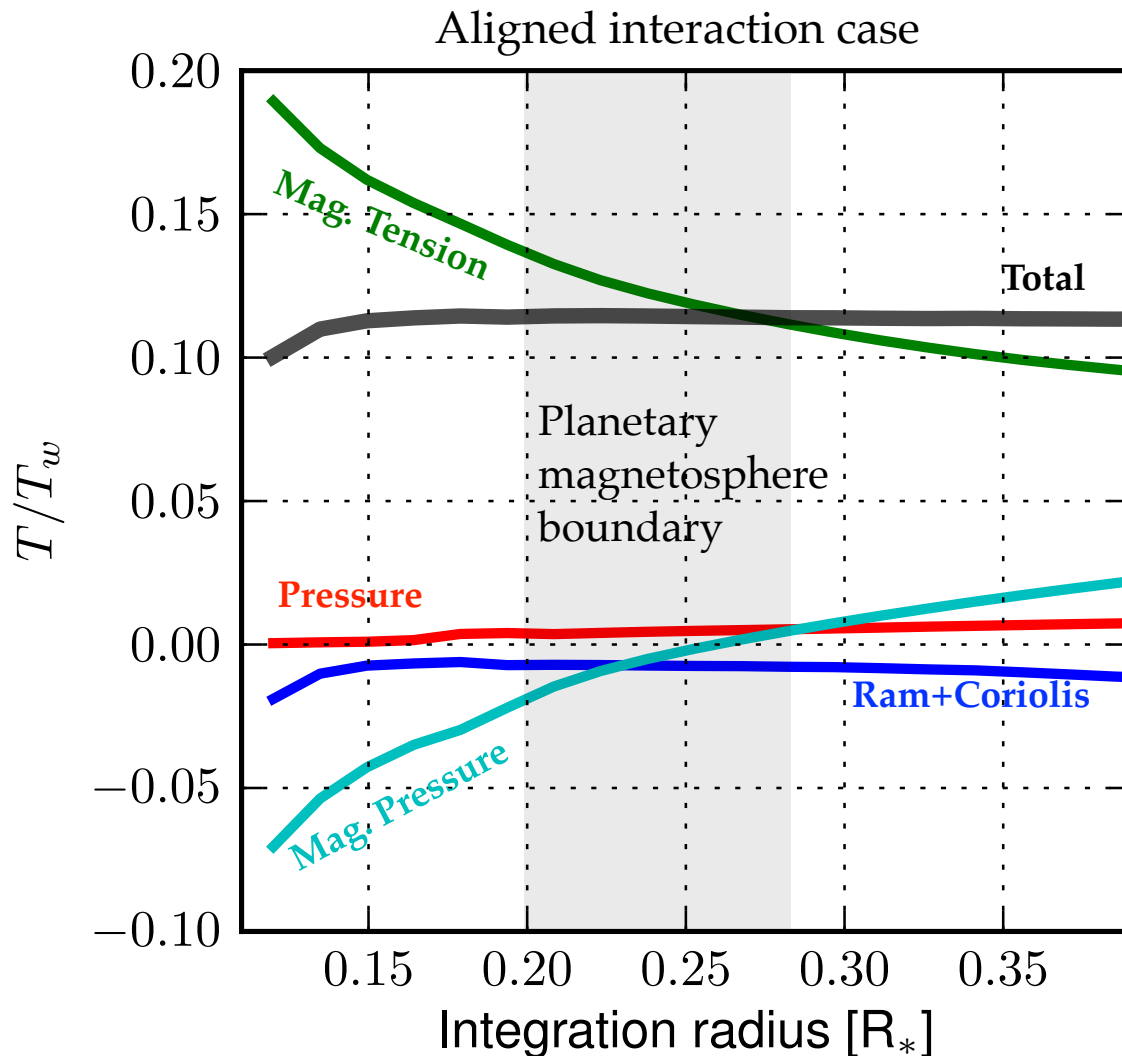


The two magnetic contributions dominate in the anti-aligned case

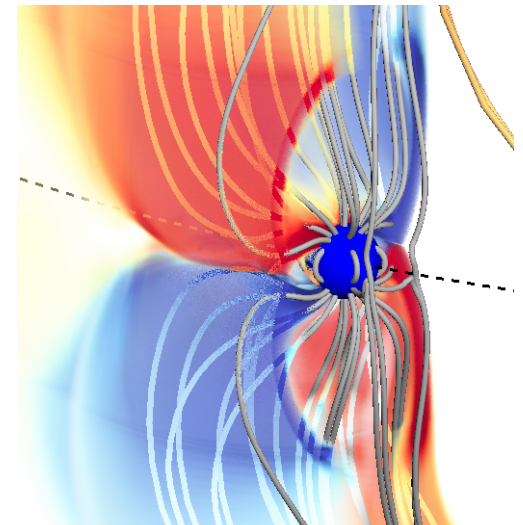


For a standard T Tauri-Star ( $\dot{M} \sim 10^5 \dot{M}_{\text{sun}}$ ),  $\tau_{\text{mig}} \sim 850 \text{ Myr}$

# Magnetic torque applied to close-in planets



The magnetic tension (field lines connecting the planet to the star) is the main contributor to the planet migration



For a standard T Tauri-Star ( $\dot{M} \sim 10^5 \dot{M}_{\text{sun}}$ ),  $\tau_{\text{mig}} \sim 140 \text{ Myr}$

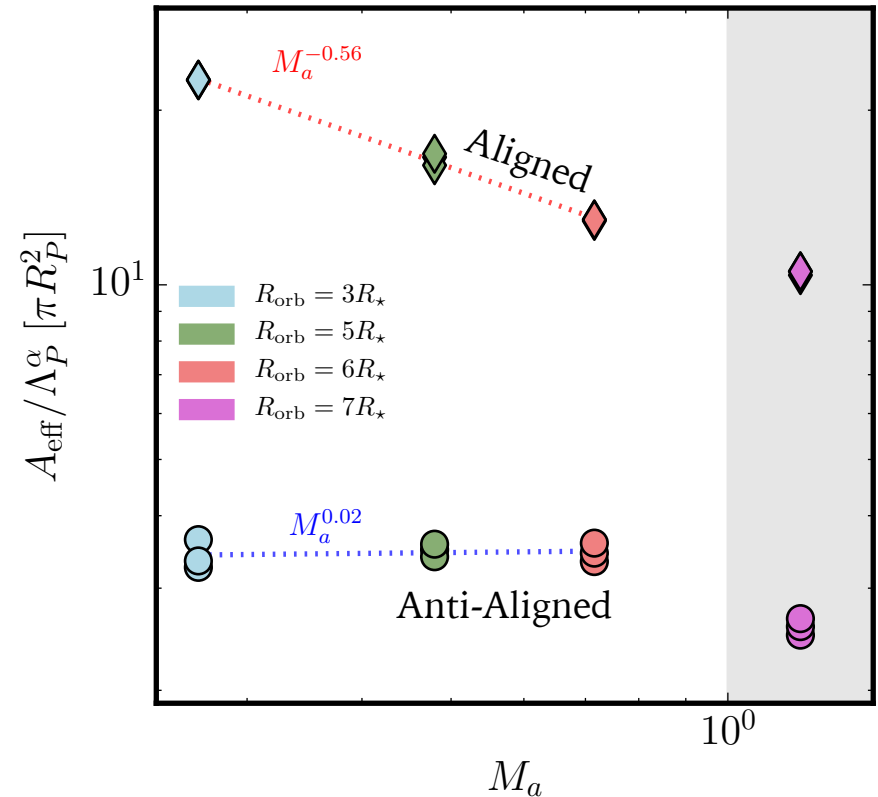
# Magnetic torque: parametrization

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Wind-dependant parameters  
(pressure, magnetic field, density)

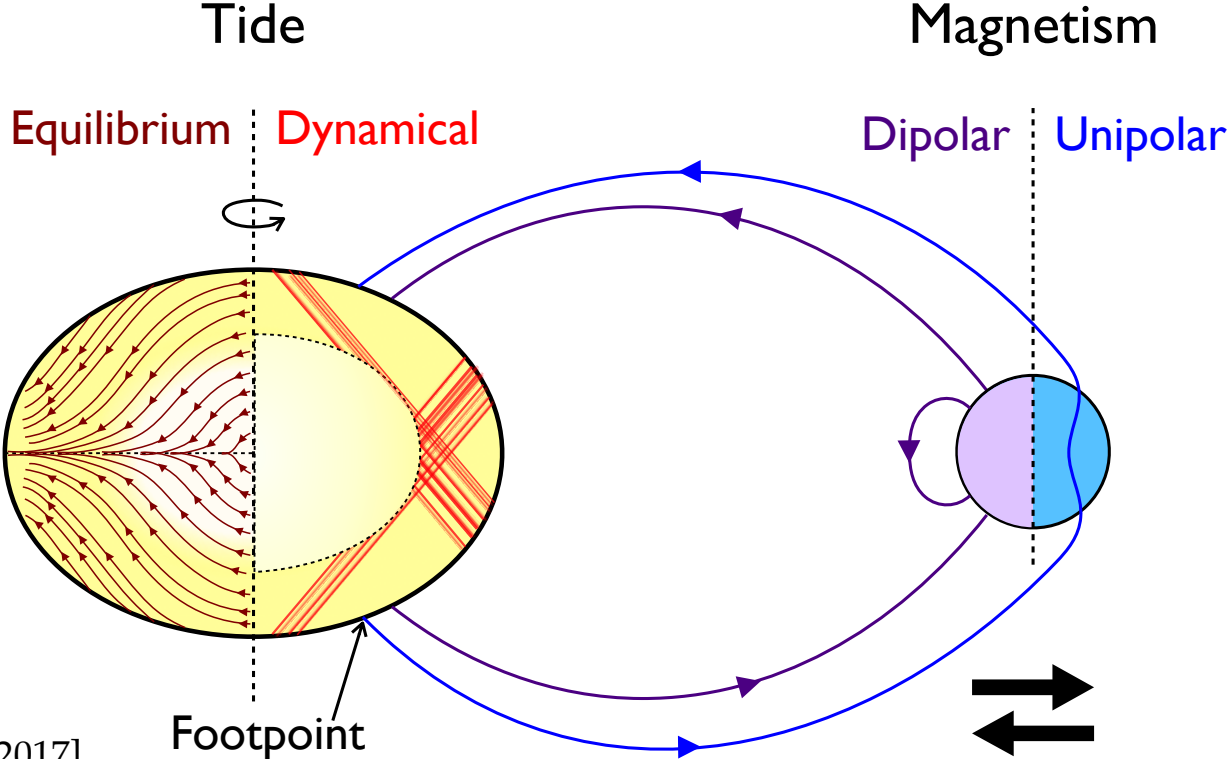
Planet-dependant parameters  
(radius, orbital radius, composition)

Magnetosphere/wind pressure ratio  
(planetary magnetic field, wind)



