

# Formation of atmospheres

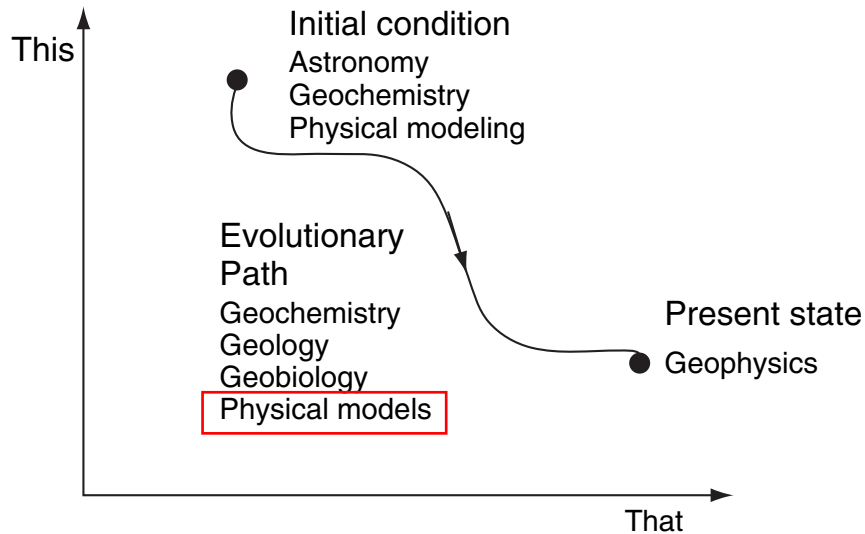


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*Hélène Massol - GEOPS - U. Paris Saclay*



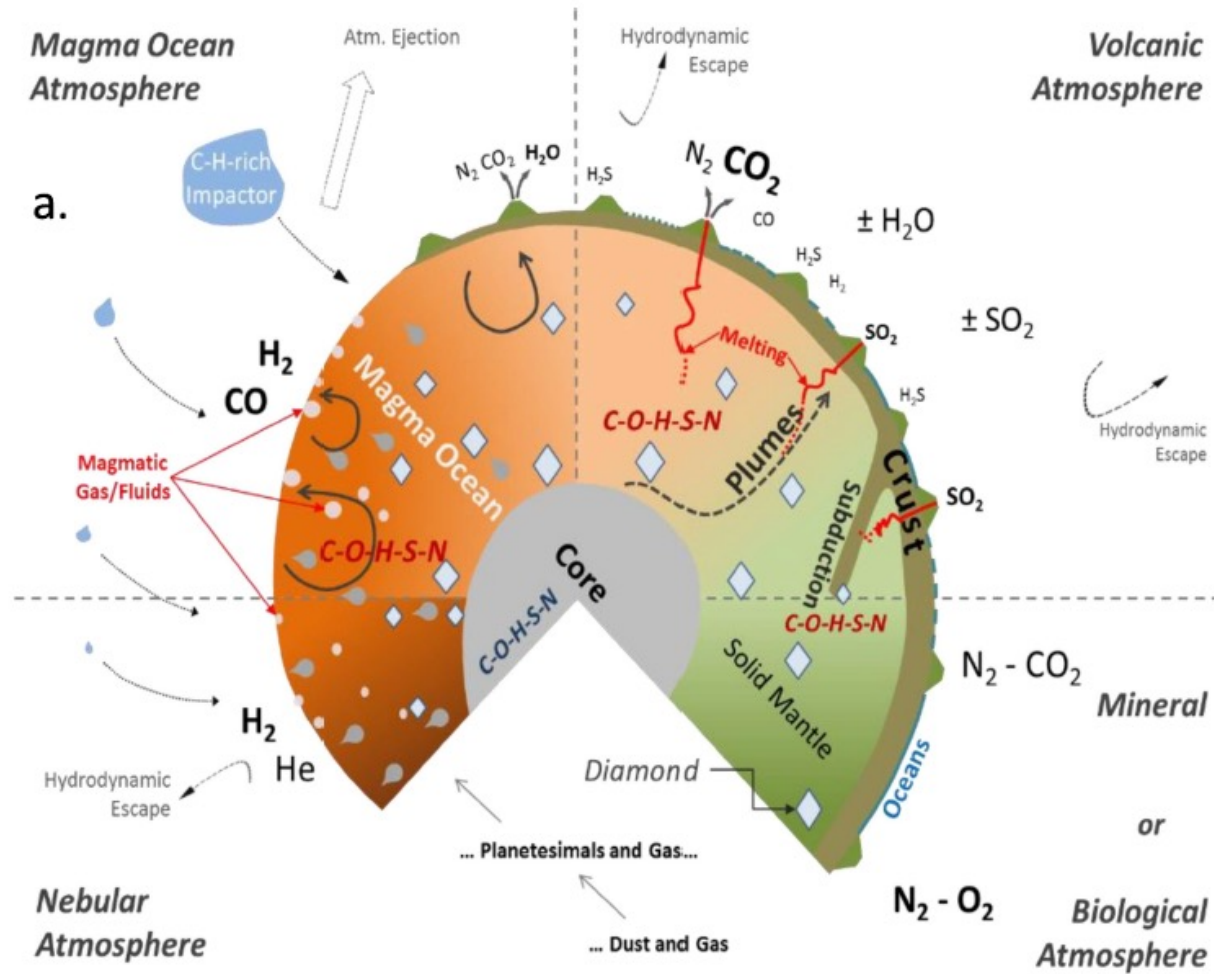


- Introduction / constraints
- Heat and volatile transfer during MO Phase
- Examples, results and applications to exo-planets

**Figure 1** Conceptual view of Earth evolution, identifying the three crucial elements (the initial condition, the evolution path, and the present state) and the sciences that contribute to their understanding. The axes are unimportant, since the diagram is merely a 2-D slice of a multidimensional phase space. They might represent temperature or composition, for example.

*Stevenson, 2014*

# Introduction / constraints: primary atmosphere



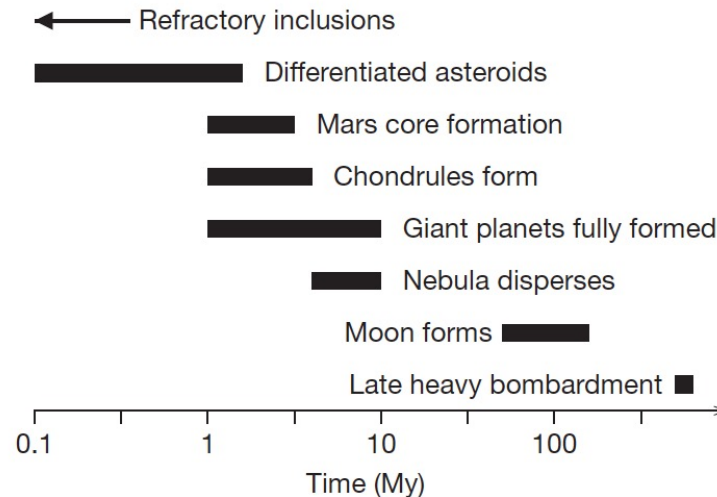
Gaillard et al., 2021

# Introduction / constraints: chronology of events

Dissipation of gas in the protoplanetary gas confirmed around 10 My by datation using short lived Pd-Ag system on Iron meteorites (*Hunt et al., 2022*)

Constraints on the time of formation of Earth

Hf-W + U pb: ~30 – 100 My and Mars : 1 – 10 My measured on martian meteorites

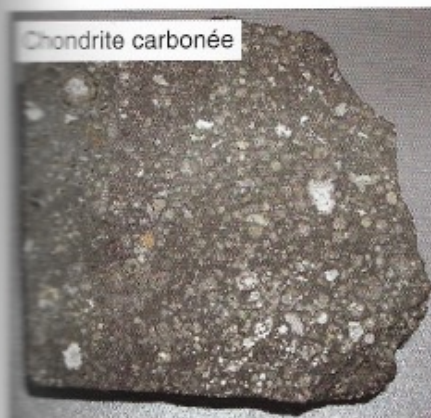


**Figure 1** The timing of events in the early Solar System (Data and models: [Dauphas and Pourmand, 2011](#); [Hutcheon et al., 2009](#); [Movshovitz et al., 2010](#); [Qin et al., 2008](#); [Sung et al., 2009](#); [Touboul et al., 2007](#)).

*Chambers, 2014*



# Introduction / constraints: different types of meteorites



Solidified magma



From core-mantle boundary



Fragments of planetary cores



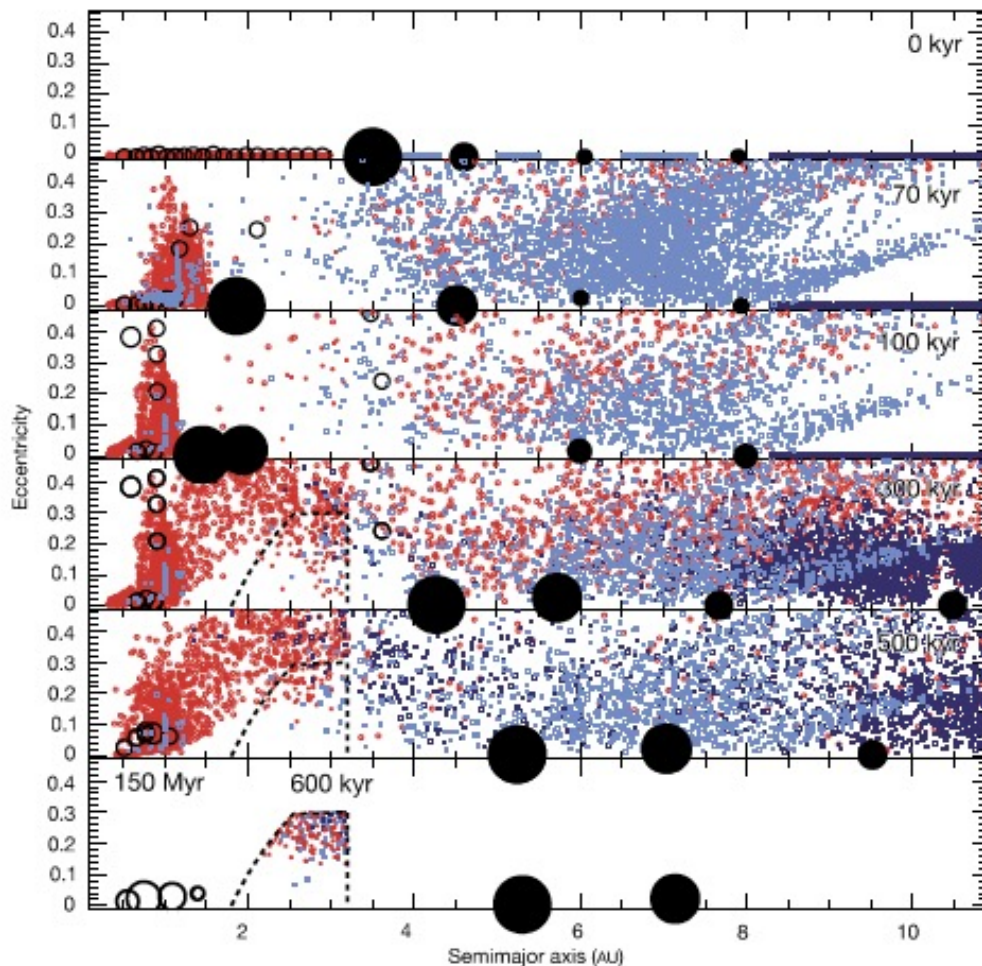


Fig. 3 Figure from Walsh et al. (2011) showing a simulation of the Grand Tack scenario. The top three panels