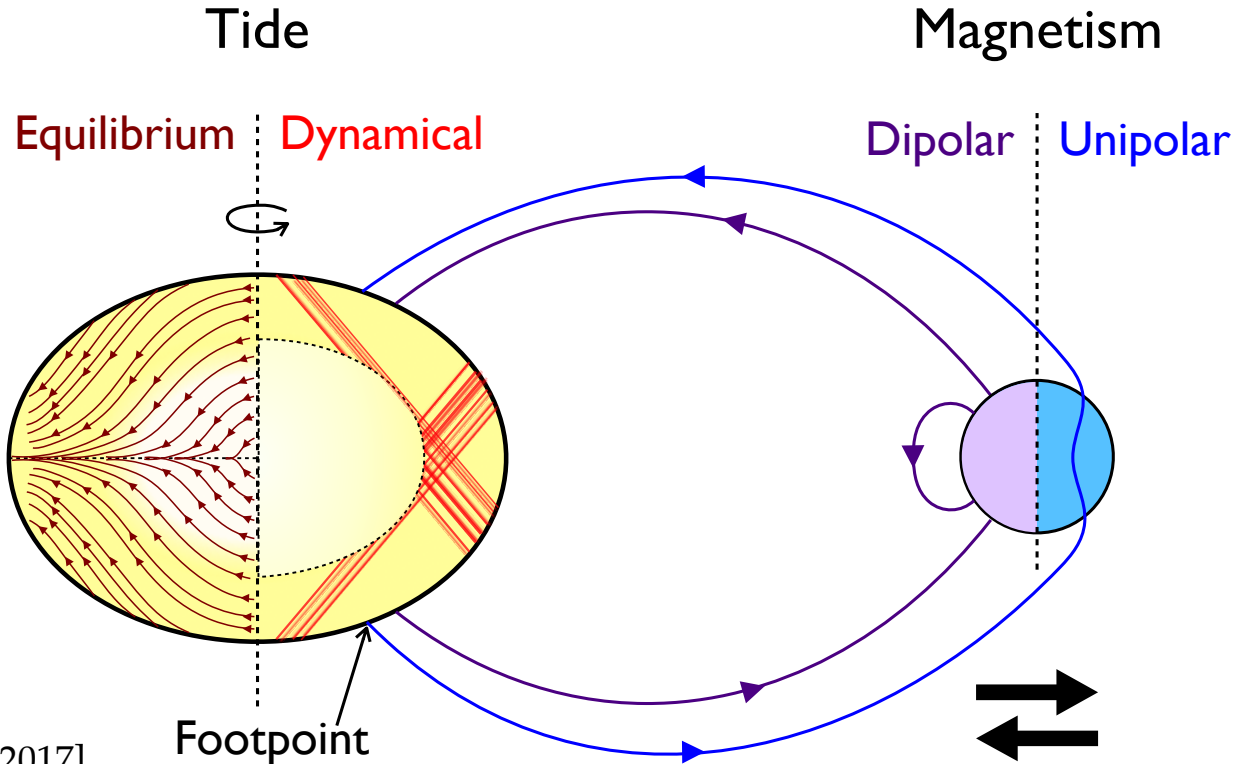


# Let's compare tidal and magnetic torques



[Strugarek+, ApJL 2017]

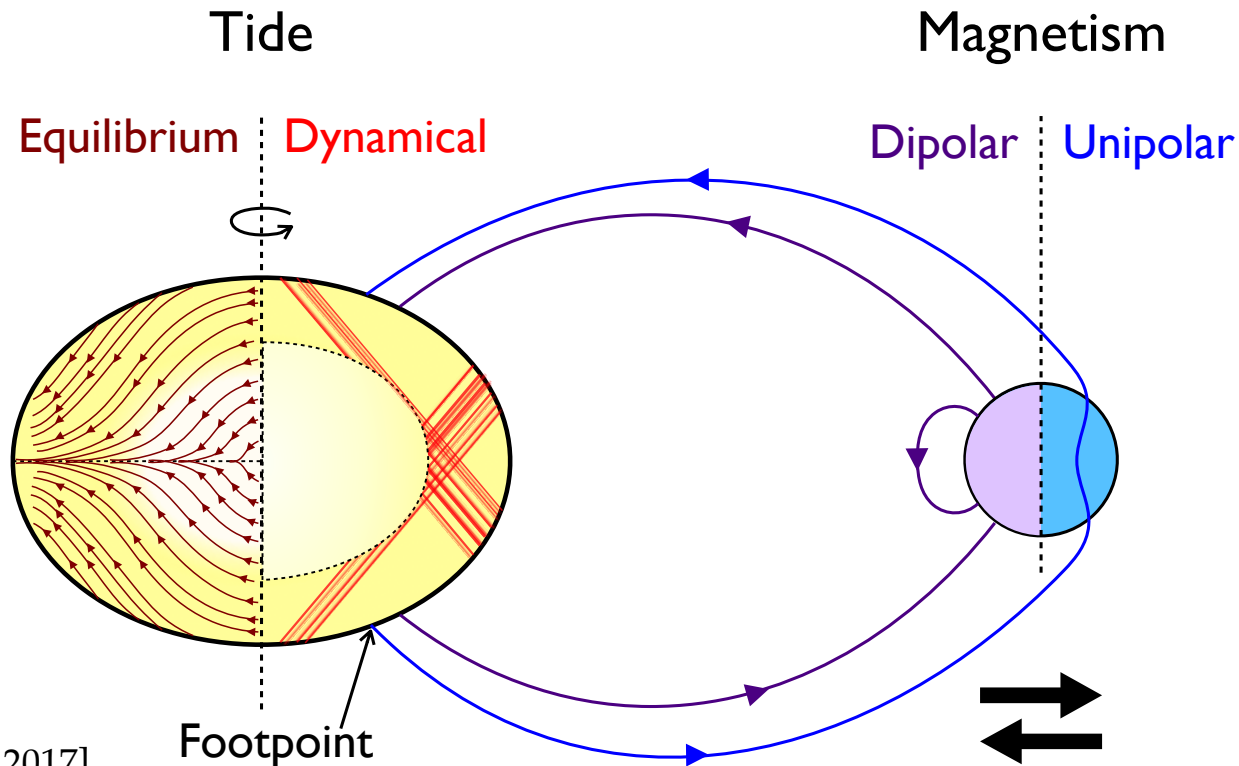
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Stellar evolution code (STAREVOL) from Gallet+ 2017

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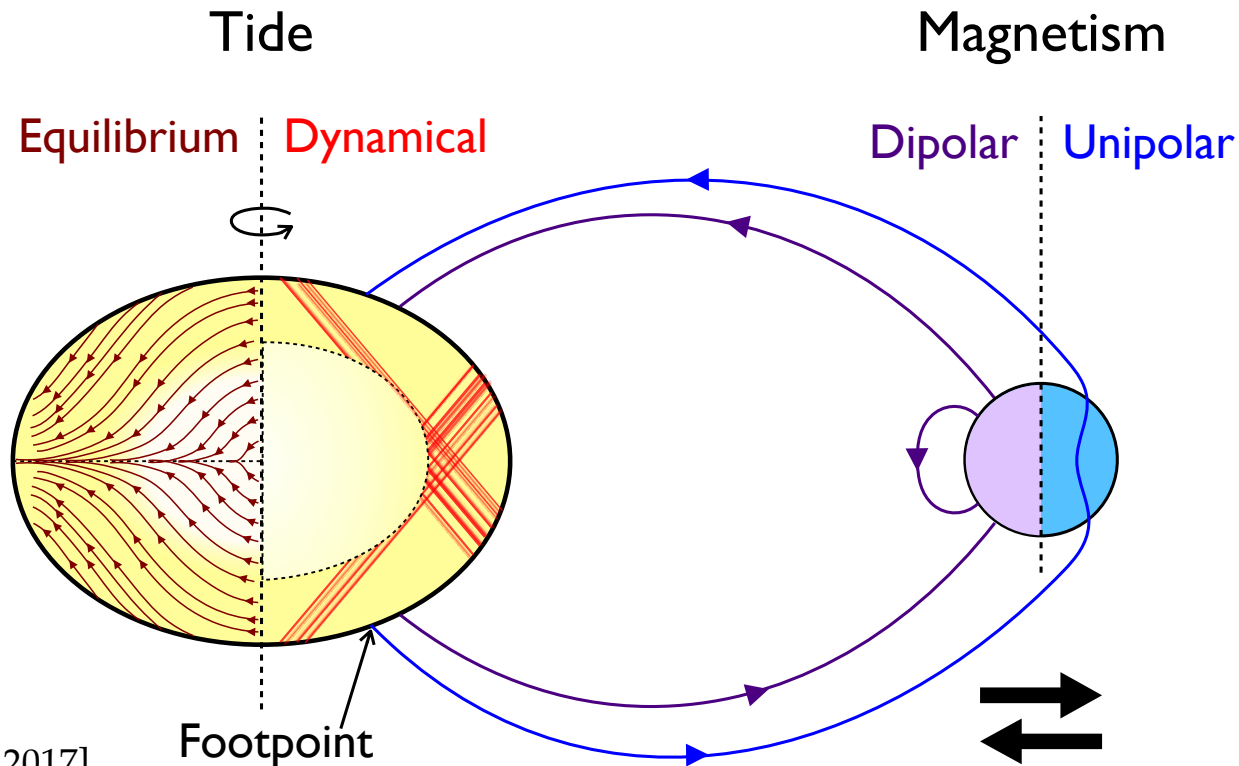


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Stellar evolution code (STAREVOL) from  
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Stellar wind model  
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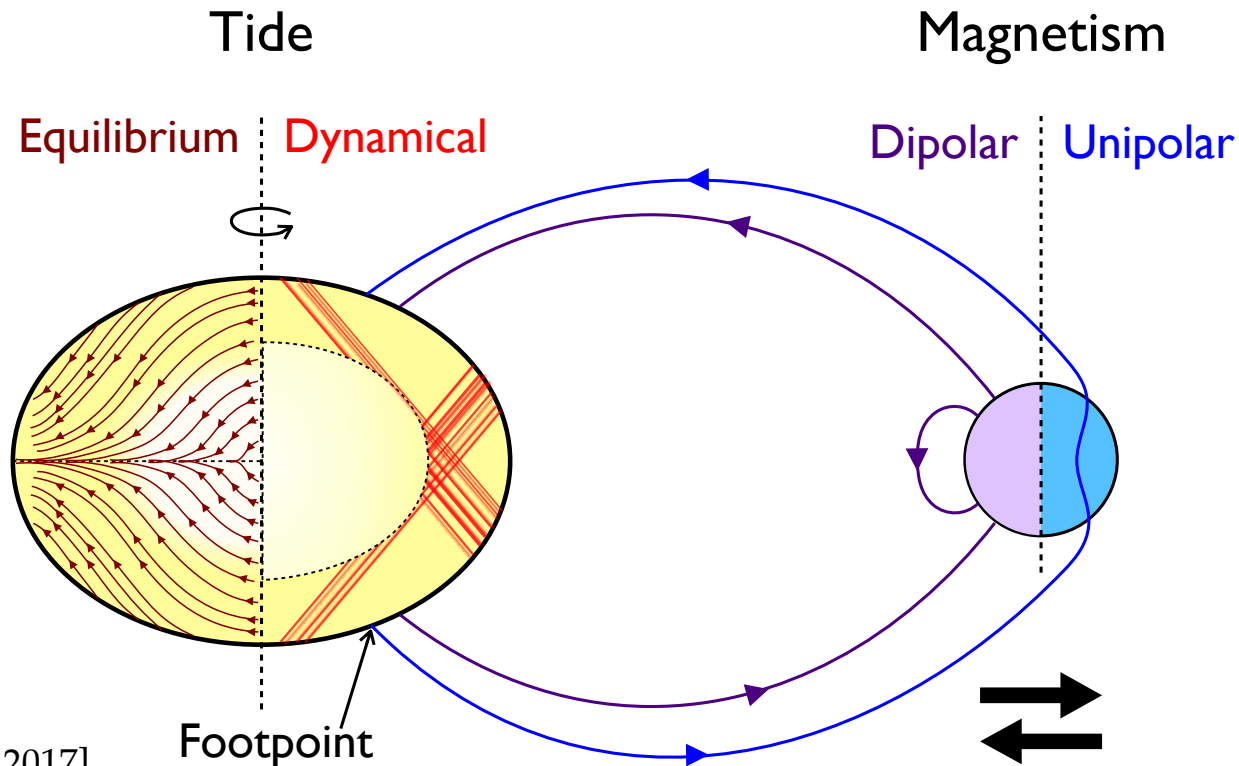
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Tidal torque formulation  
Remus+ 2012, Mathis 2015, Mathis+ 2016

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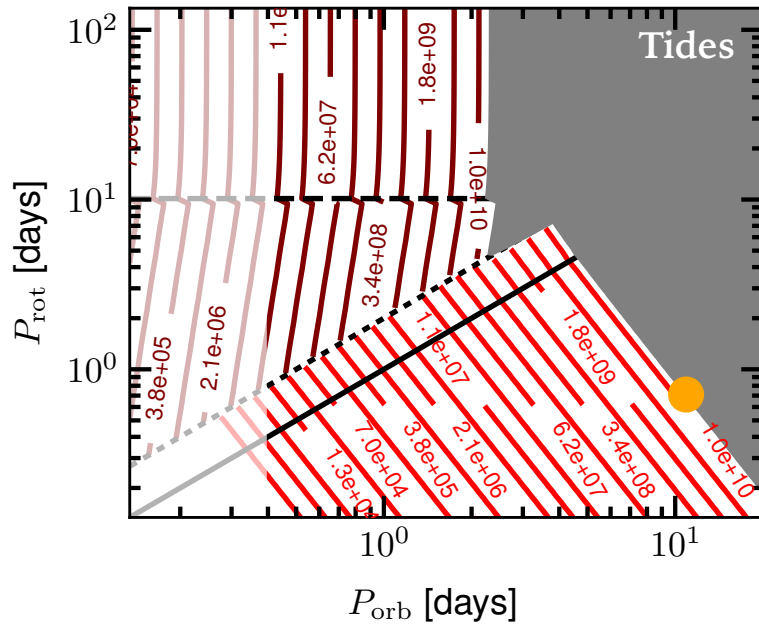
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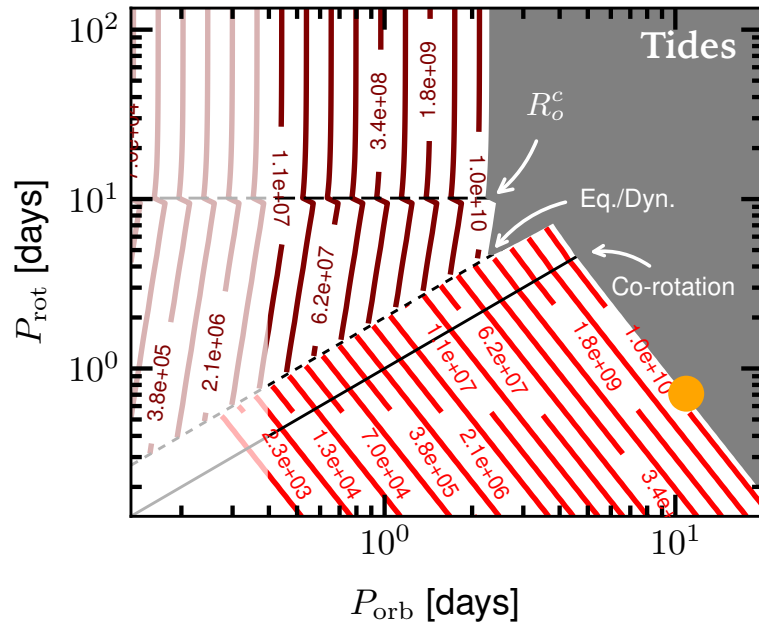
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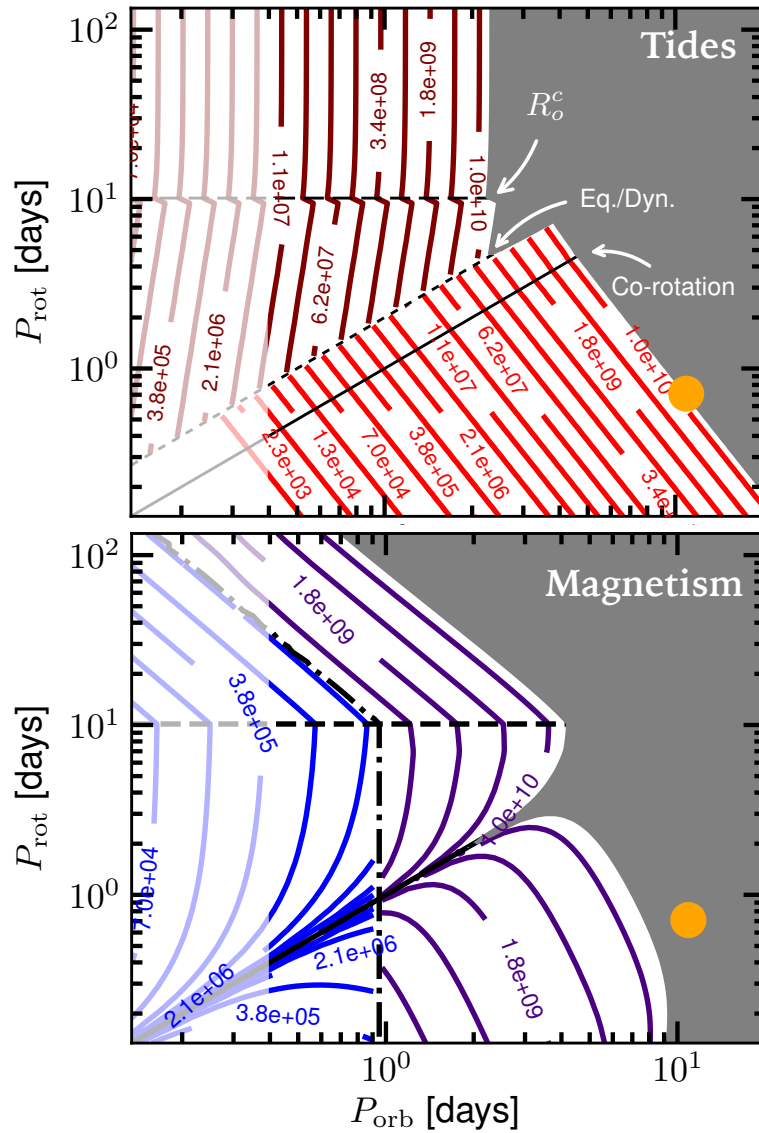
# Application: $\Pi$ auri star & hot Jupiter (e.g. Tap 26)



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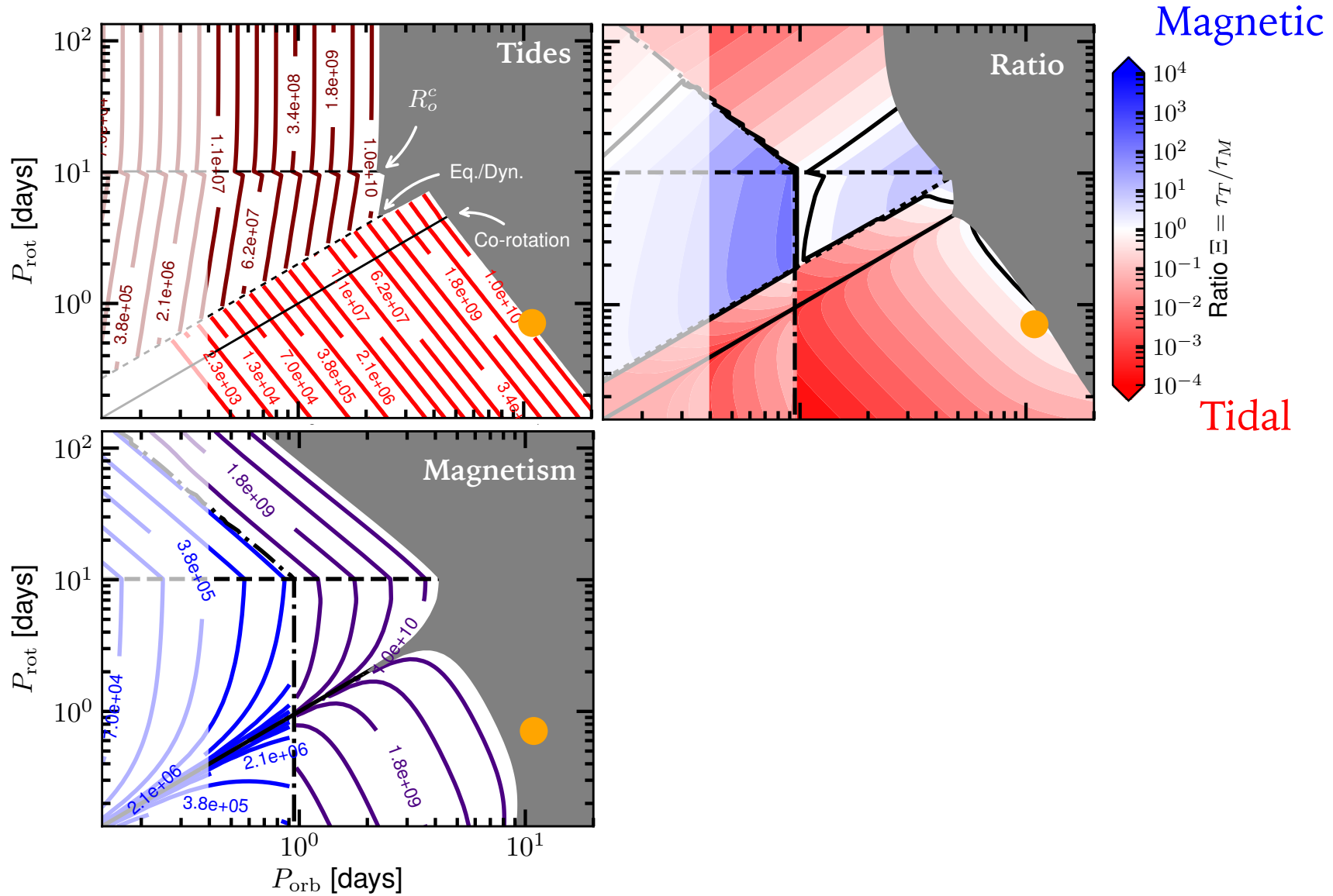