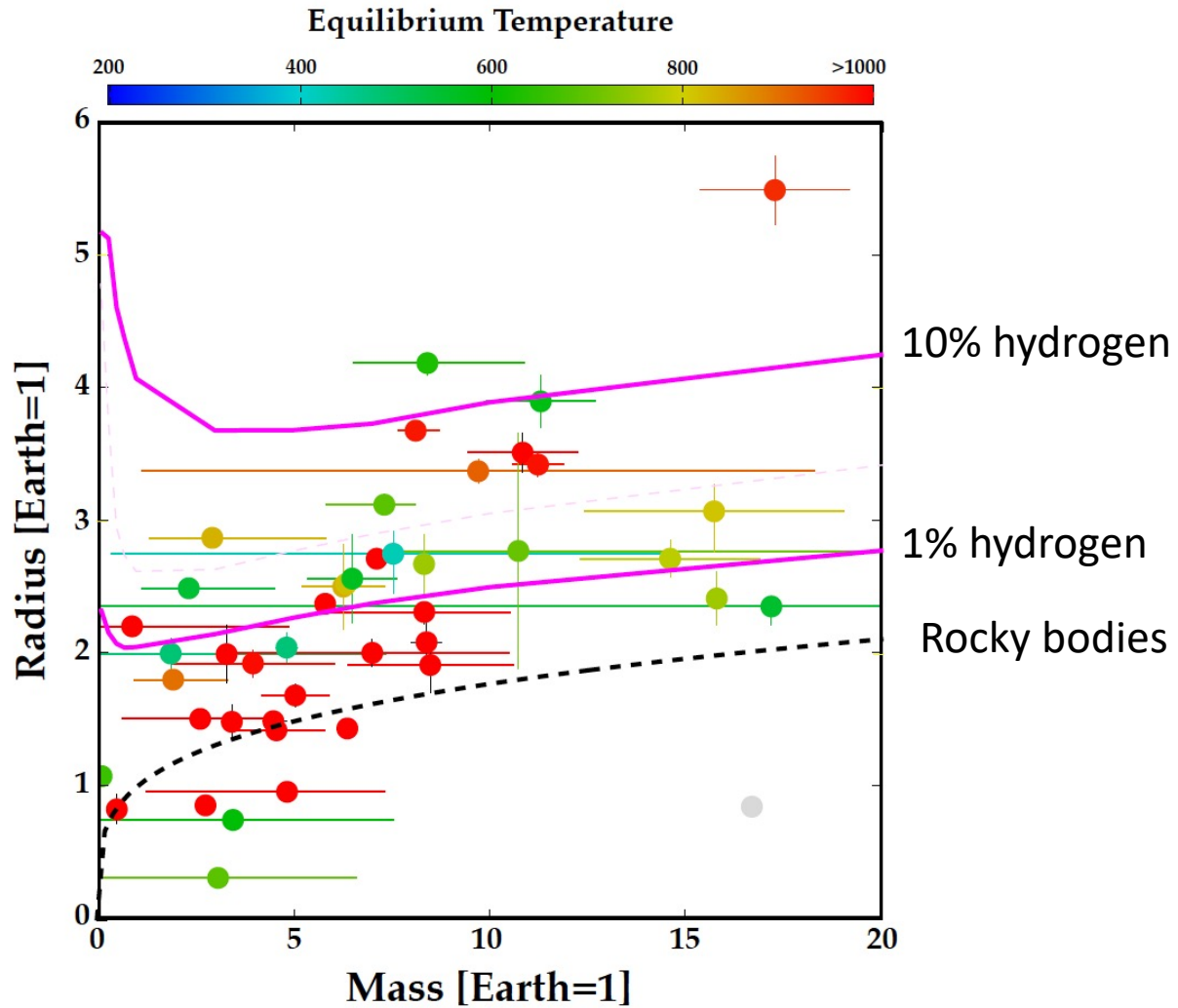
These ages compatible with accretion models

O'Brien, 2018

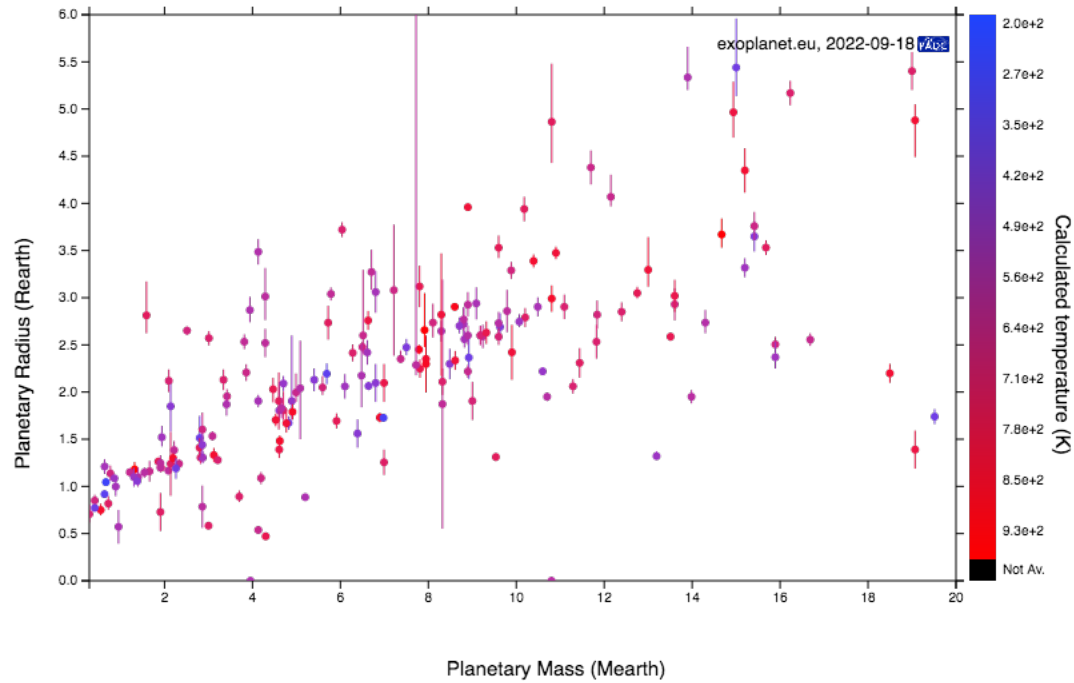
# Introduction / constraints:

Super Earth less than  $20 M_{\oplus}$

Diversity in H content



# Introduction / constraints:



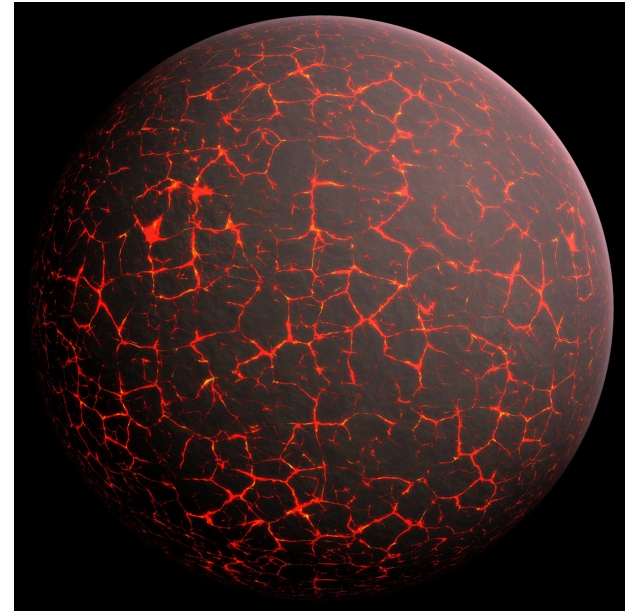


## Cause of magma ocean

- Blanketing by a dense primary atmosphere
- Impacts
- Radiogenic elements (Al<sub>26</sub>)

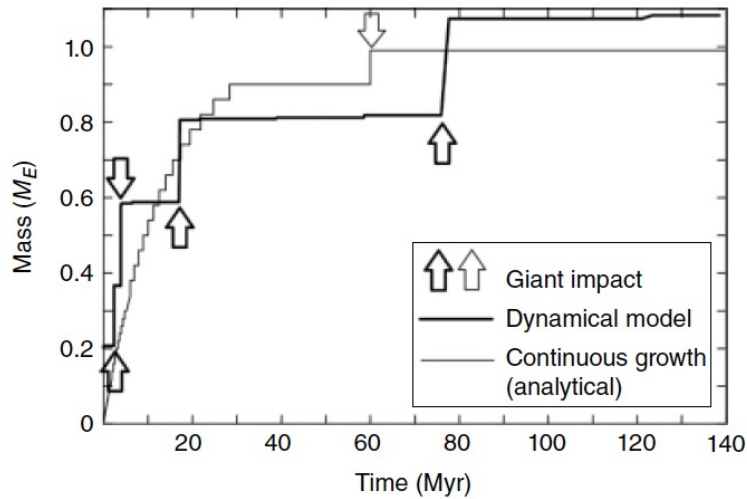


Image Credit: NASA/JPL-Caltech

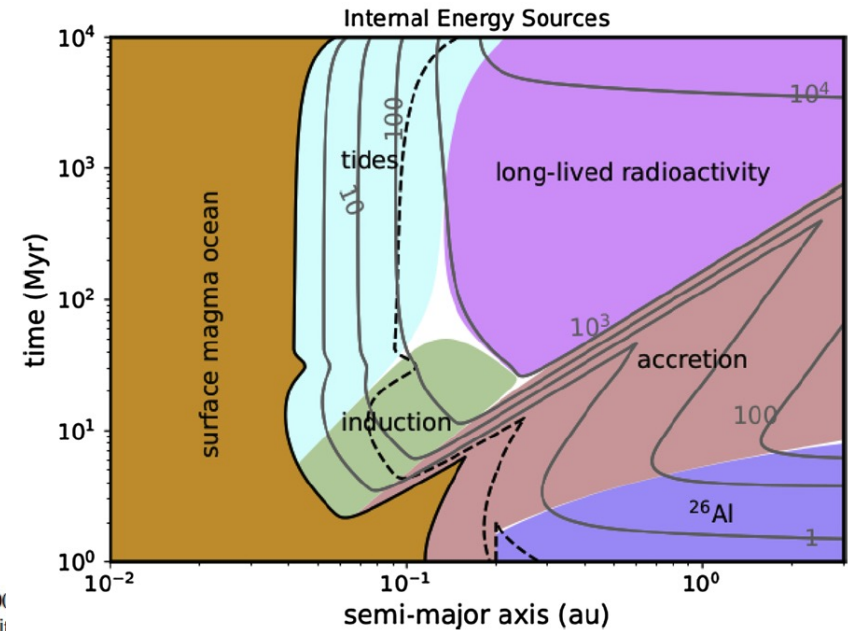


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<https://commons.wikimedia.org/w/index.php?curid=37984303>

# Introduction / constraints: Magma oceans



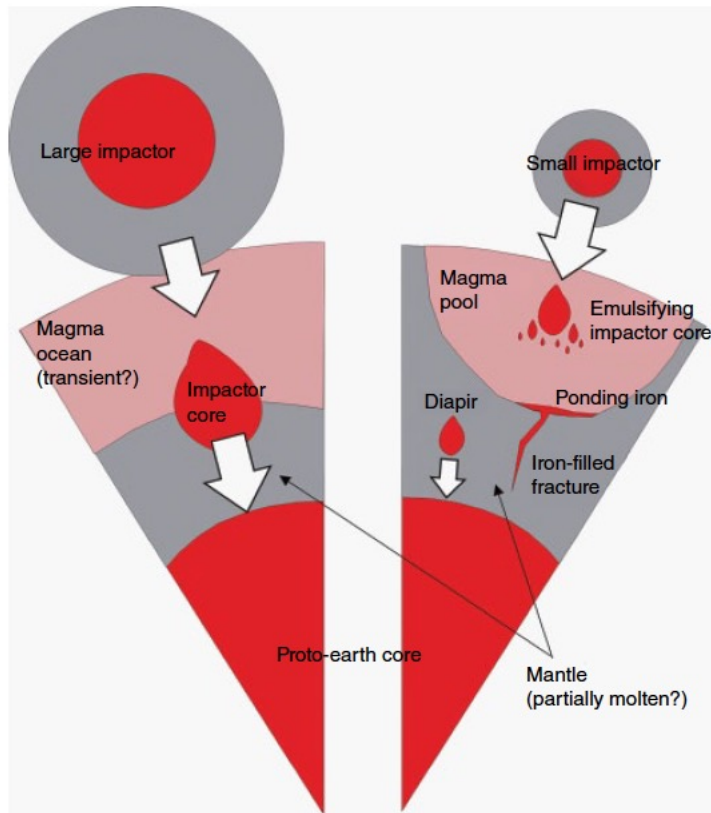
**Figure 5.1** Bold line shows a typical growth curve from an N-body simulation [run2a of *Raymond et al., 2001*], arrows indicate giant impacts. The thin line assumes growth at an exponentially decaying rate (equation 5.4) with timescale  $\tau$  of 10 Myr and a later final Moon-forming impact. Exponential growth models with longer e-folding timescales have been used to reproduce the Hf-W and U-Pb isotope systematics of the Earth [*Halliday, 2004; Rudge et al., 2010*] under the assumption of incomplete re-equilibration during impacts (Section 5.2.4.1).  $M_E$  is one Earth mass.