# Application: $\tau$ Tauri star & hot Jupiter (e.g. Tap 26)

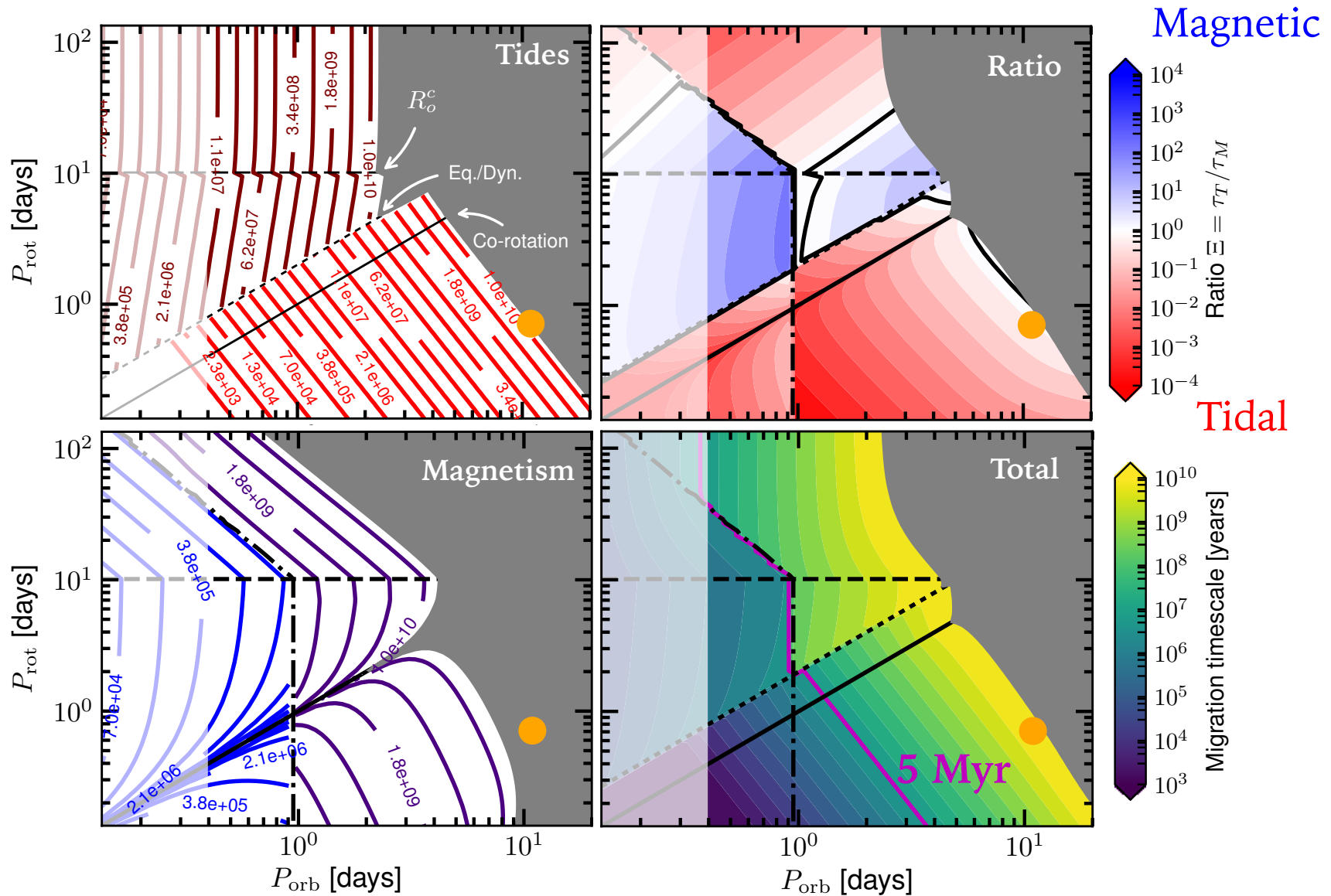




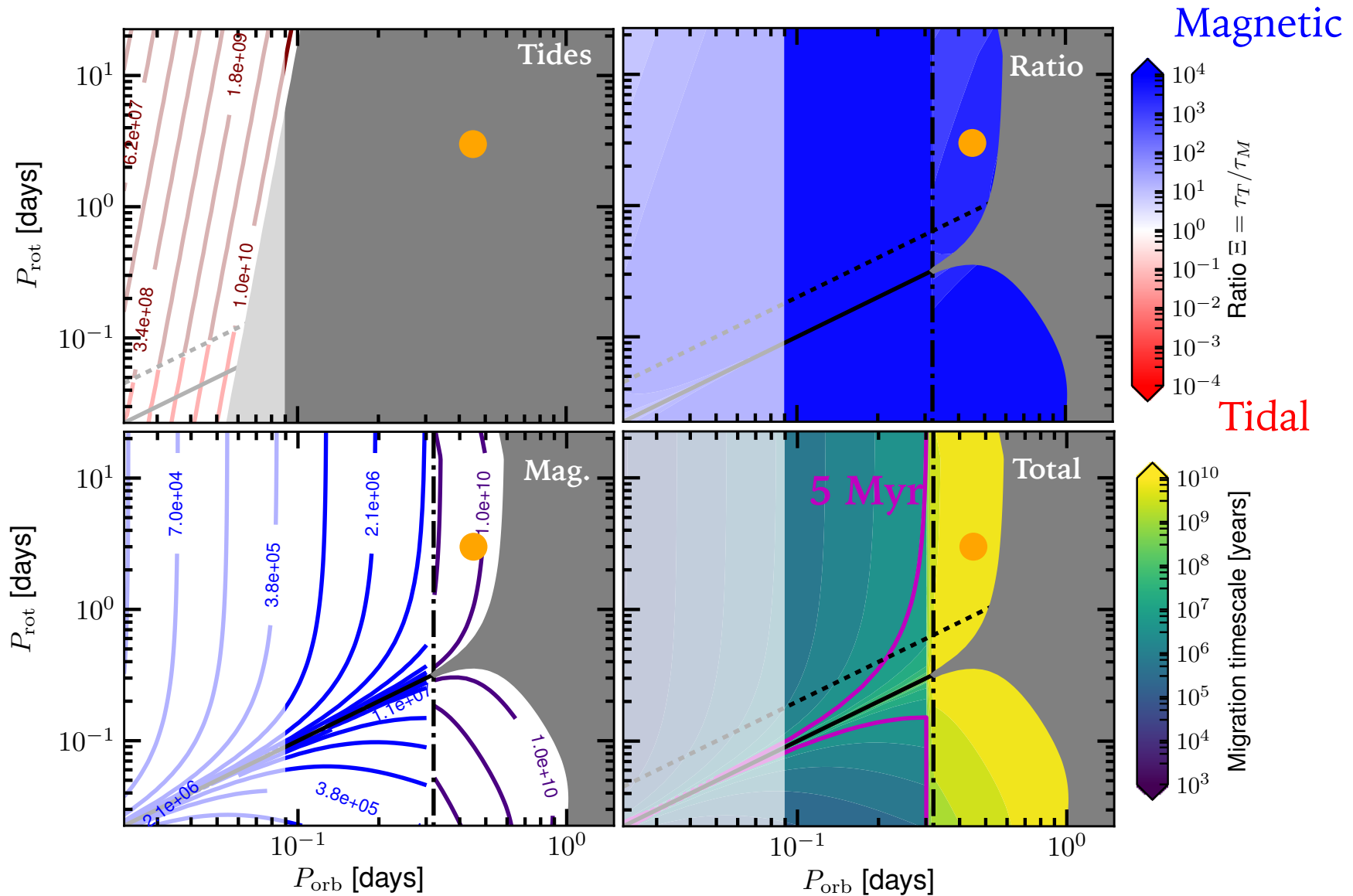
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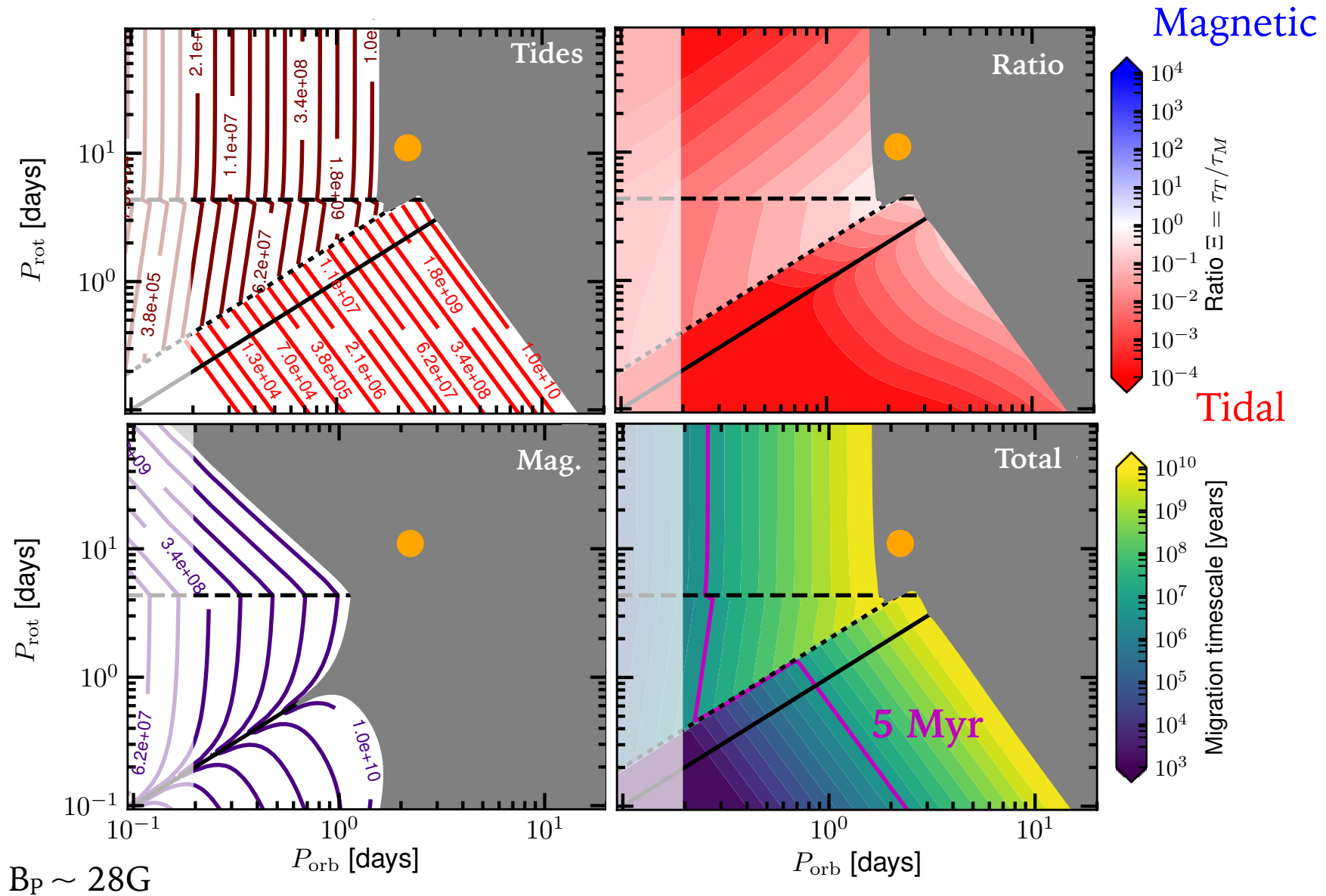
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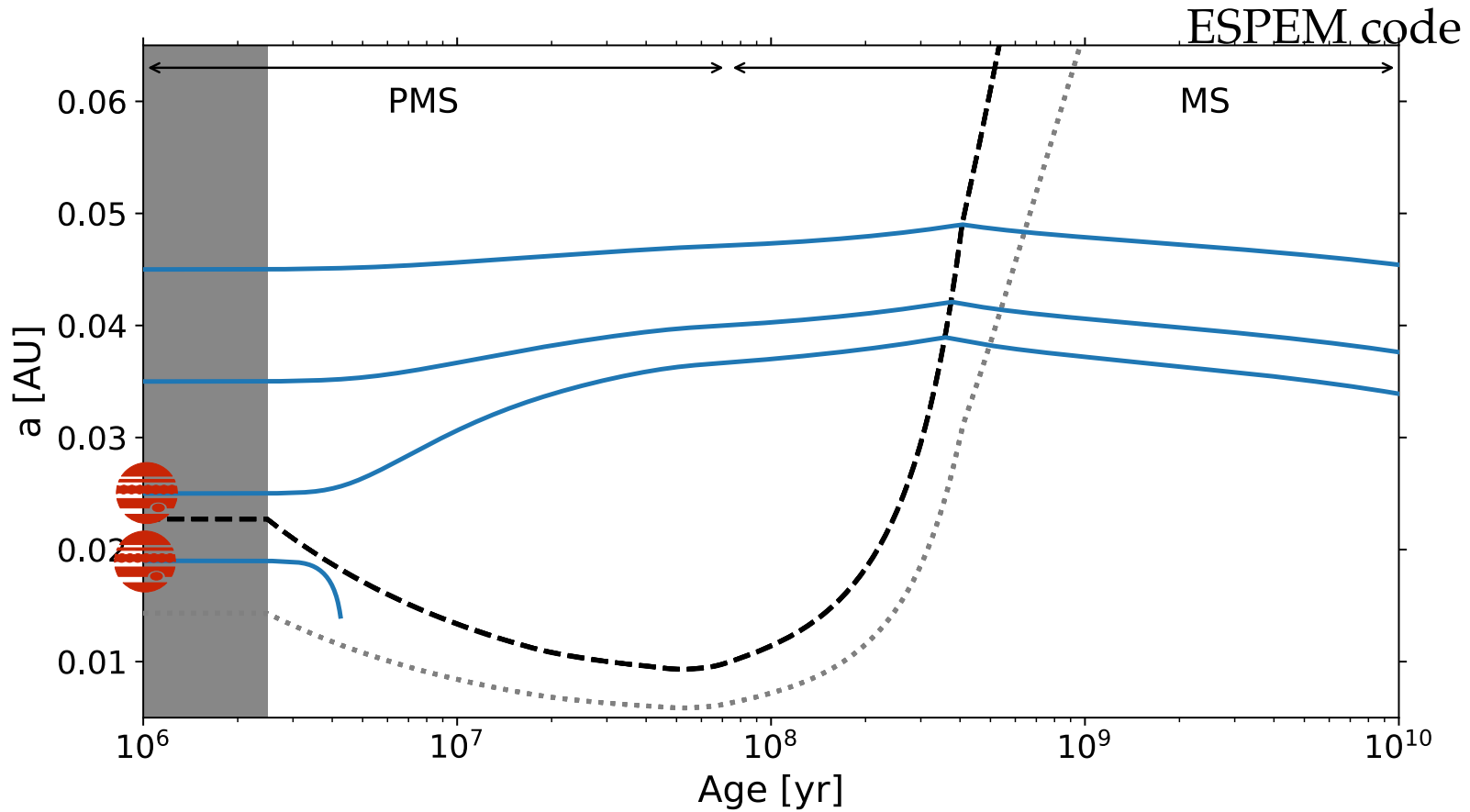
# Application: M dwarf & Earth-like planet (e.g. Kepler-42)



# Application: K star & hot Jupiter (e.g. HD 189733)

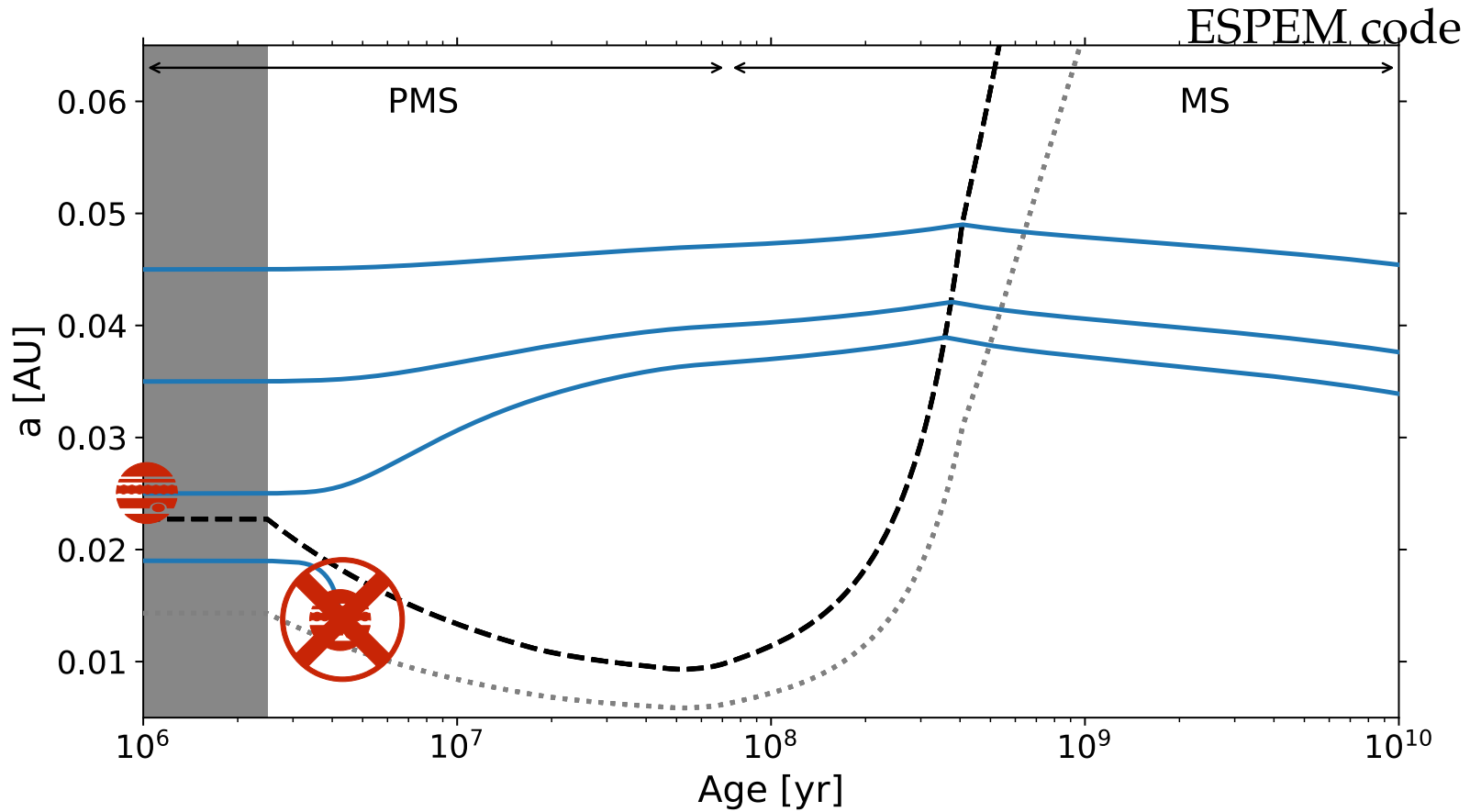


# Close-in planets migration due to tidal and



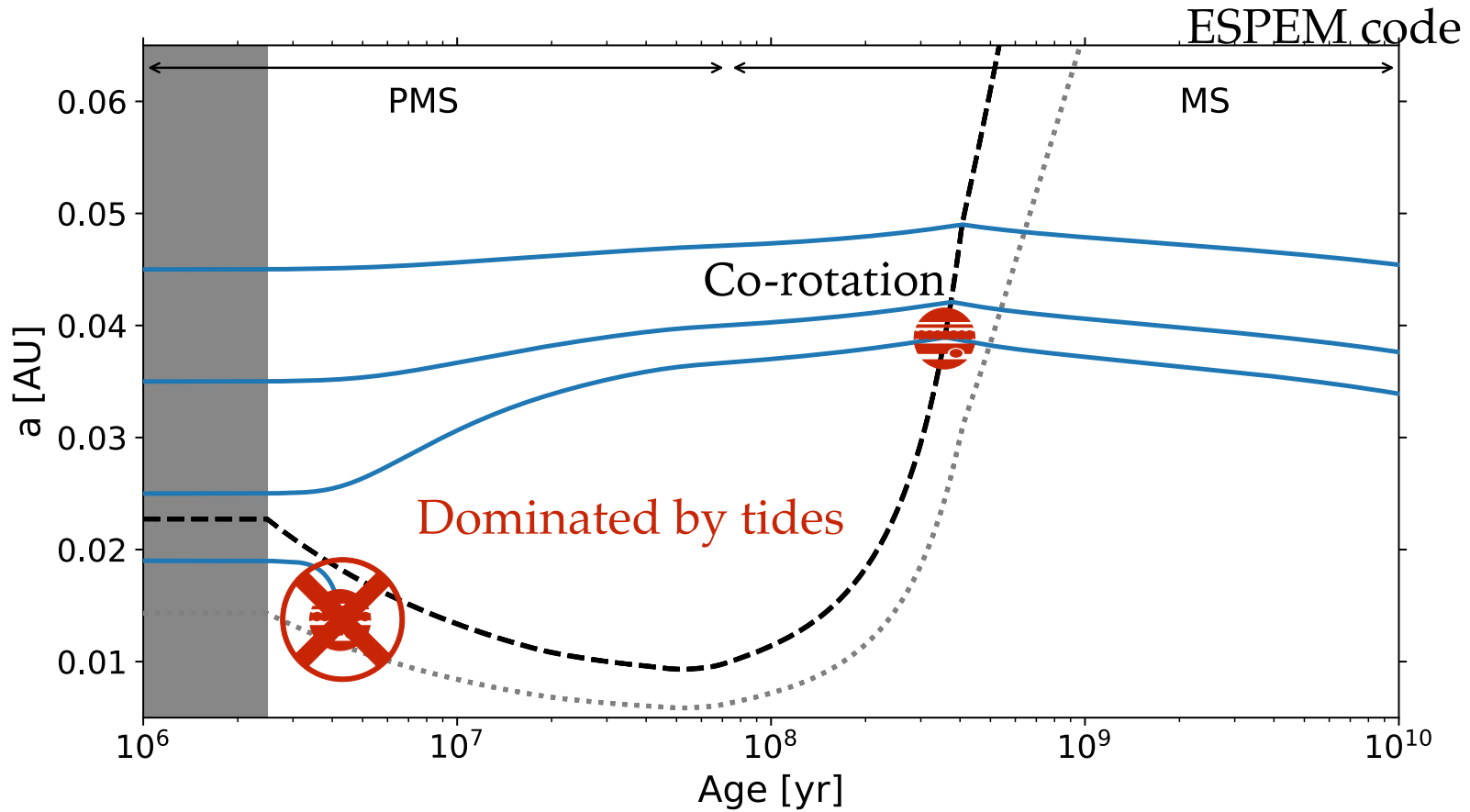
[Benbakoura+ 19; Ahuir+ 21]