*Chao, 2021*

*Nimmo & Kleine 2015*

# Introduction / constraints: Magma oceans



Nimmo & Kleine 2015

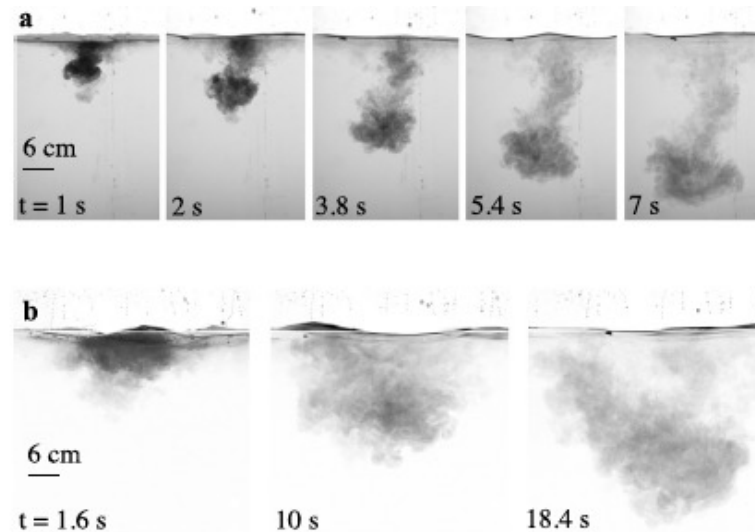
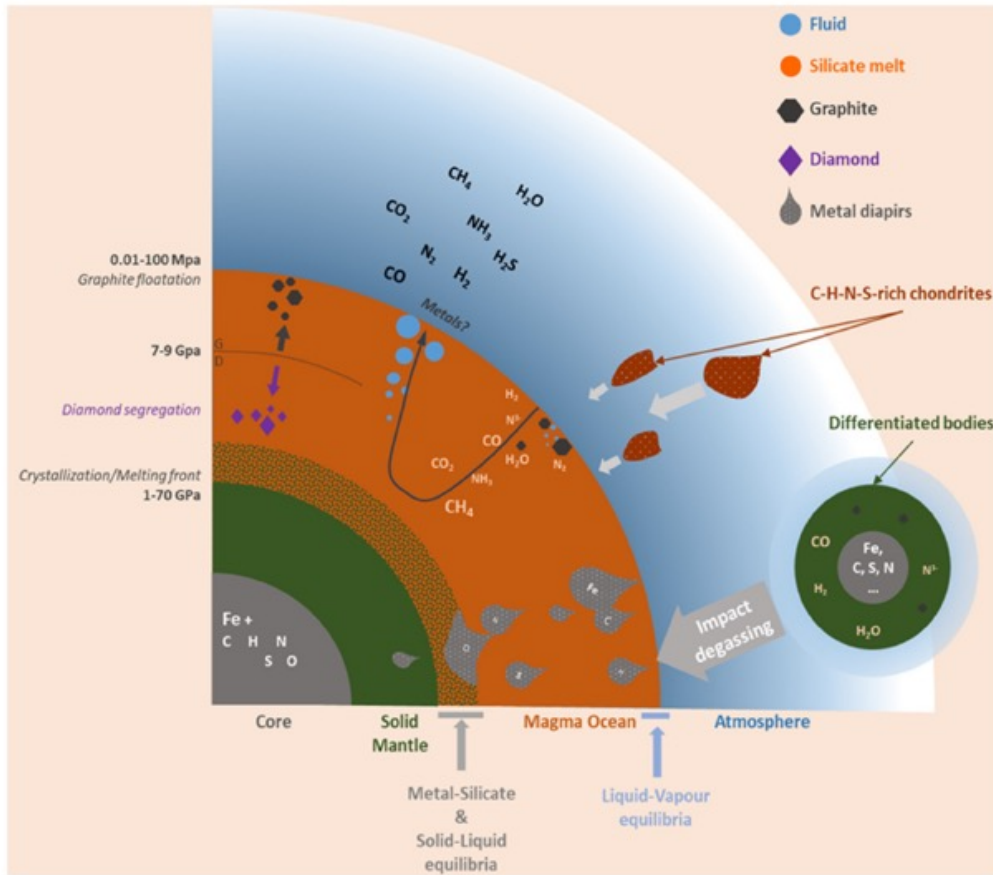


Fig. 3. Late-time thermal stage following the impact shown in Fig. 2. (a) Low impact Froude number,  $Fr \approx 6$ ,  $U \approx 1.3 \text{ ms}^{-1}$ . (b) High impact Froude number,  $Fr \approx 100$ ,  $U \approx 5.4 \text{ ms}^{-1}$ . In both experiments  $P \approx 0.02$  and the impactor radius  $R \approx 3 \text{ cm}$ . The experiments in (a) and (b) correspond to supplementary movies S1 and S2, respectively.

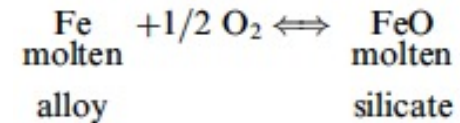
Landeau et al., 2021

# Introduction / constraints: Magma oceans

Certaines petites planètes sans atmosphères primitives vont pouvoir former une atmosphère secondaire Par dégazage.