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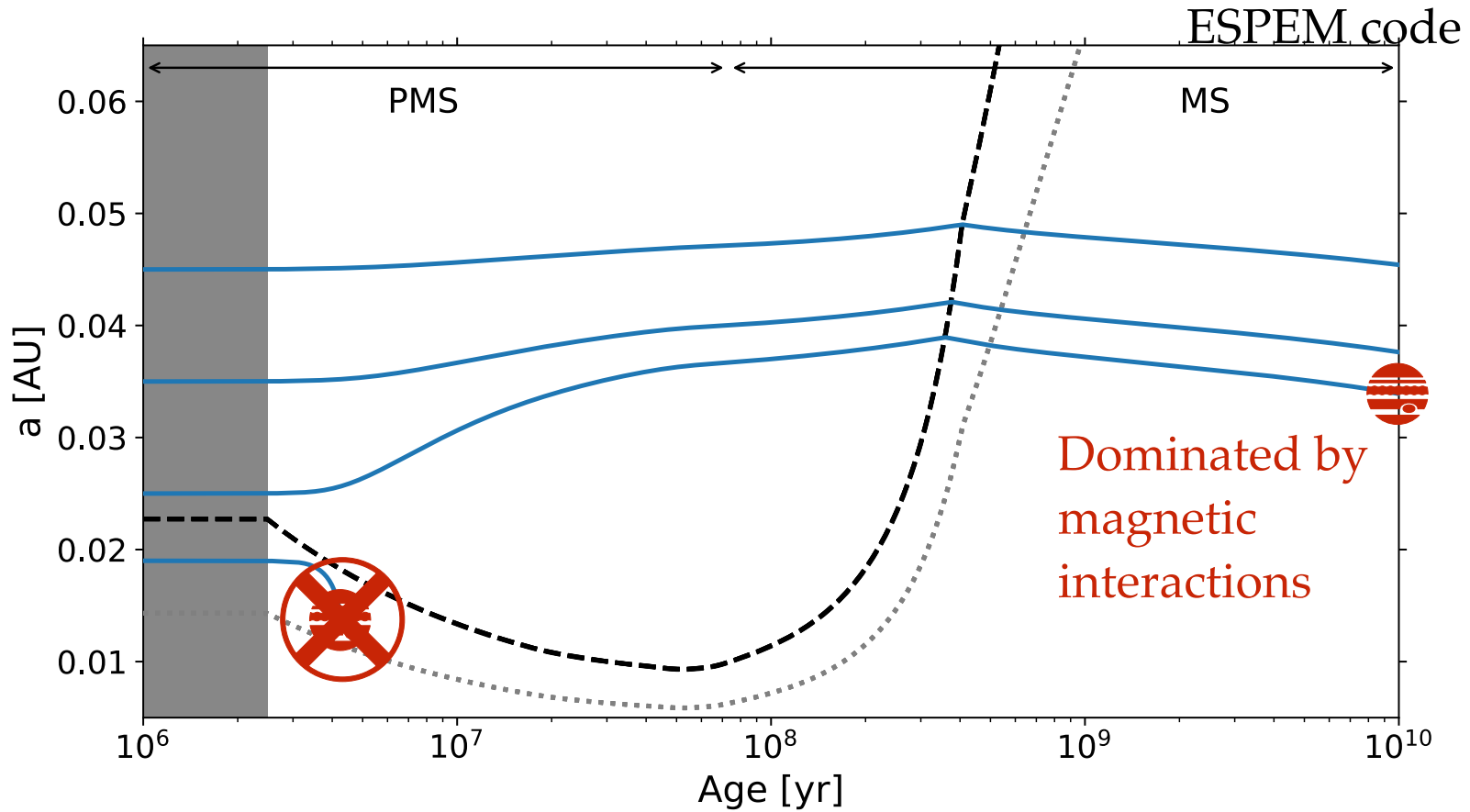
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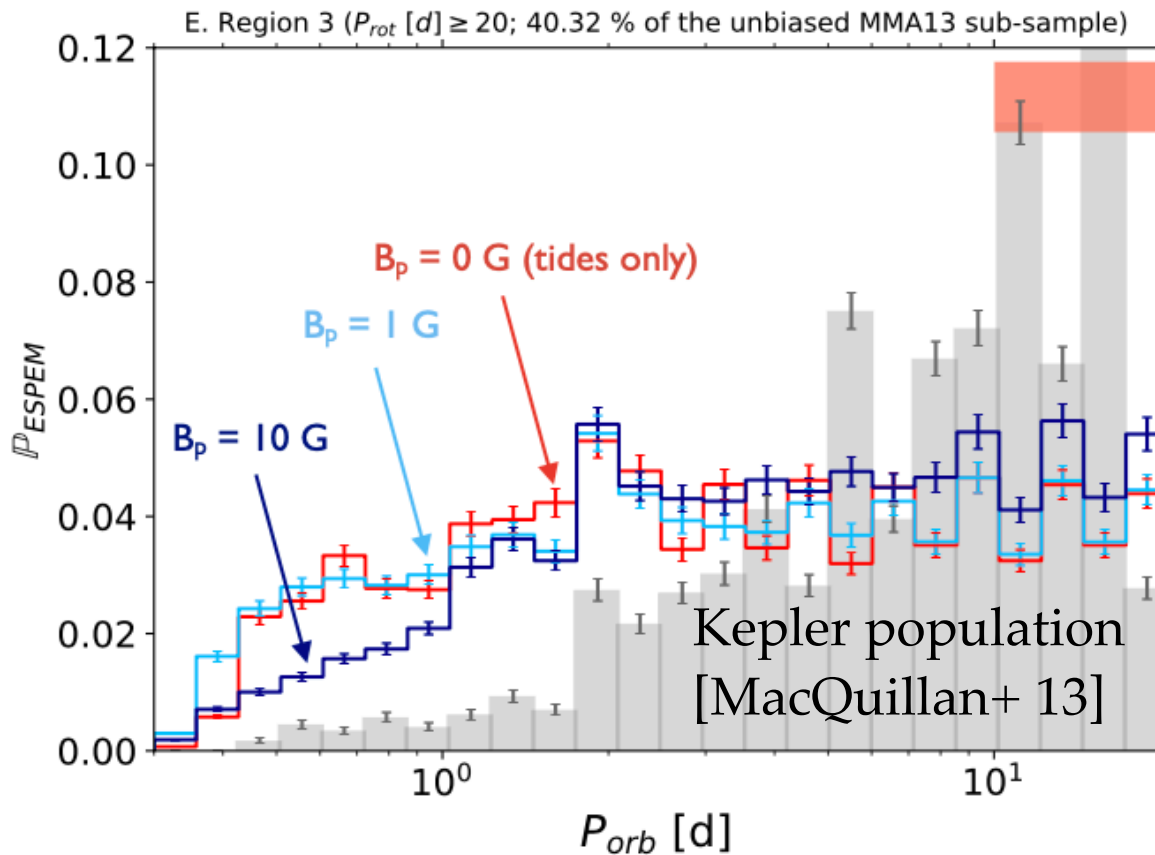
# Close-in planets migration due to tidal and



[Benbakoura+ 19; Ahuir+ 21]



# Synthetic population vs Kepler-field population



Magnetic effects shape the distribution of close-in **super-Earths** ( $P_{orb} < 1d$ ) around **slow rotators** ( $P_{rot} > 5d$ )

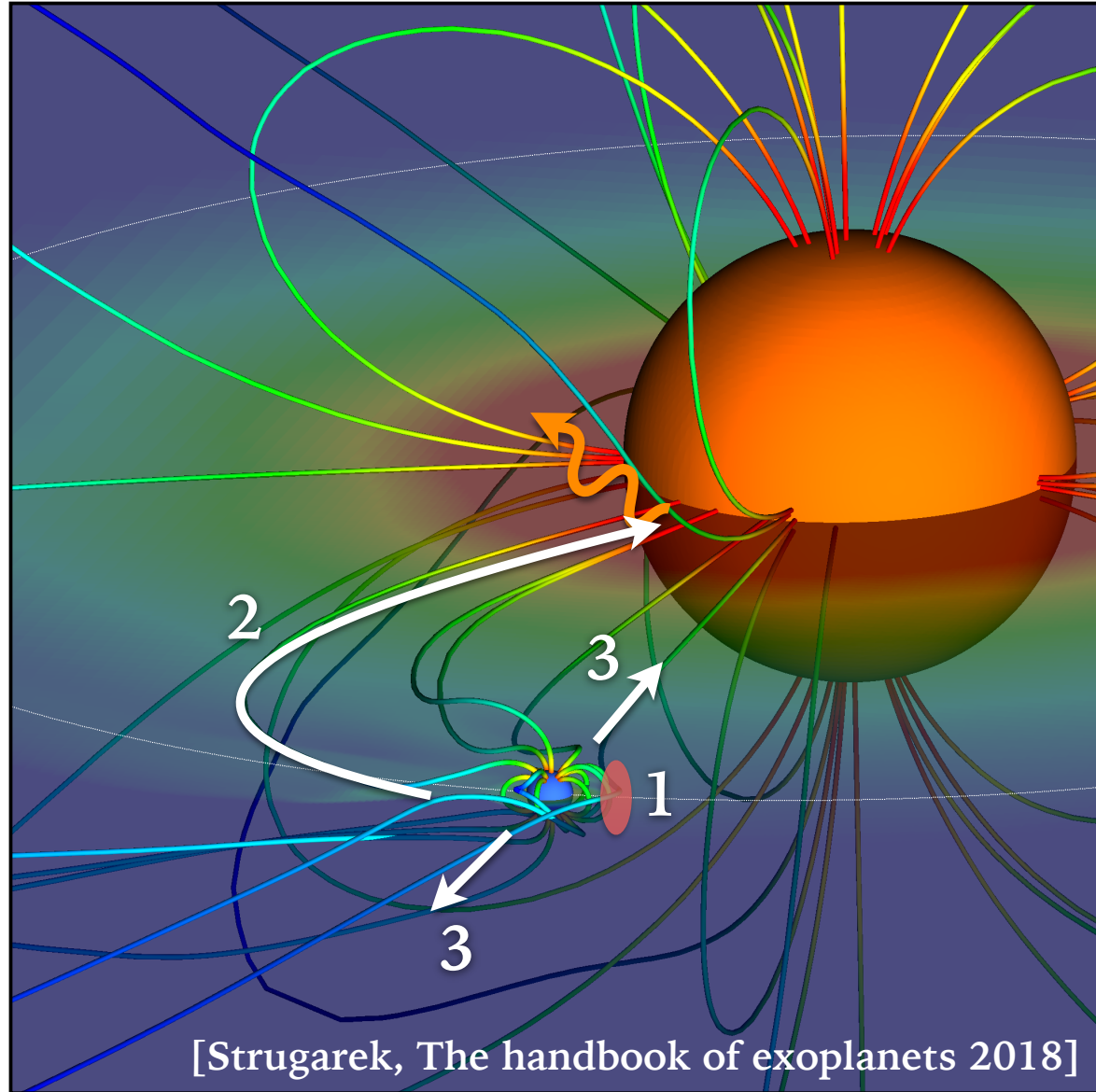
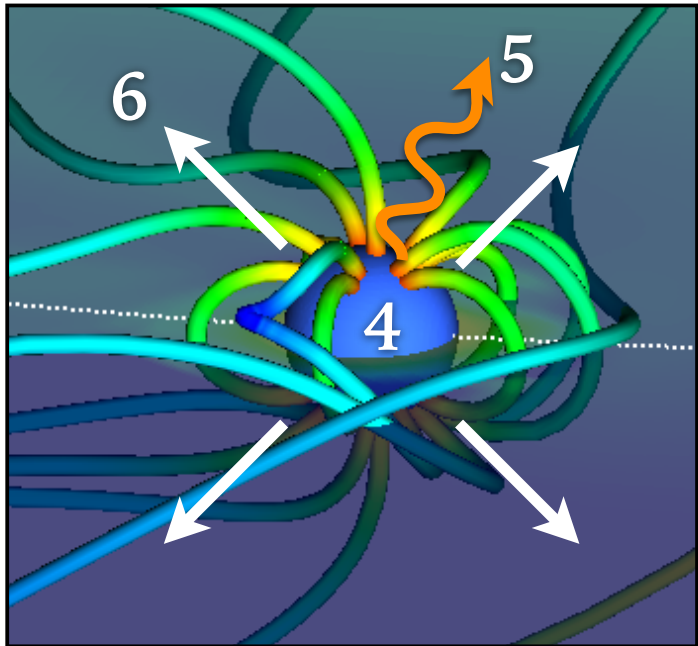
But we still have an excess of hot planets in our synthetic sample:

- Initial conditions effects?
- Extreme planetary dynamo?
- Additional Star-Planet interactions?