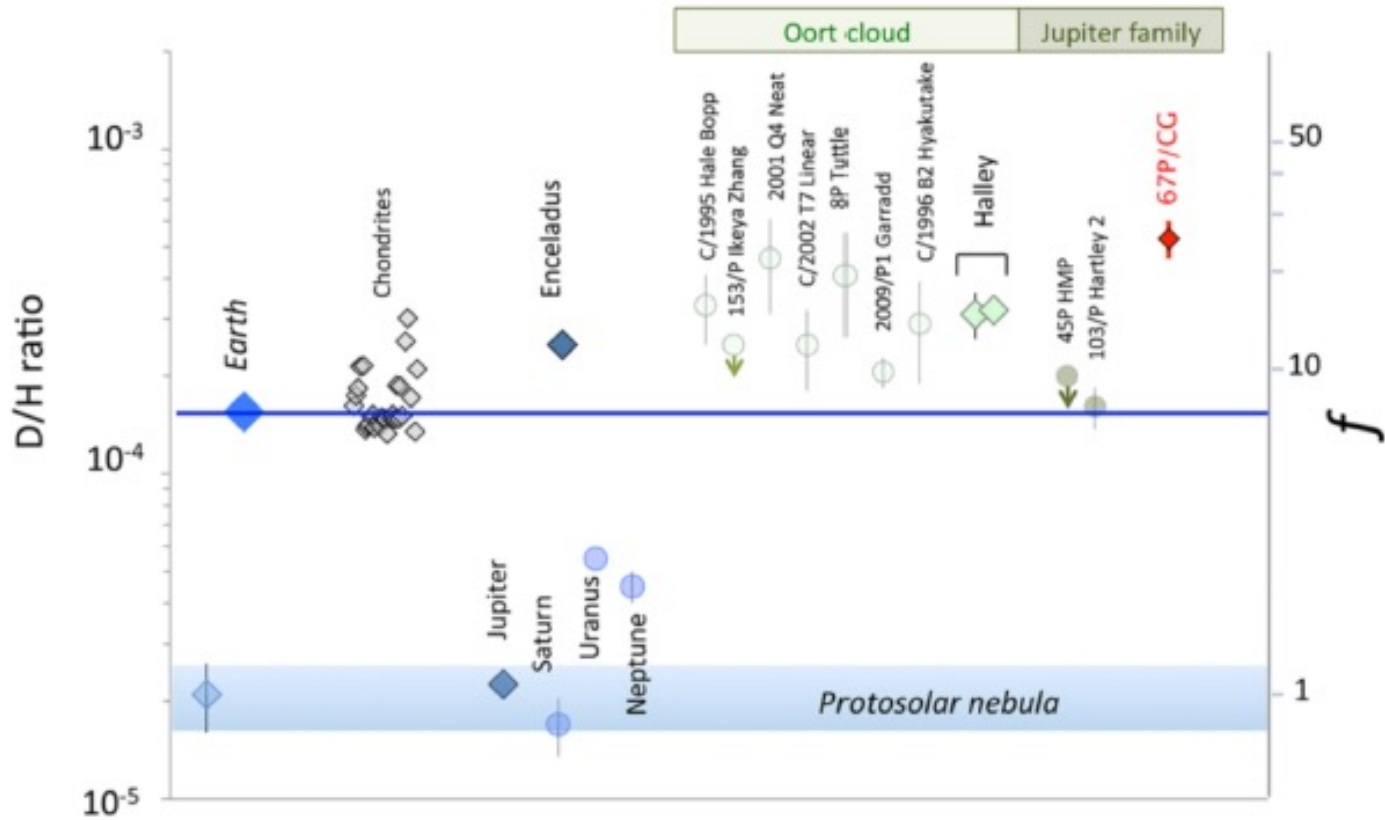
Gaillard et al. 2022



Hirschmann, 2022

IW:  $f_{\text{O}_2}$

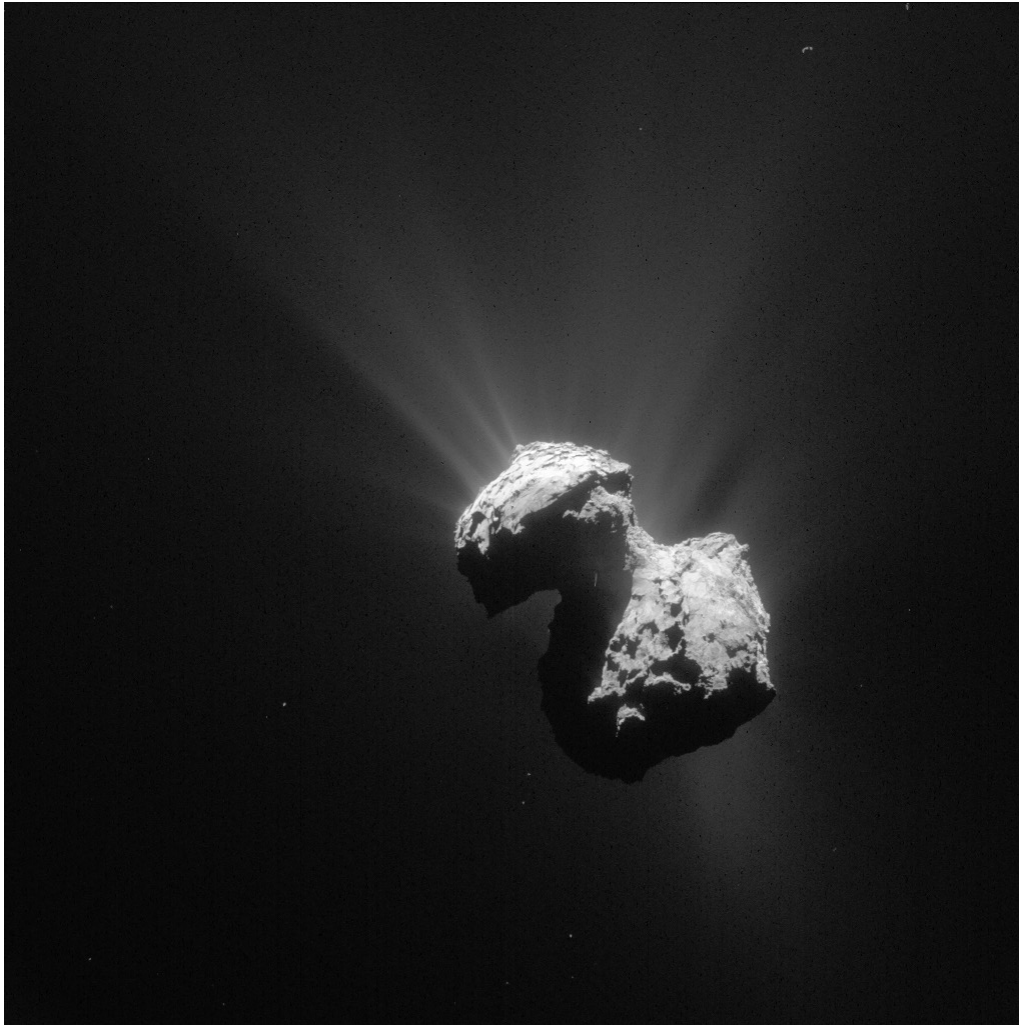
# Introduction / constraints: Volatiles



Altwegg et al., 2015

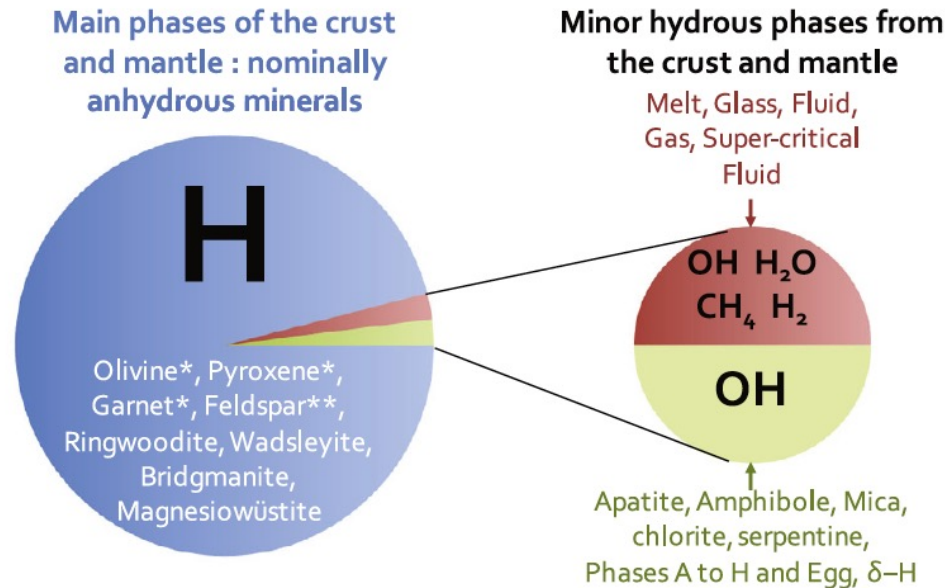


**CometWatch - NavCam images, ESA 7/7/2015**



Tchourioumov gerasymenko

~18.3 / +81 EO

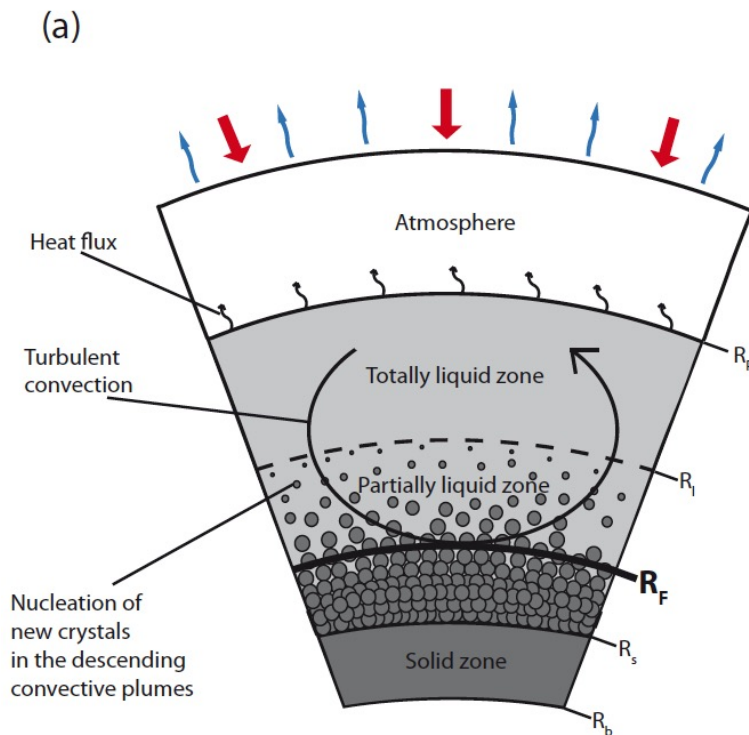


**Fig. 1** Sketch illustrating the “water” species present in the various phases of the Earth’s mantle and crust. The size of the pie slices represents the approximate volume percentage of the phases. The main reservoir of water in the mantle is nominally anhydrous minerals where hydrogen (H) enters their lattice in defects, and bonds to structural oxygen (*blue field*) (Bell and Rossman 1992b). \*Olivine, pyroxene and garnet can incorporate water as H<sub>2</sub> under reduced conditions (Yang et al. 2016). \*\*K-feldspar can sometimes include water as H<sub>2</sub>O and NH<sub>4</sub> (Johnson and Rossman 2004)

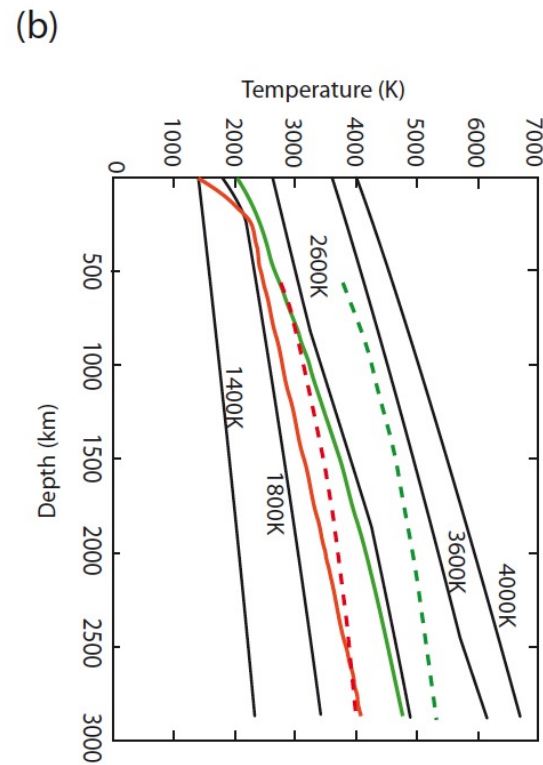
*Peslier et al., 2017*

- Introduction / constraints
- Heat and volatile transfer during MO Phase
- Examples, results and applications to exoplanets

# Heat and volatile transfer during Magma Ocean Phase



Legend: Net Solar flux Escape processes



$$SC = \int_{R_b}^{R_p} \rho C_p \frac{dT_p}{dt} r^2 dr = R_p^2 Q_S + R_B^2 Q_B + H_{int} + H_{sol}.$$

Massol et al., 2016

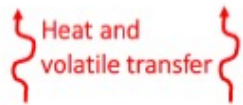
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# Heat and volatile transfer during Magma Ocean Phase

## Coupled model

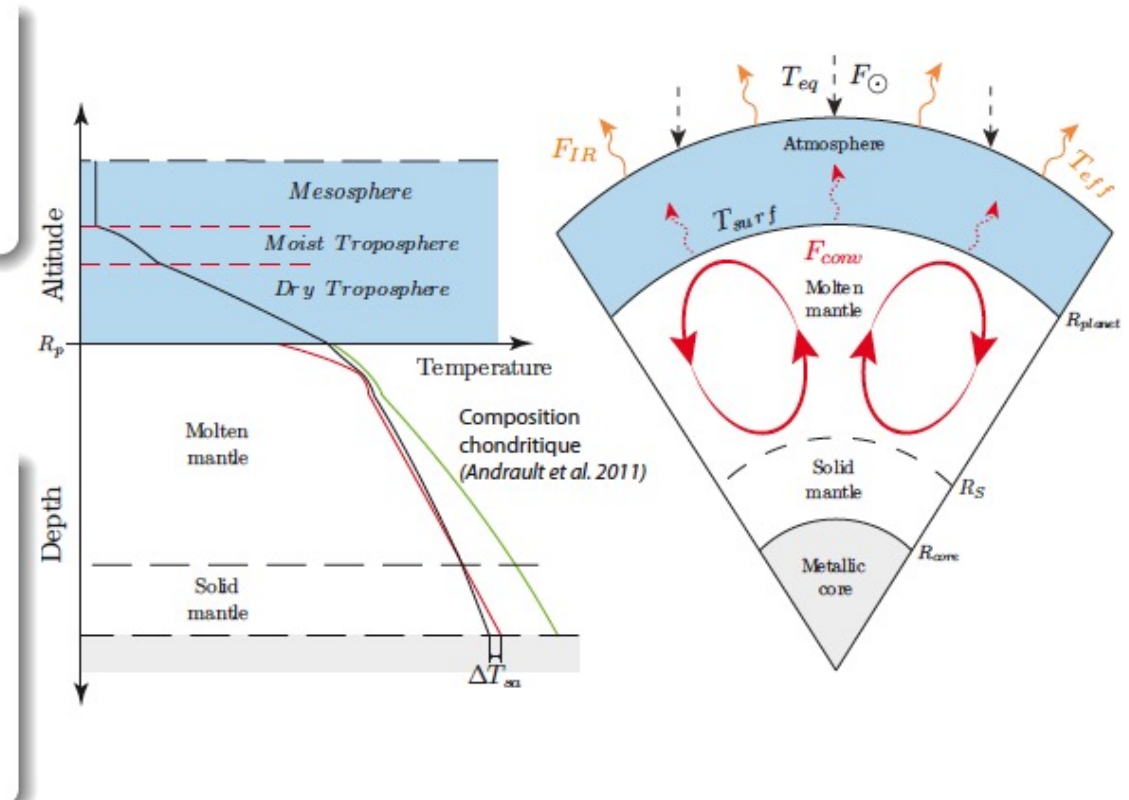
### ATMOSPHERE

- Radiative-convective (Marcq, 2012&2017)
- 1D Plan parallel



### MANTLE

- Parameterized convection,  $F_{conv} \propto Ra^{1/3}$
- Bottom up solidification



Salvador et al., 2017

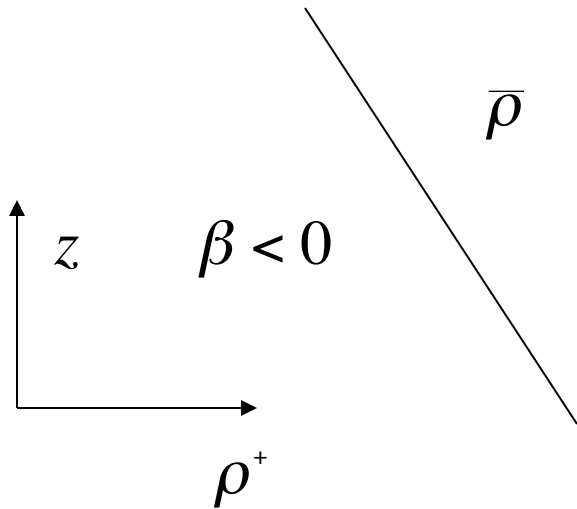


# Heat and volatile transfer during Magma Ocean Phase: Ra

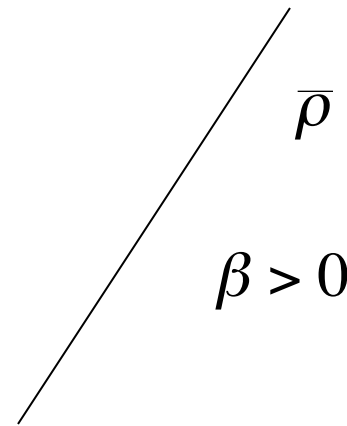
$$u \equiv 0$$

$$\bar{\rho} = \rho_0(1 + \beta z)$$

Statically stable



Statically unstable

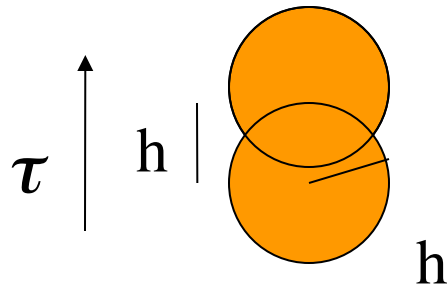


Important physical parameters:

Dimensions

$g$	$\nu$	$K$	$\beta$
$LT^{-2}$	$L^2T^{-1}$	$L^2T^{-1}$	$L^{-1}$

Spherical particle of radius  $h$   
Displaced upward of  $h$  in a time  $\tau$




- Characteristic velocity **U**?
- Potential energy gained **E**?
- How much **time** for the particle to lose buoyancy?
- What is the **energy D** dissipated by viscous forces?