[Ahuir+ 21]

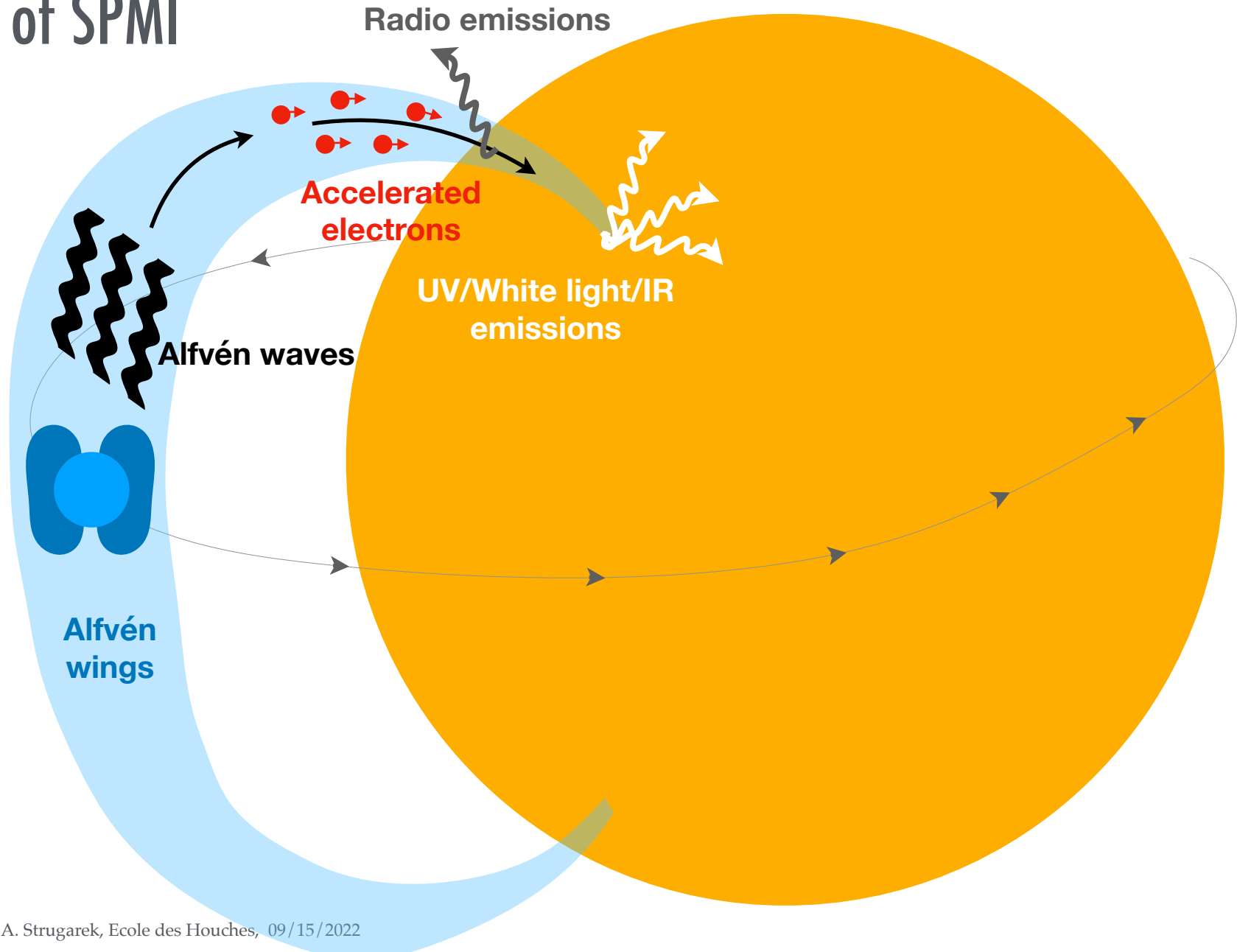
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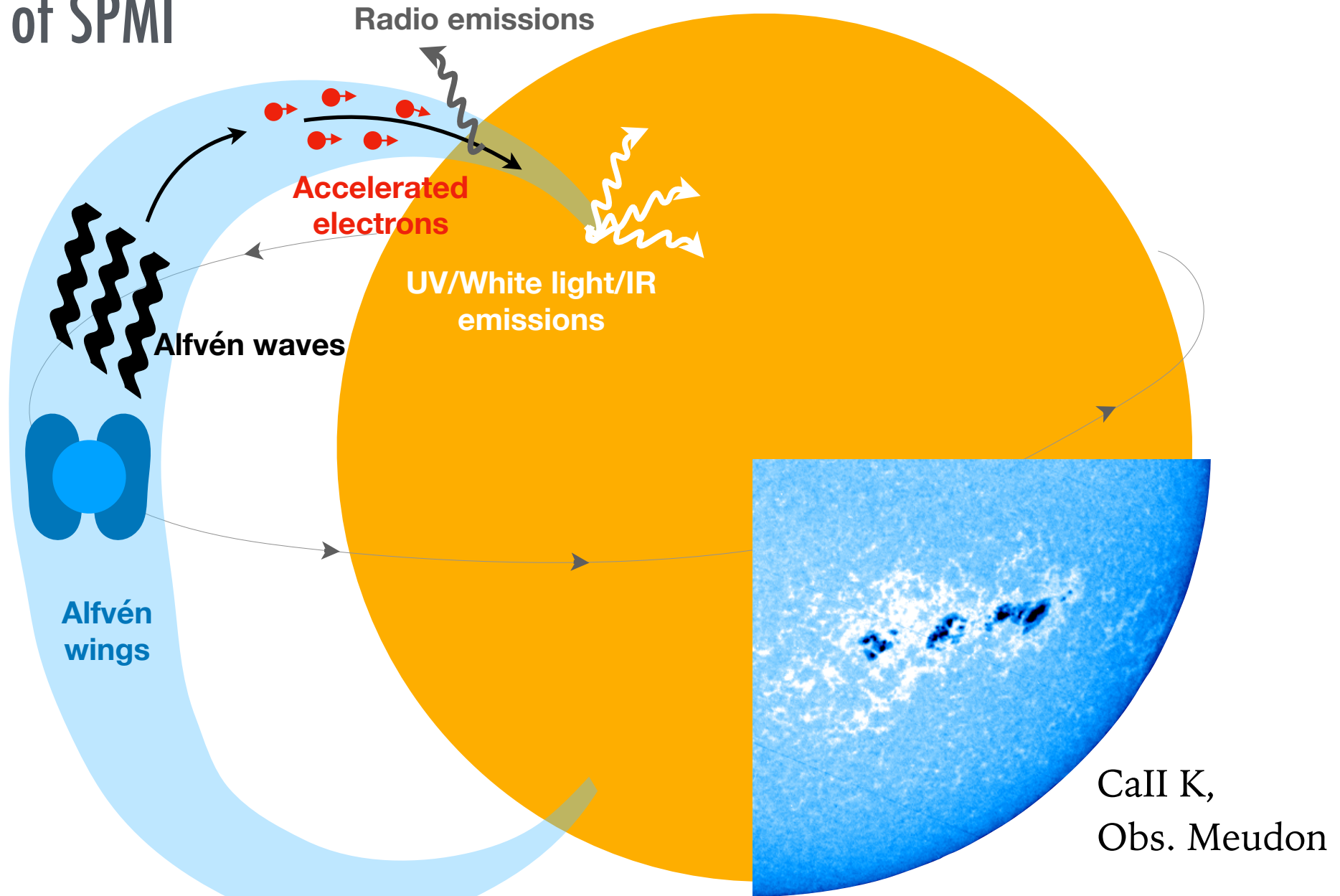


[Strugarek, The handbook of exoplanets 2018]

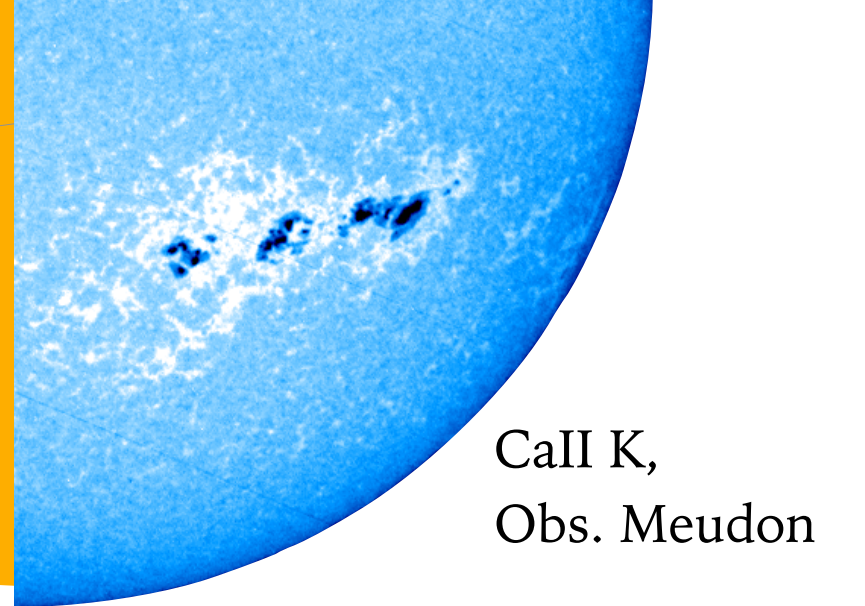
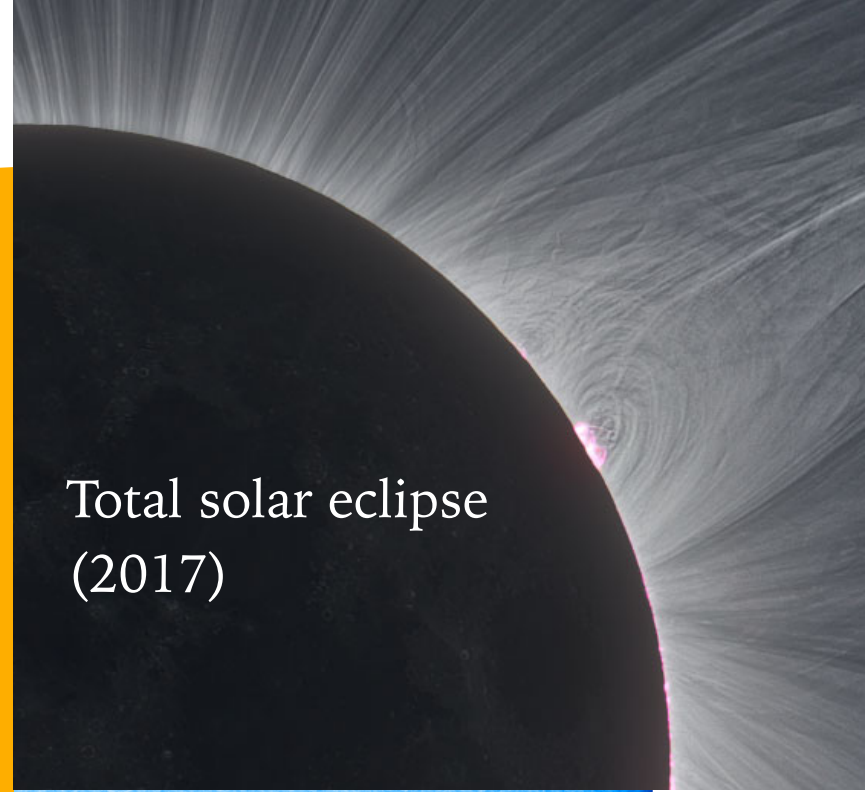
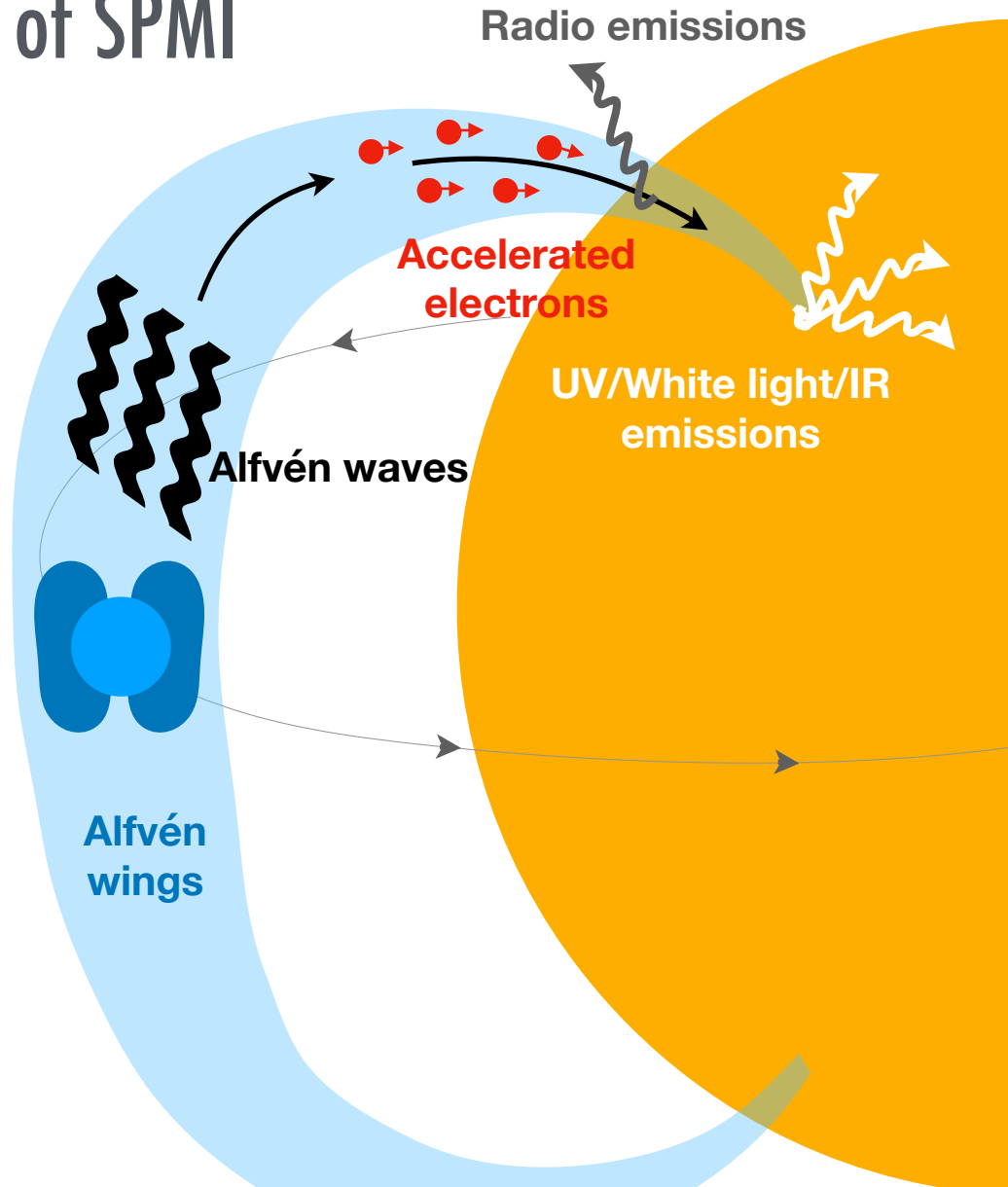
# Perspectives: the missing aspects of SPMI