Characteristic velocity :  $U \propto \frac{h}{\tau}$

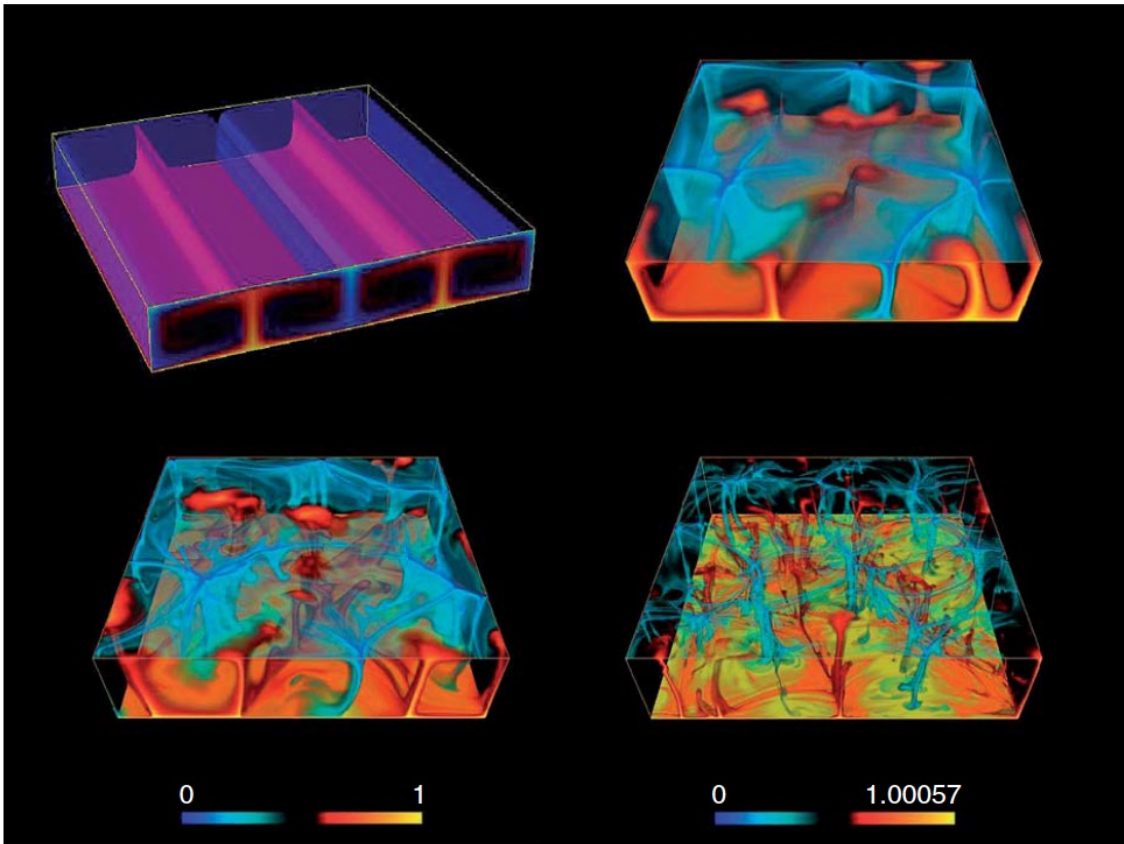
Potential energy gained :  $E \propto (\Delta\rho h^3)gh$   
 $\propto \rho_0\beta h^3 gh$   
 $\propto \rho_0\beta gh^5$

The particle lose buoyancy in a time:  $\tau \propto \frac{h^2}{\kappa}$

Energy D dissipated by viscous forces :  $D \propto \mu \frac{U}{h} h^2 h \propto \mu \frac{h^3}{\tau} \propto \mu \kappa h$

**Instability if E>D**   $\rho_0\beta gh^5 > \mu \kappa h$  i.e.:  $Ra = \frac{\beta gh^4}{\kappa \nu} > 1$

Convective patterns as a function of Rayleigh number:



$$\text{with } \beta = \frac{\alpha \Delta T}{h}$$

$$Ra = \frac{\alpha g \Delta T h^3}{\kappa \nu}$$

up to  $Ra=30$  for entirely liquid MO

**Figure 4** Convection patterns of a fluid heated from below at Rayleigh number  $10^5$ ,  $10^6$ ,  $10^7$ ,  $10^8$ . The temperature color bars range from 0 (top boundary) to 1 (bottom boundary). The Boussinesq approximation was used (numerical simulations by F. Dubuffet). The increase in Rayleigh number corresponds to a decrease of the boundary layer thicknesses and the width of plumes. Only in the case of the lowest Rayleigh number (top left) is the convection stationary with cells of aspect ratio  $\sim \sqrt{2}$  as predicted by marginal stability. For higher Rayleigh number, the patterns are highly time dependent.

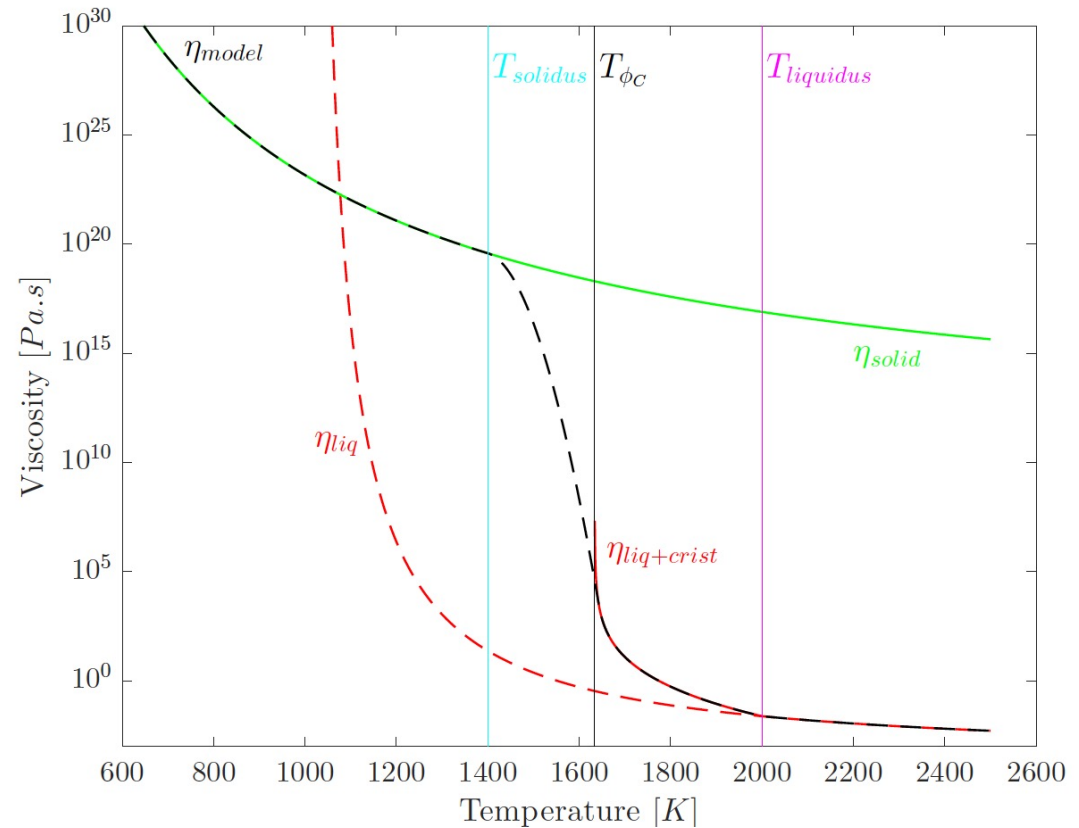


# Heat and volatile transfer during Magma Ocean Phase

## Viscosity

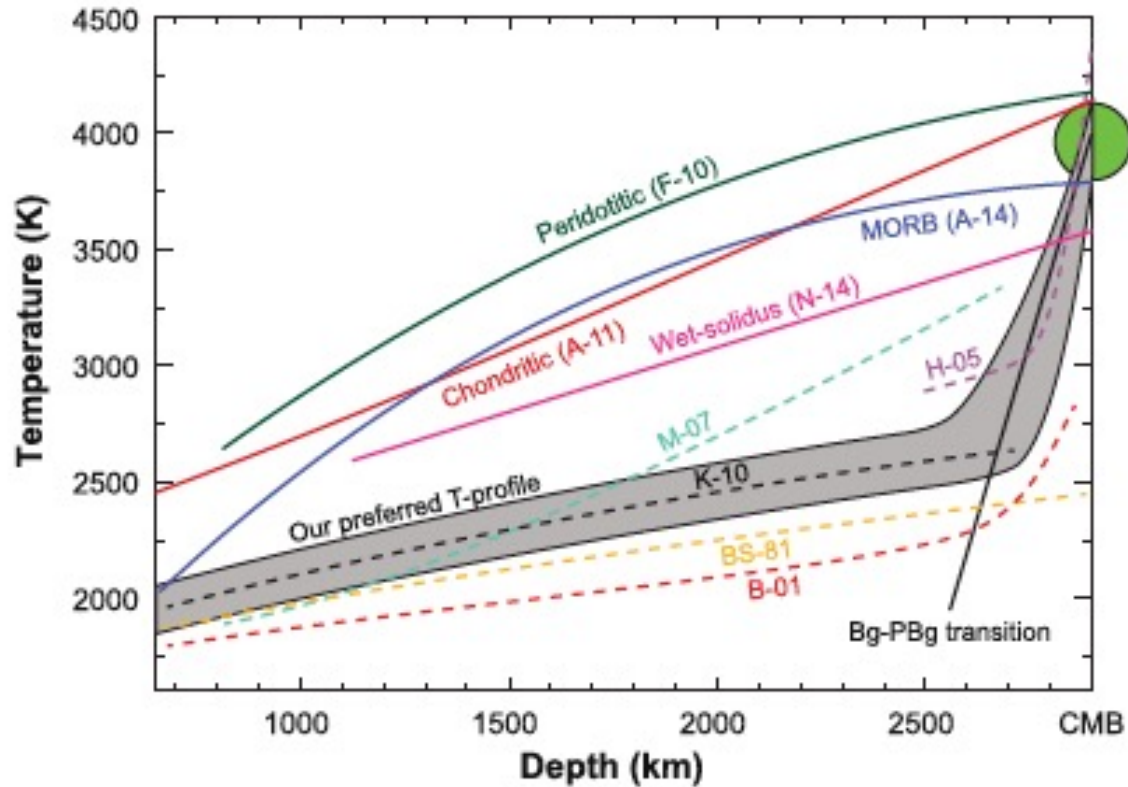
After Takei and Holtzmann 2009

Interpolation between  $\eta_s$  (Karato et Wu, 1993)  
et  $\eta_l$  (Karki et Stixrude, 2010)



Salvador et al., 2017

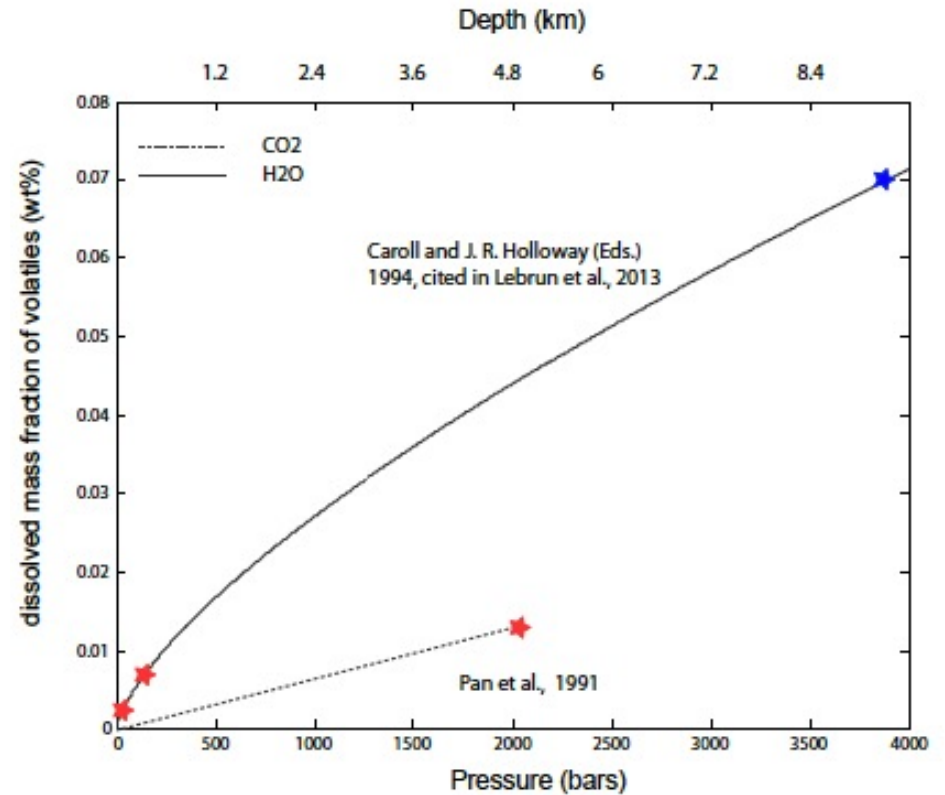
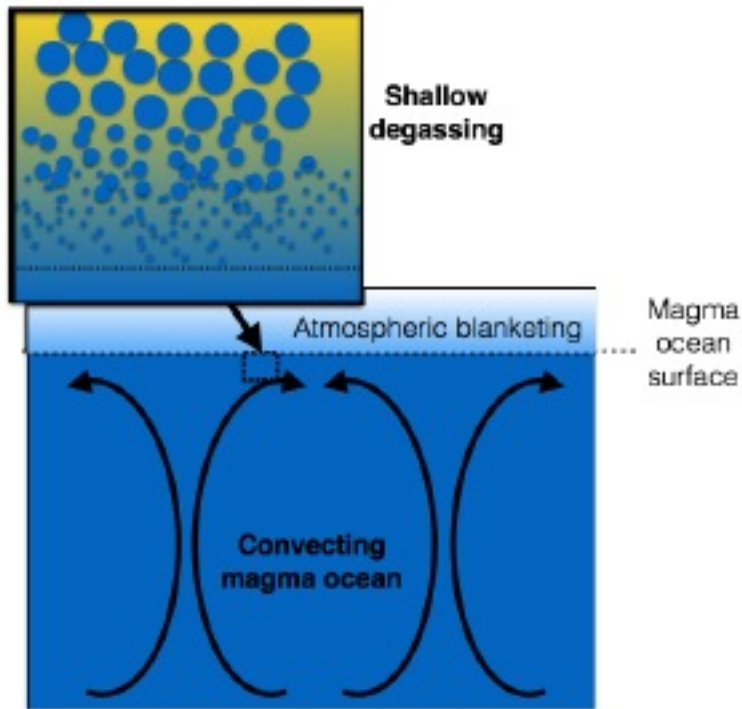
# Heat and volatile transfer during Magma Ocean Phase



Andraut et al., 2016

# Heat and volatile transfer during Magma Ocean Phase

## Exsolution



*More references in Iacono et al., 2012*

Mass conservation of volatiles

$$M_{vol}^{solid} + M_{vol}^{liquid} + M_{vol}^{atm} = X_{vol,t=0} M_{t=0}^{MO}$$

$$M_{vol}^{atm} = \frac{4\pi R_p^2}{g} \left( \frac{\mu_{vol}}{\bar{\mu}} \right) P_{vol}(X_{vol}^{liquid})$$

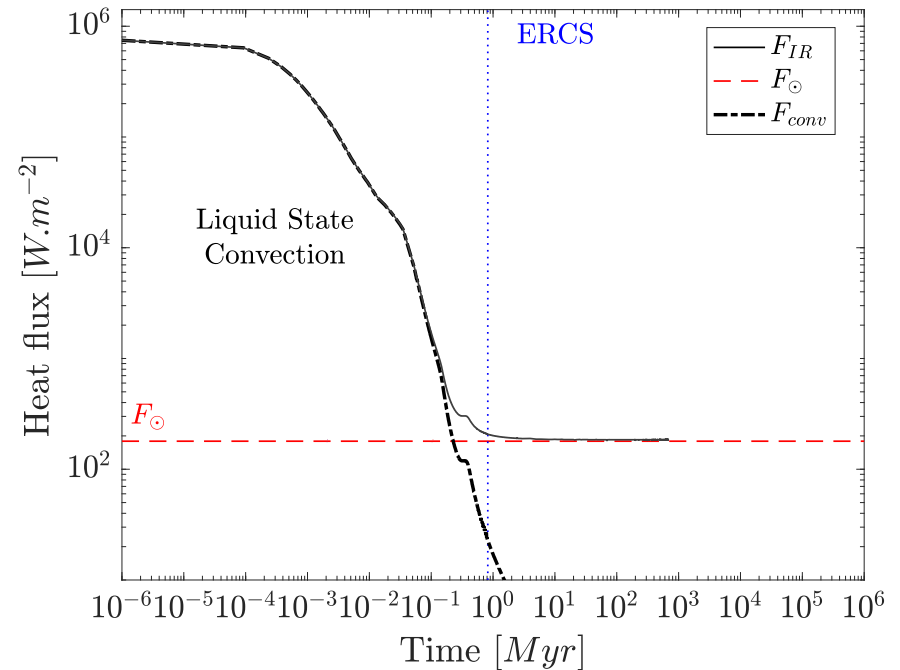
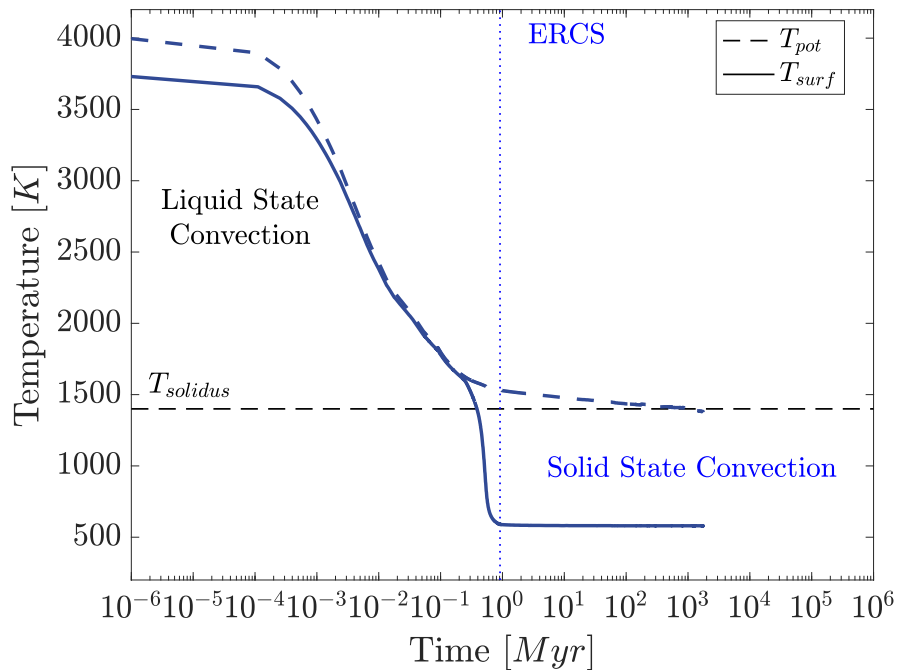
## Definition: End of Rapid Cooling Stage

- Limit between two evolution stages:
  - Controlled by  $F_{conv}$ , vigorous mantle convection, strong cooling, degassing of the atmosphere
  - Controlled by  $F_{\odot}$ , quasi steady-state:  $T_{surf}$  and  $P_{atm}$ , other geodynamic regime
- ERCS taken at time when  $F_{conv} = \frac{1}{10} F_{\odot}$



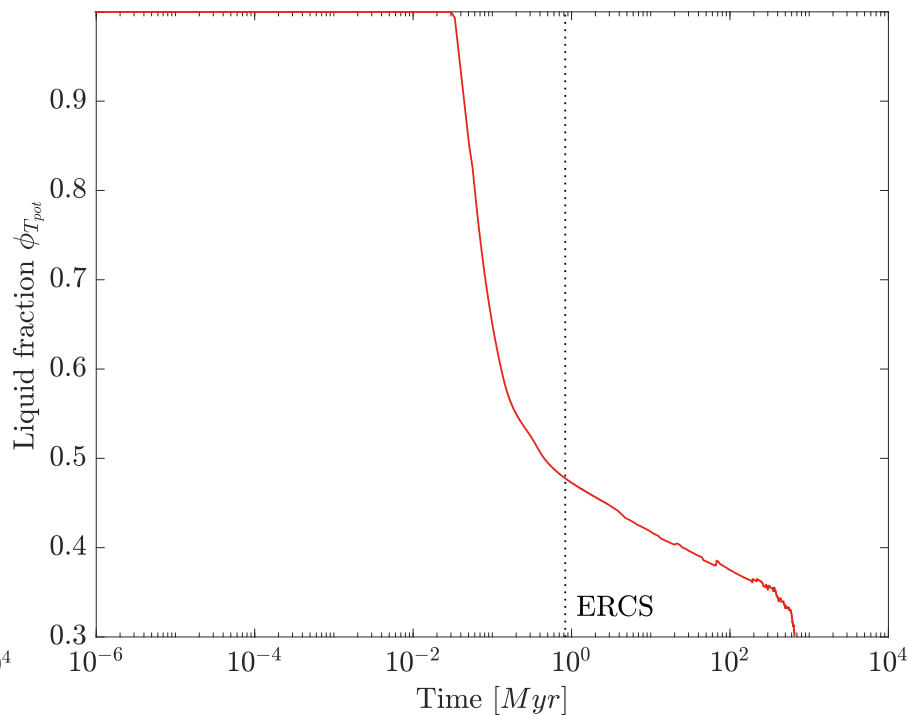
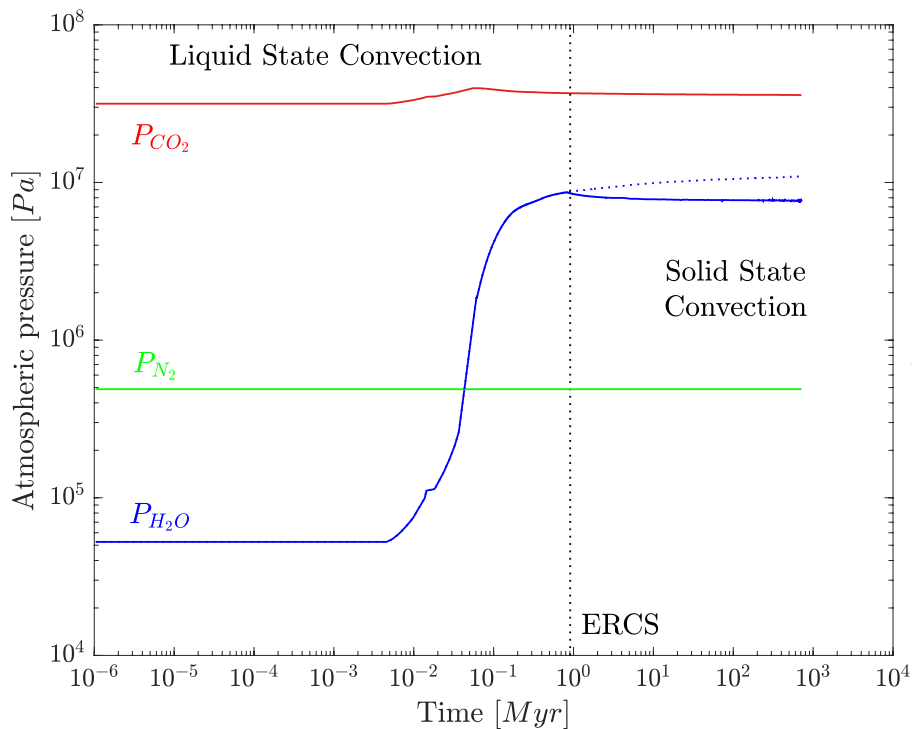
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## Time evolution



CO<sub>2</sub>: 600 ppm; H<sub>2</sub>O: 0.4 EO; D=1 AU  $\alpha=0.2$

## Degassing



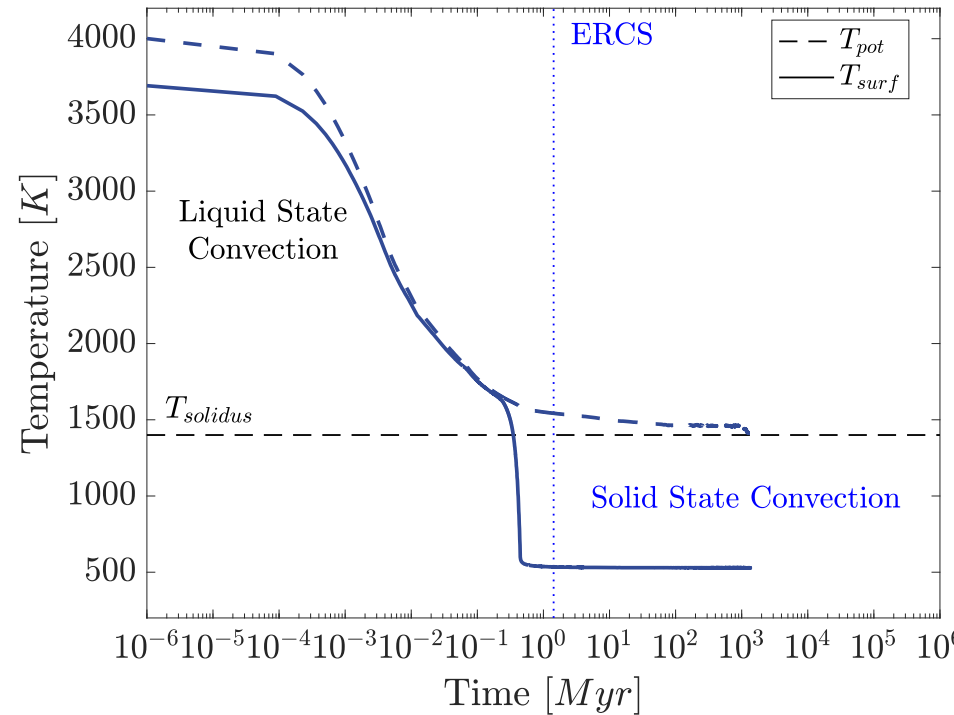
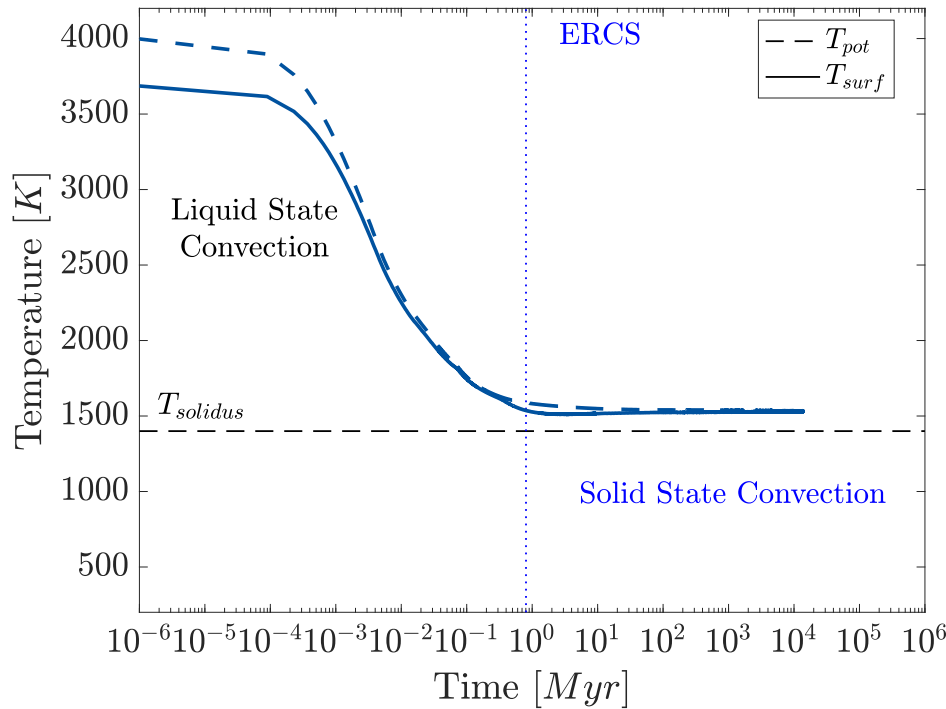
$CO_2$ : 600 ppm;  $H_2O$ : 0.4 EO;  $D=1$  AU  $\alpha=0.2$