# Perspectives: the missing aspects of SPMI



# Perspectives: the missing aspects of SPMI



# Third 'atelier ExoSystèmes' for the french community

January 10-12 2023, Marseille

@NASA-JPL / Caltech

SOC

F. Cantaloube

M. Deleuil

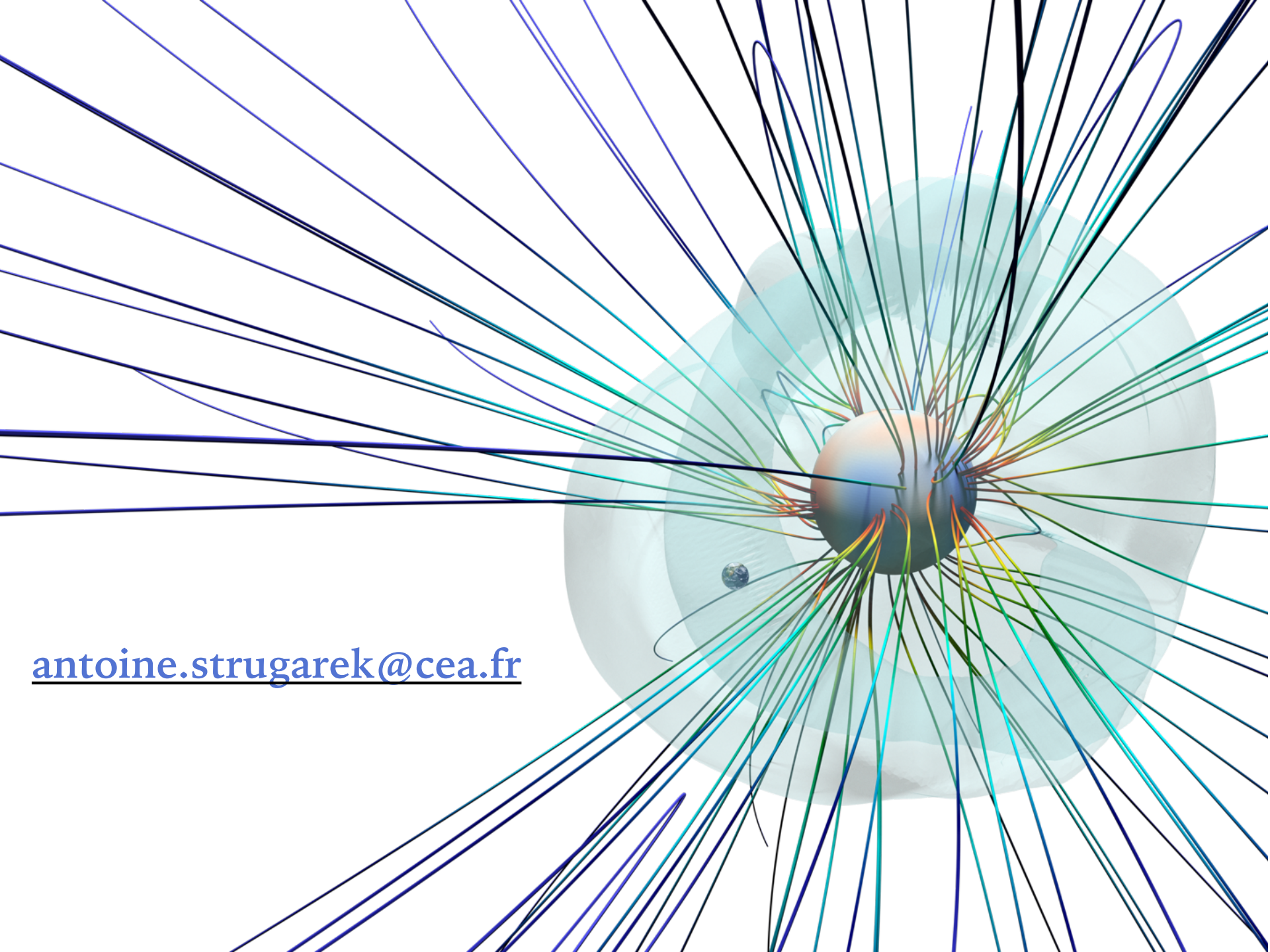
Q. Kral

P.O. Lagage

J. Leconte

A. Santerne

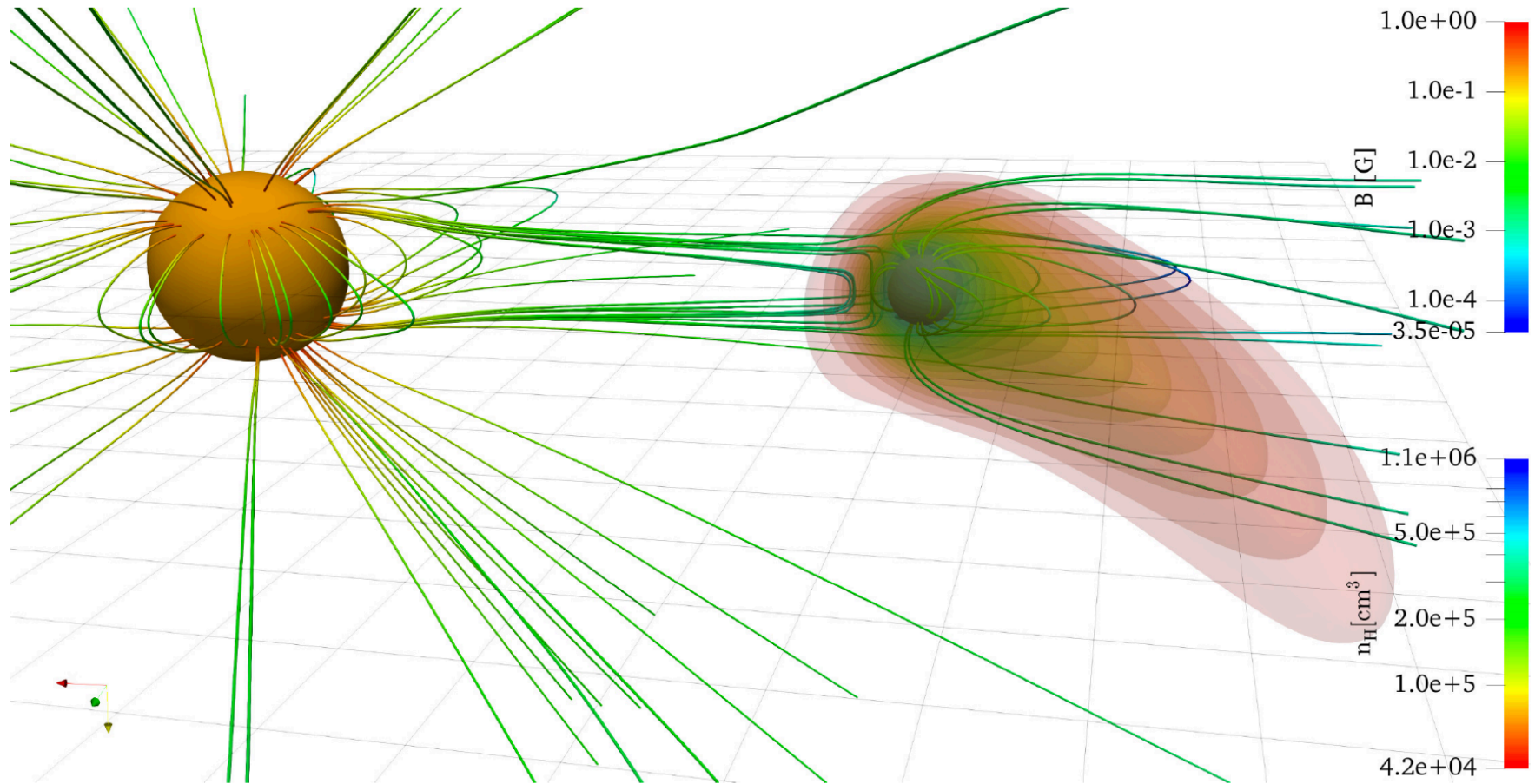
A. Strugarek



[antoine.strugarek@cea.fr](mailto:antoine.strugarek@cea.fr)

# Lyman-alpha absorption in HD 209458b

*Villarreal d'Angelo + 18*: MHD model including a neutral escaping atmosphere of H





# Lyman-alpha absorption in HD 209458b

Villarreal d'Angelo + 18: MHD model including a neutral escaping atmosphere of H