# Examples, results and applications

Distance to the sun

0,77

1,3



End of Rapid Cooling Stage

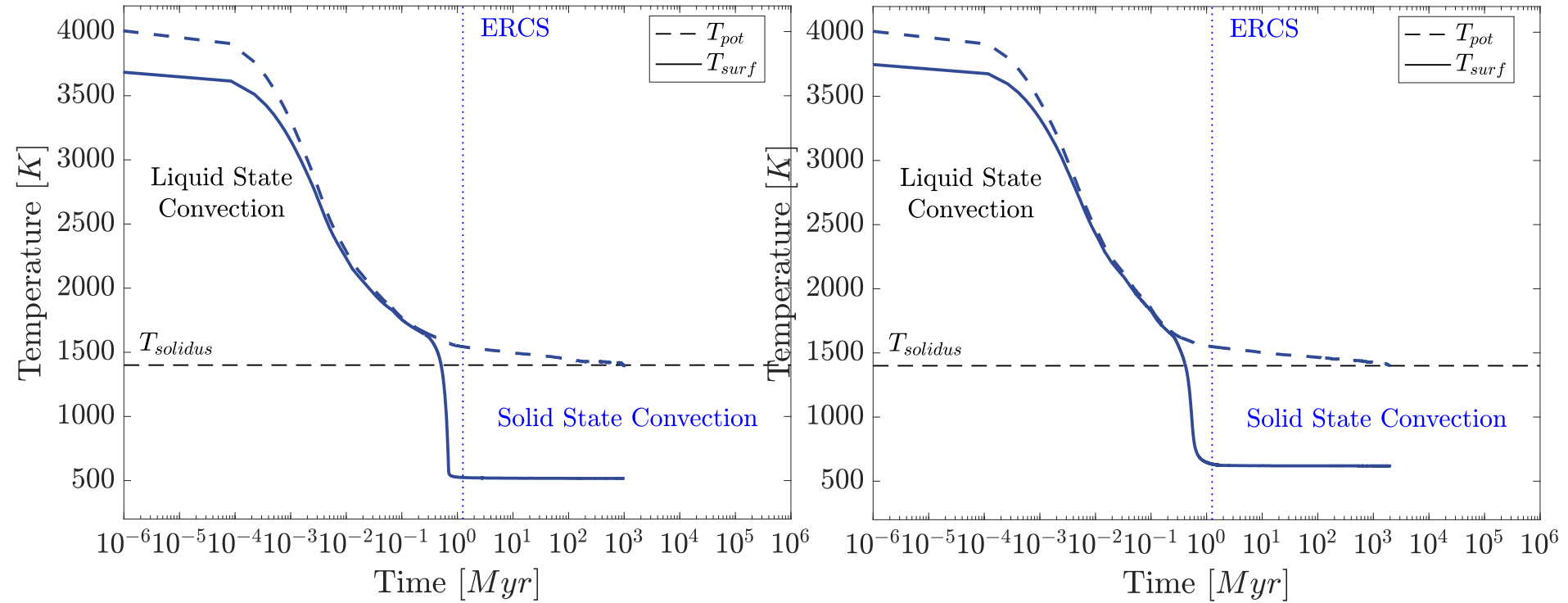
ERCS :  $F_{conv} \leq 0.1 \times F_{\odot}$

Volatiles

$H_2O = 0.6$  MEO,  $CO_2 = 500$  ppm

# Examples, results and applications

## CO<sub>2</sub> content

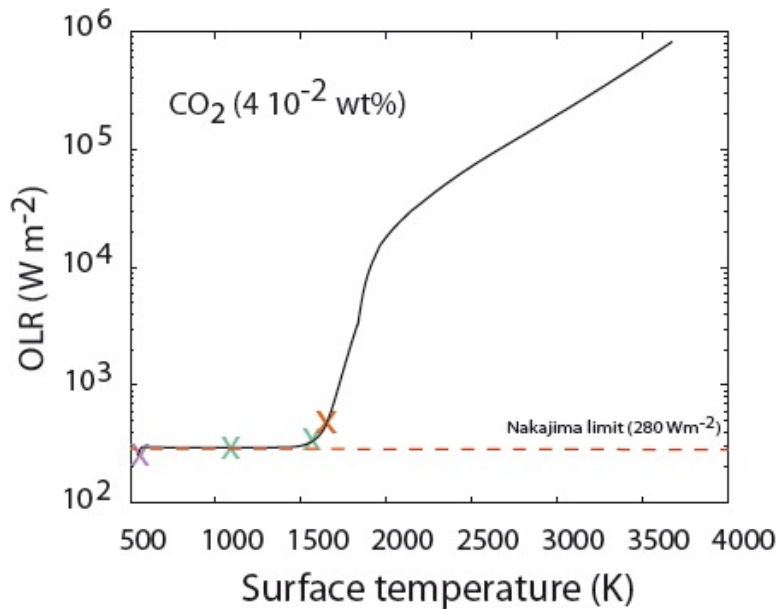


CO<sub>2</sub>: 400 ppm; H<sub>2</sub>O: 0.6 EO; D=1 AU  $\alpha=0.2$

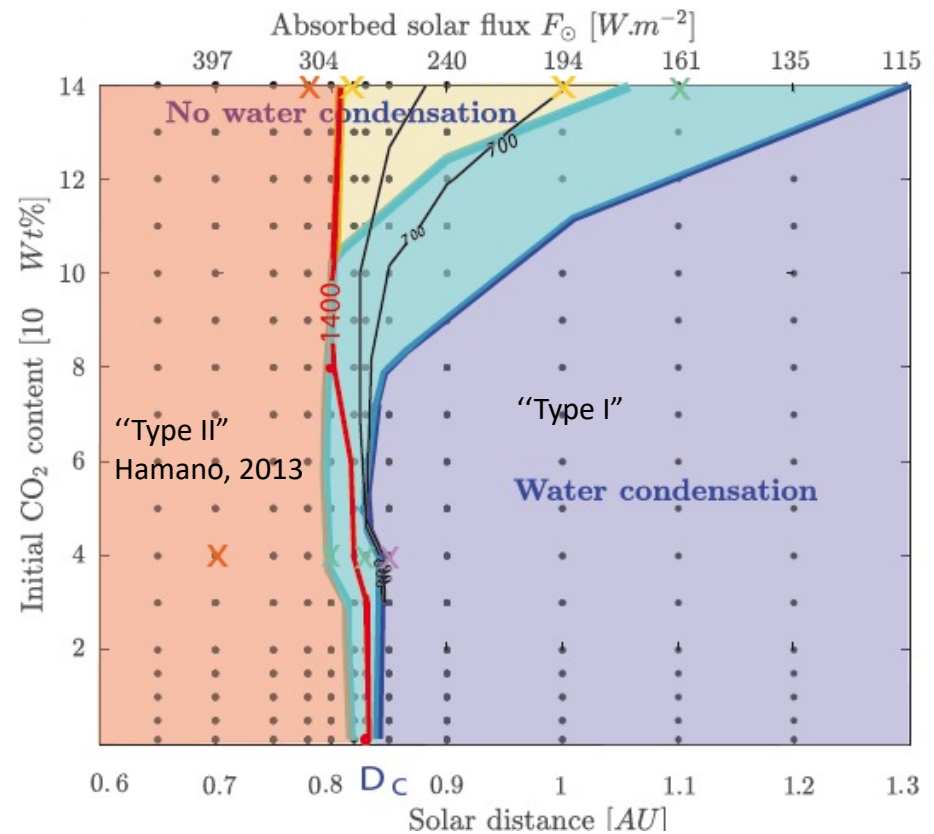
CO<sub>2</sub>: 1200 ppm; H<sub>2</sub>O: 0.6 EO; D=1 AU  $\alpha=0.2$

# Examples, results and applications

$H_2O = 0.6$  MEO,  $\alpha = 0.2$



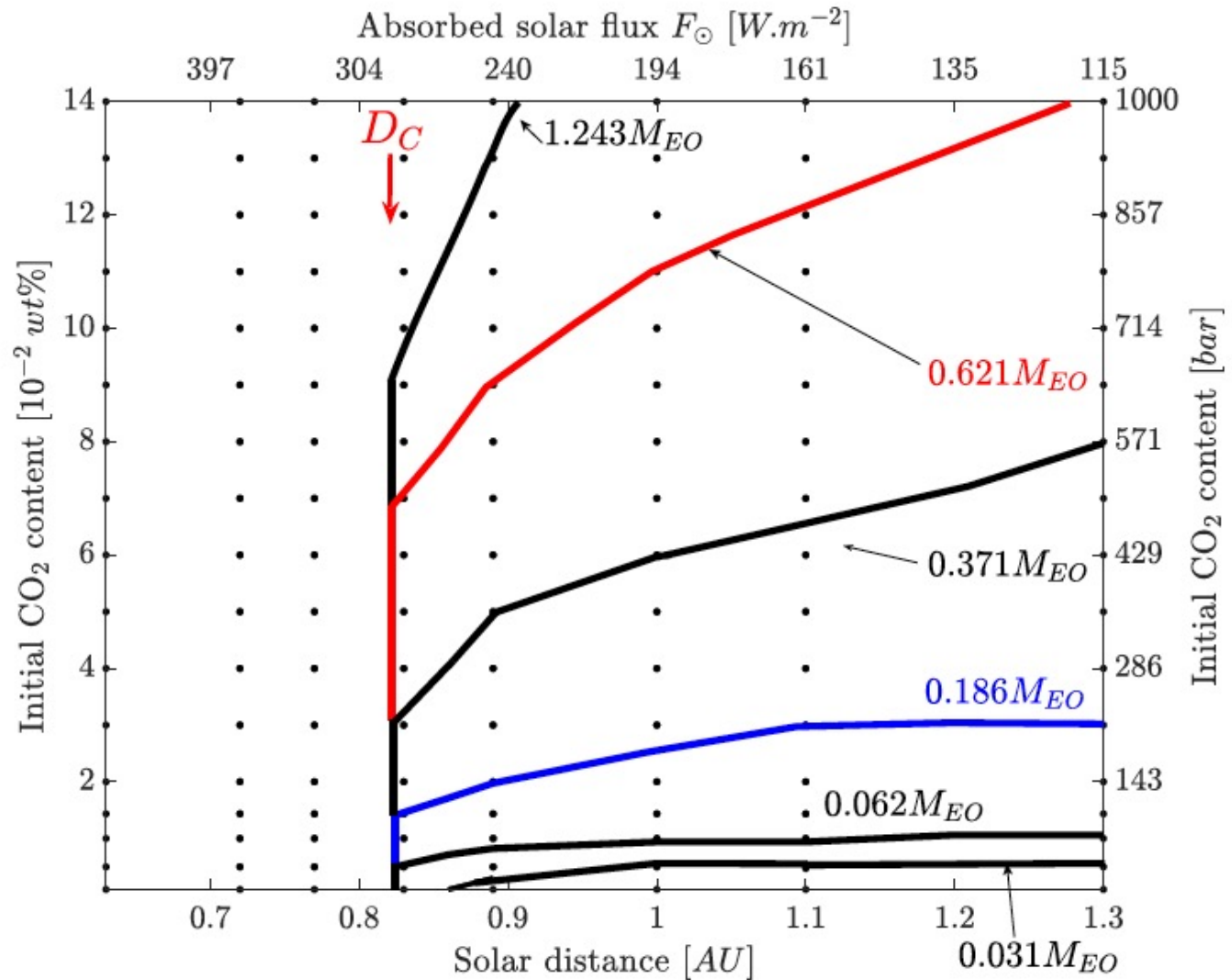
- partial melt
- no condensation
- condensation post ERCS
- condensation at ERCS





# Examples, results and applications

For different initial water content



## Scaling

volatiles

$$x^* = \frac{x[0]_{CO_2}}{x[0]_{CO_2} + x[0]_{H_2O}}$$

Distance to the star

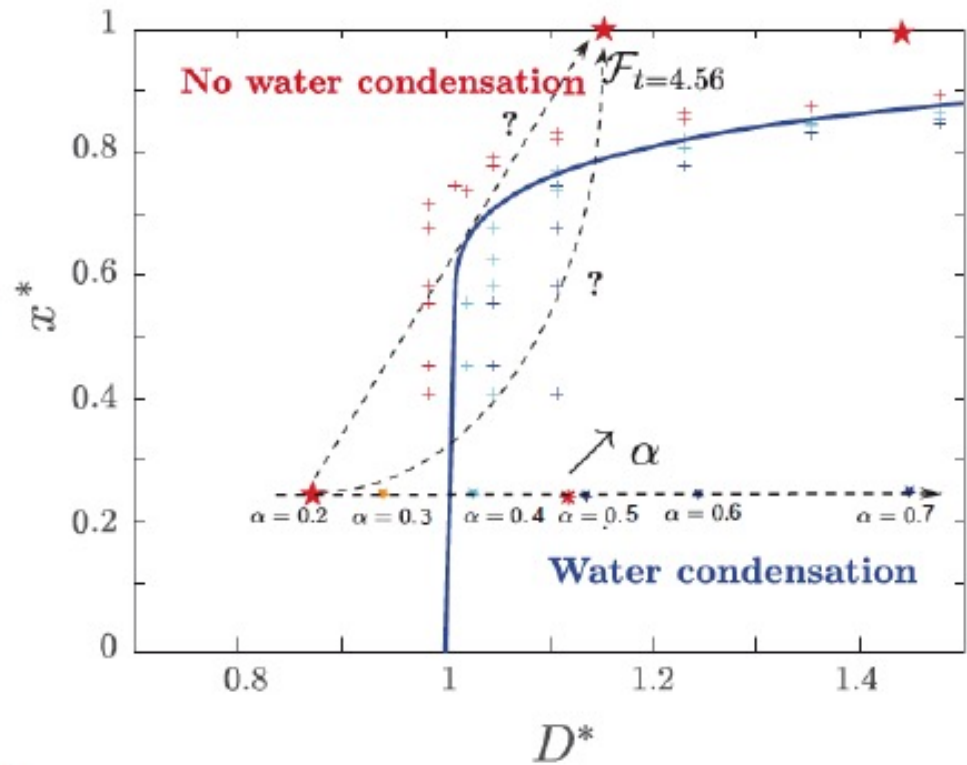
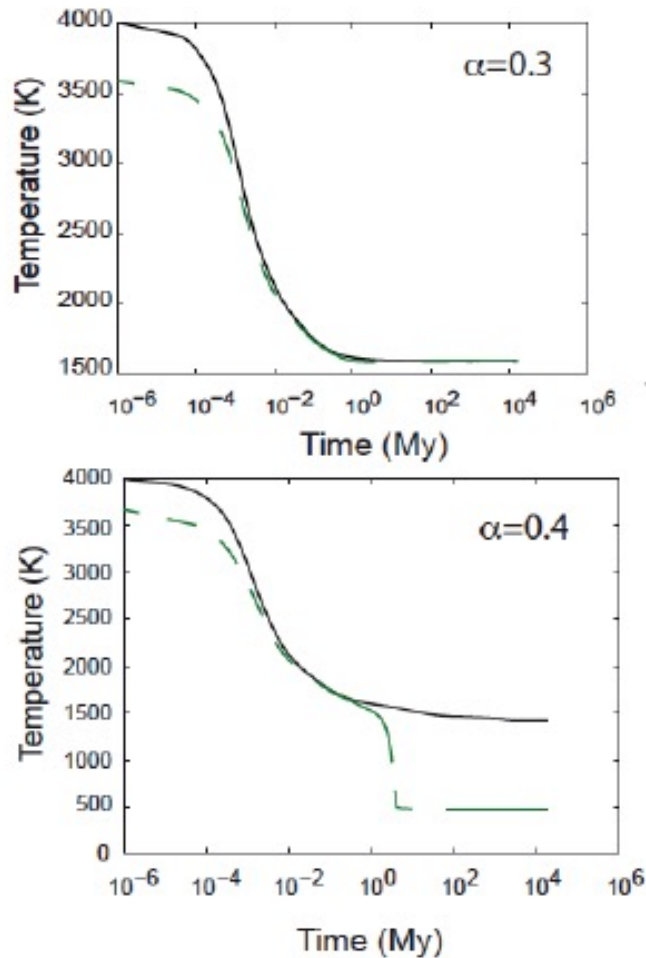
$$D_C = \sqrt{\frac{F_0(1-\alpha)}{F_{IR}}}$$

$F_0$  being the solar constant for the young sun and  $F_{IR}$  the Nakajima limit ( $280 \text{ W m}^{-2}$   
In Marcq et al. model used here)

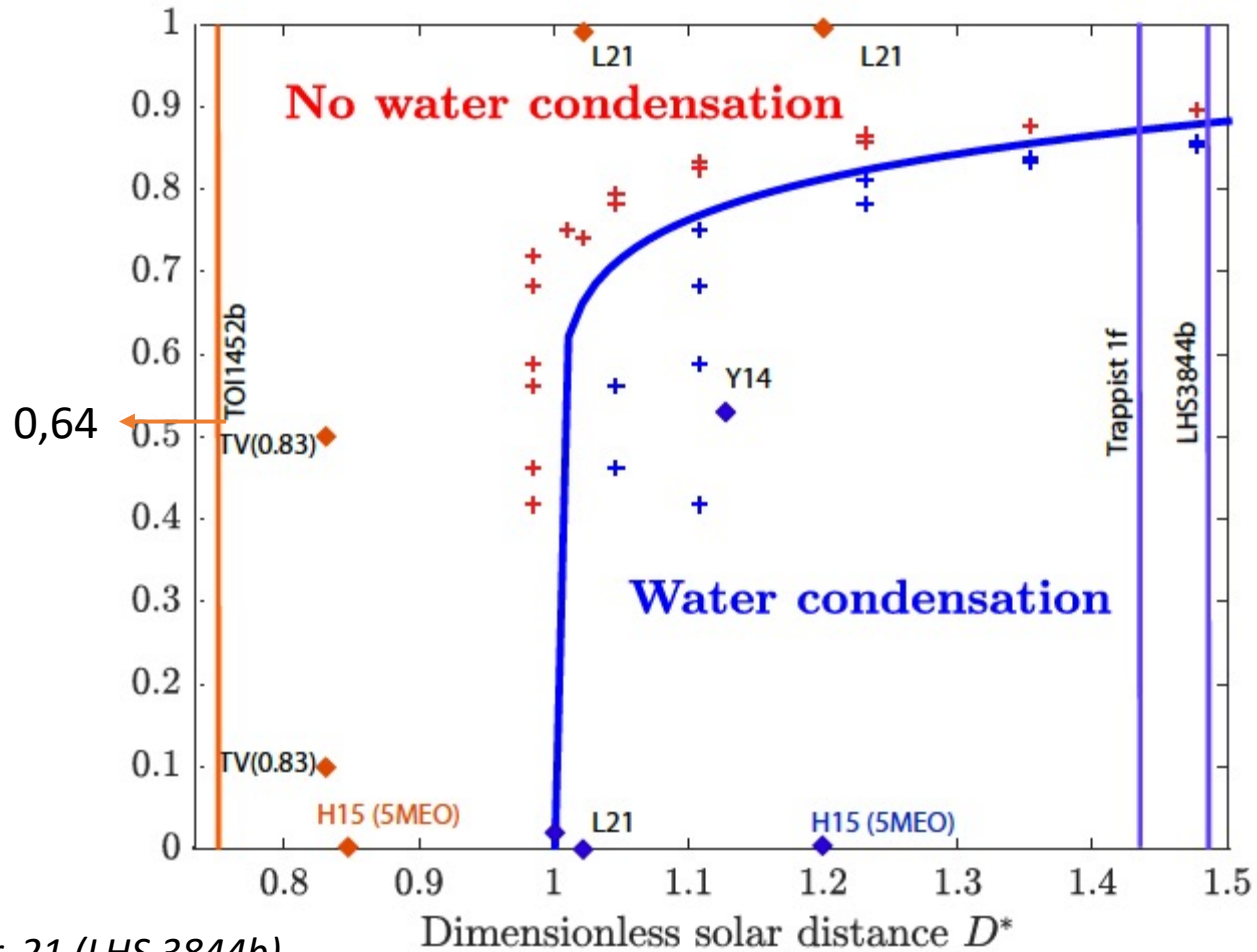
$$D^* = \frac{D}{D_C} = 0.814$$

## Venus

For a 0.4 albedo it requires 0.13 Earth Ocean ( $x^* \approx 0.65$ )



# Examples, results and applications



Meier, 21 (LHS 3844b)

Dorn and Lichtenberg, 21 (Trappist 1f)

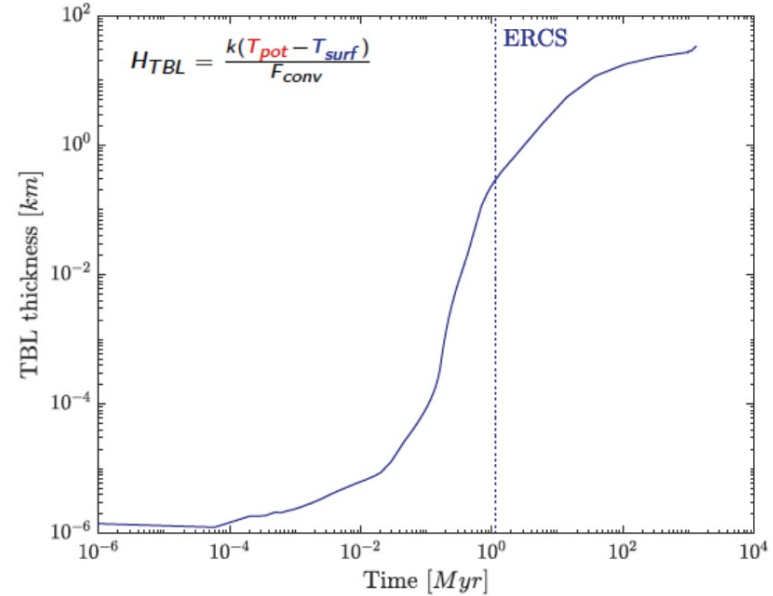
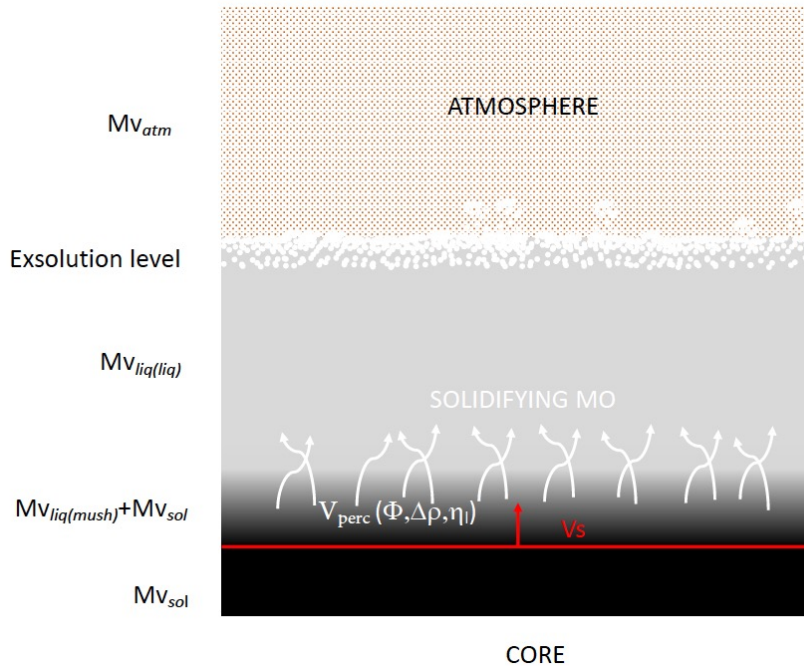
Cadieux, 22 (TOI1452b)

Massol et al. in prep.

# Conclusions perspectives

Perspectives:

- Post MO Heat and volatile transfer
- Surface conditions and fracturation



Application to exoplanets: remarks, questions?